

LONG-SLIT SPECTROSCOPY OF PARSEC-SCALE JETS FROM DG TAURI

HEEYOUNG OH^{1,2}, TAE-SOO PYO^{3,4}, IN-SOO YUK², AND BYEONG-GON PARK^{1,2}

¹Korea University of Science and Technology, 217 Gajeong-ro, Yuseong-gu, Daejeon 305-350, Korea
hyoh@kasi.re.kr, bgpark@kasi.re.kr

²Korea Astronomy and Space Science Institute, 776 Daedeokdae-ro, Yuseong-gu, Daejeon 305-348, Korea
yukis@kasi.re.kr

³Subaru Telescope, National Astronomical Observatory of Japan, 650 North A'ohoku Place, Hilo, HI 96720, USA
pyo@subaru.naoj.org

⁴School of Mathematical and Physical Science, SOKENDAI (The Graduate University for Advanced Studies), Hayama, Kanagawa 240-0193, Japan

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Abstract: We present observational results from optical long-slit spectroscopy of parsec-scale jets of DG Tau. From HH 158 and HH 702, the optical emission lines of H α , [O I] $\lambda\lambda$ 6300, 6363, [N II] $\lambda\lambda$ 6548, 6584, and [S II] $\lambda\lambda$ 6716, 6731 are obtained. The kinematics and physical properties (i.e., electron density, electron temperature, ionization fraction, and mass-loss rate) are investigated along the blueshifted jet up to 650'' distance from the source. For HH 158, the radial velocity ranges from -50 to -250 km s⁻¹. The proper motion of the knots is $0''.196 - 0''.272$ yr⁻¹. The electron density is $\sim 10^4$ cm⁻³ close to the star, and decreases to $\sim 10^2$ cm⁻³ at 14'' away from the star. Ionization fraction indicates that the gas is almost neutral in the vicinity of the source. It increases up to over 0.4 along the distance. HH 702 is located at 650'' from the source. It shows ~ -80 km s⁻¹ in the radial velocity. Its line ratios are similar to those at knot C of HH 158. The mass-loss rate is estimated to be about $\sim 10^{-7} M_{\odot}$ yr⁻¹, which is similar to values obtained from previous studies.

Key words: stars: formation — stars: outflows — Herbig-Haro objects — stars: individual: DG Tau — techniques: spectroscopy

1. INTRODUCTION

In the formation of stars and planetary systems, studying outflows and mass-accretion is essential (Hartigan et al. 1995). Through the analysis of the kinematics and physical conditions of outflows from young stellar objects (YSOs), we are able to understand the jet launching mechanism and the interaction with ambient material. Parsec-scale jets are useful to study star formation in somewhat long time scale and large spatial distribution. There are a number of parsec-scale jets known (e.g., Ray 1987; Bally & Devine 1994; Ogura 1995; Bally et al. 1995). One of the extreme cases is the HH 222 system, suggested as a giant Herbig-Haro flow spread over 5.3 parsec in length (Reipurth et al. 2013).

Mundt & Fried (1983) discovered a jet-like outflow from DG Tau. DG Tau is an early evolutionary stage CTTS (Class II) with mass of $0.67 M_{\odot}$ (Hartigan et al. 1995) and age of about 3×10^5 yr (Beckwith et al. 1990). McGroarty & Ray (2004) suggested that it is the driving source of a parsec-scale jet including HH 158, HH 702, and HH 830. The later two HH objects are located at more than 9' away from DG Tau. In their follow up observations, McGroarty et al. (2007) reported that HH 702 showed receding motion with respect to DG Tau. They also concluded that HH 830 is not ejected from DG Tau. Eisloffel & Mundt (1998)

showed that the HH 158 outflow extends to $\sim 11''$ toward the southwest direction and estimated the direction of the proper motion of knots from DG Tau as $\sim 226 \pm 7^{\circ}$. They also calculated the jet inclination angle as $\sim 32^{\circ}$ with respect to the line of sight. In addition to an optical jet, Rodríguez et al. (2012) and Lynch et al. (2013) detected a radio knot at $\sim 7''$ from the central source. High spatial resolution observations with a space telescope and ground based large telescope with an adaptive optics system showed that the blueshifted HH 158 jet can be traced to within $\sim 0''.1$ from the star (Bacciotti et al. 2000, 2002; Pyo et al. 2003; Maurri et al. 2014). Lavalley et al. (1997) reported that the redshifted jet with the length of $\sim 1''$. Pyo et al. (2003) showed that there is a $0''.7$ gap between the redshifted jet and the star in the [Fe II] $1.644 \mu\text{m}$ emission, which indicates the presence of an optically thick disk. The two velocity components are shown in the optical and near infrared forbidden emission lines (Hartmann & Raymond 1989; Solf & Böhm 1993; Hamann 1994; Hartigan et al. 1995; Pyo et al. 2003, 2006).

In this paper, we report the results of optical long-slit spectroscopy on HH 158 and HH 702 outflows. We investigate the electron density, electron temperature, and ionization fraction from close vicinity of the central source to the knot at 14'' of HH 158 jet and HH 702 at 650'' from DG Tau. We note that this is the first time to obtain optical spectroscopy of the HH 702 outflow. We also summarize the proper motions of the knots in

Table 1
Log of Observations.

Object	Date (UT)	P.A. ($^{\circ}$)	Exposure time (s)	Slit width ($''$)	Spectral resolution ($\lambda/\Delta\lambda$)	Spectral coverage (\AA)
HH 158	2012 Nov. 15	223	1800	2.9	2070	5900–7050
HH 702	2012 Nov. 15	193	2800	2.9	2070	5900–7050

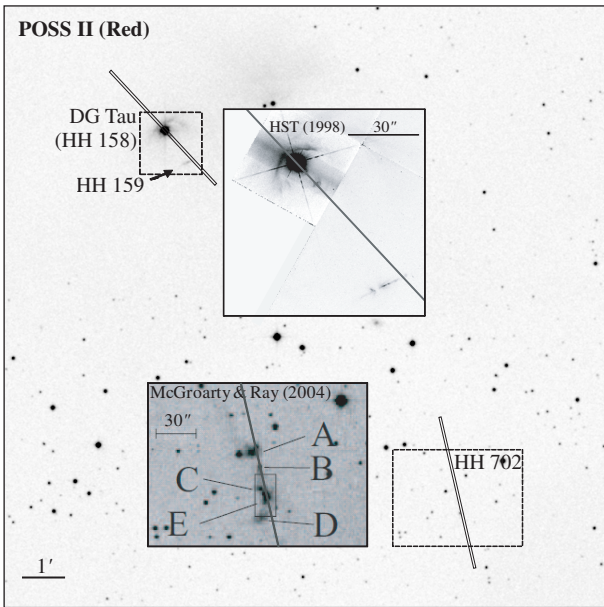


Figure 1. Slit positions on the object field. The center of the $2''.9(W) \times 3''.6(L)$ slit is set on HH 158 and HH 702. The position angles are 223° and 193° for HH 158 and HH 702, respectively. The background image is from the Second Palomar Observatory Sky Survey (POSS II). The HST image of HH 158, 159 and $H\alpha$ image of HH 702 from McGroarty & Ray (2004) are inserted as insets. Knot A – E of HH 702 are marked.

the jet with the positions during the last 30 years. In Section 2, we describe the observation and data reduction. In Section 3, we show the line intensity and ratios in the position-velocity map. Sections 4 and 5 are the discussion and summary, respectively.

2. OBSERVATION AND DATA REDUCTION

The observation was conducted on November 15, 2012 (UT) at BOAO (Bohyunsan Optical Astronomy Observatory) with an optical long-slit spectrograph installed on the Cassegrain focus of a 1.8-m telescope. The wavelength coverage was 5900 – 7050 \AA with 1200 gmm^{-1} grism, which was chosen to achieve the emission lines of $H\alpha$, [S II] $\lambda\lambda 6716, 6731$, [N II] $\lambda\lambda 6548, 6584$, [O I] $\lambda\lambda 6300, 6363$. The slit length and width were 3.6 arc-minute and 2.9 arc-second, respectively. The resultant spectral resolution was ~ 2070 (150 km s^{-1}) with the dispersion of $0.41 \text{ \AA pixel}^{-1}$. The CCD camera size is $4k \times 4k$ and the pixel scale along the slit is $0''.45 \text{ pixel}^{-1}$. Exposures on comparison lines were taken with the FeNeArHe lamp before and after the

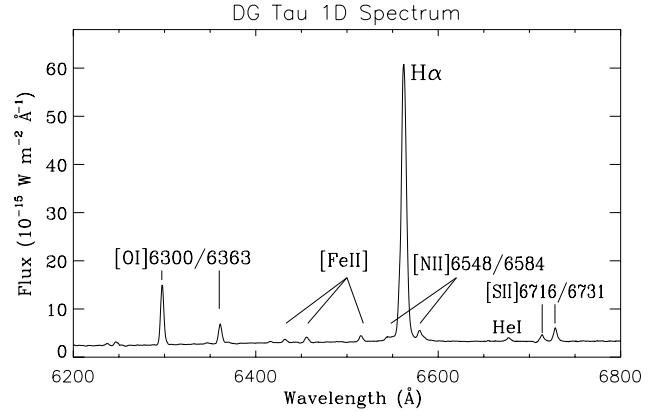


Figure 2. Integrated 1D spectrum of the central region of DG Tau.

exposure of each object, for wavelength calibration.

Table 1 shows the observation log of HH 158 and HH 702. Figure 1 shows the slit positions. The position angle (PA) of the HH 158 region including the central star is determined to be 223° (McGroarty & Ray 2004) after verification through an HST image. In the case of HH 702 located at $\sim 11'$ from the star, a PA of 193° was selected to obtain the emissions as much as possible. The exposure times were 1800 s and 2800 s for HH 158 and HH 702, respectively. HR 1544 (A1 type) was observed for telluric line correction and flux calibration.

We reduced the data with standard IRAF (Image Reduction and Analysis Facility) packages. Bias was subtracted from every frame. The flat-fielding was carried out using a normalized flat made by the APFLATTEN task. Bad pixels and cosmic-ray events were corrected. The spectra were extracted using the APALL task. The wavelength calibration and distortion correction were accomplished by the IDENTIFY, REIDENTIFY, FITCOORDS, TRANSFORM tasks. There were 1 and 2 pixel shifts between the comparison frames before and after the exposure for HH 158 and HH 702, respectively. Since the instrument is mounted on the Cassegrain focus, the instrumental displacement due to changes of the telescope elevation angle cause those shifts. We used the average of two lamp images to achieve a more accurate spectral calibration. There were also wavelength shifts between object frames. We corrected these shifts with the OH airglow lines in the object frames. After that, the sky emission lines were subtracted using the BACKGROUND task. The wavelength sensitivity correction and flux calibration were done with the standard star, HR 1544 (A1, $V=4.35$, $T_{\text{eff}}=9150 \text{ K}$).

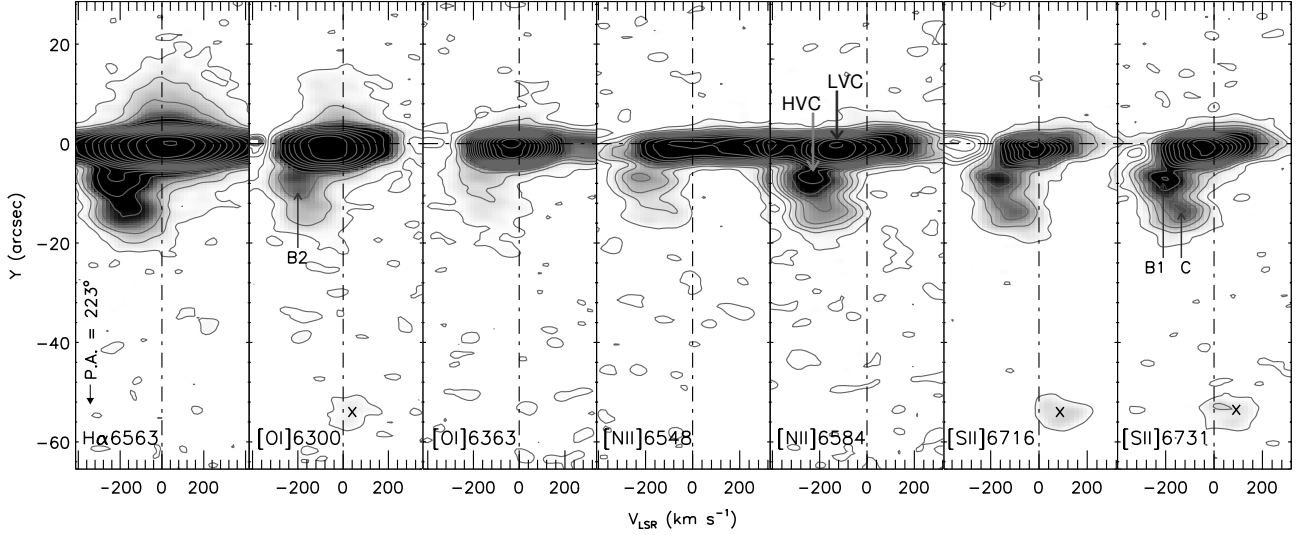


Figure 3. Continuum-subtracted PVDs of the emission lines of HH 158. Contours are drawn from 0.001 (3σ) to the peak levels of 5.2 ($\text{H}\alpha$), 0.98 ($[\text{O I}] \lambda 6300$), 0.33 ($[\text{O I}] \lambda 6363$), 0.17 ($[\text{N II}] \lambda 6548$ and $\lambda 6584$), 0.10 ($[\text{S II}] \lambda 6716$) and 0.22 ($[\text{S II}] \lambda 6731$) $\times 10^{-15} \text{ W m}^{-2} \text{ \AA}^{-1}$ with an equal interval in a logarithmic scale. The position angle (223°) of the slit is marked in the left panel. The original spectrum has Gaussian Smoothed with a σ value of 1.5×1.5 . Knots B1, B2 (Lavalley et al. 1997), and C (Eislöffel & Mundt 1998) are identified at $Y \sim -6''$ and $-14''$. The faint emissions at $Y = -54''$ in $[\text{O I}] \lambda 6300$, $[\text{S II}] \lambda 6716$ and $[\text{S II}] \lambda 6731$ are parts of the HH 159 outflow ejected from DG Tau B. These are labeled with the ‘x’ symbols in the figure.

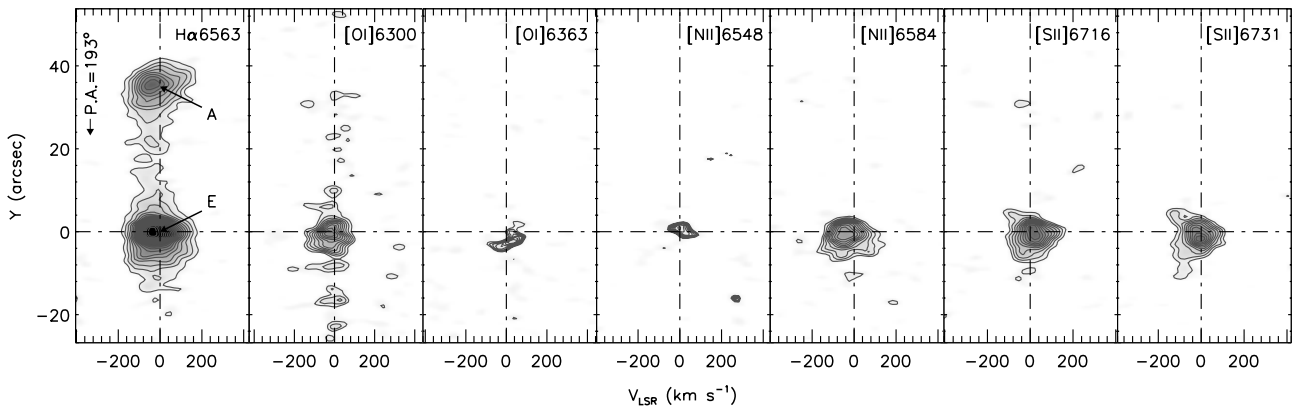


Figure 4. PVDs of the emission lines of HH 702. Contours are drawn from 0.05 to the peak levels of 1.5 ($\text{H}\alpha$), 0.45 ($[\text{O I}] \lambda 6300$), 0.12 ($[\text{O I}] \lambda 6363$), 0.11 ($[\text{N II}] \lambda 6548$), 0.38 ($[\text{N II}] \lambda 6584$), 0.60 ($[\text{S II}] \lambda 6716$), and 0.58 ($[\text{S II}] \lambda 6731$) in $10^{-17} \text{ W m}^{-2} \text{ \AA}^{-1}$ with an equal interval in a logarithmic scale. The position angle (193°) of the slit is marked in the left panel. Northeast is up and southwest is down in the diagram. The knots A and E of McGroarty & Ray (2004) are identified.

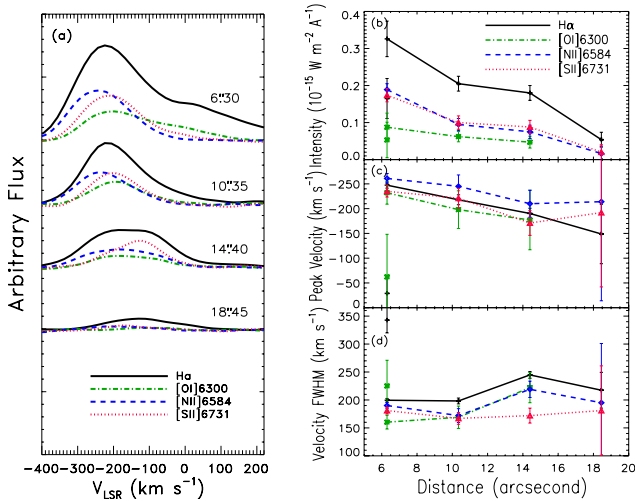


Figure 5. (a) Variation of the line profile, (b) peak intensity, (c) peak velocity, and (d) velocity width (FWHM) of the each emission line of HH 158 with the distance from DG Tau.

We removed a deep $H\alpha$ absorption line in the spectrum of the standard star using the SPLIT task before applying the wavelength sensitivity correction.

3. RESULTS

3.1. Kinematics

Figure 2 shows the 1D spectrum extracted from the central source region within $\pm 5''.4$ from DG Tau. In the wavelength range of $6200 - 6800\text{\AA}$, the $H\alpha$ $\lambda 6563$, $[O\ I]\ \lambda\lambda 6300, 6363$, $[N\ II]\ \lambda\lambda 6548, 6584$, and $[S\ II]\ \lambda\lambda 6716, 6731$ emissions were detected. The $[Fe\ II]$ emissions at $\lambda 6433, \lambda 6456, \lambda 6516\ \text{\AA}$ and the atomic Helium line at $\lambda 6678\ \text{\AA}$ are also identified. The $H\alpha$ emission has the highest intensity peak level of $6 \times 10^{-14}\ \text{W m}^{-2}\ \text{\AA}^{-1}$, which is ~ 30 times higher than that of the stellar continuum. The wing of the strong $H\alpha$ line affects the line profiles of two $[N\ II]$ lines.

Figures 3 and 4 show the position-velocity diagrams (PVDs) of HH 158 and HH 702. In Figure 3, the stellar continuum was subtracted. The position at $Y = 0''$ indicates DG Tau and the peak of knot E in Figures 3 and 4, respectively. HH 702 is located at $\sim 11'$ ($= 0.45\ \text{pc}$ at $140\ \text{pc}$ distance) southwest from DG Tau, as shown in Figure 1. The Gaussian smoothing was applied in both figures with sigma of 1.5×1.5 pixels.

In Figure 3, we identify the knots B1, B2 (Lavalley et al. 1997) and C (Eisloffel & Mundt 1998) at the distance of $6''.74, 8''.5$ and $\sim 14''$ from the star. The two velocity components appear in all emission lines in Figure 3. The low velocity component (LVC) is located within $\pm 5''$. The peak velocity of the LVCs varies from ~ -80 to $-20\ \text{km s}^{-1}$. The high velocity component (HVC) is relevant at $|Y| > 6''$. Figure 5c shows that the HVC is getting slower with increasing distance, e.g., $-270\ \text{km s}^{-1}$ at $6''.3$ (knot B1) and $-150\ \text{km s}^{-1}$ at $18''.45$. The real jet velocity is 1.18 times higher than the radial velocity considering the inclination angle of 32° with re-

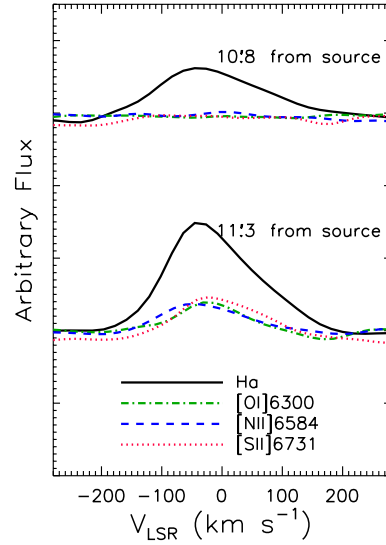


Figure 6. Line profiles at knots A (upper) and E (lower) of HH 702.

spect to the line of sight (Eisloffel & Mundt 1998), so the radial velocity of $-270\ \text{km s}^{-1}$ corresponds to the jet velocity of $-320\ \text{km s}^{-1}$. Additionally, it is noticeable that the emission lines are showing different spatial extents. In Figure 3, all emissions except $[N\ II]\ \lambda 6548$ show extensions to $Y \sim -20''$ at $3\ \sigma$ contour levels. Knot B1 is relevant in all emissions except $[O\ I]\ \lambda 6363$. Knot C is remarkable in $[S\ II]\ \lambda 6731$ and appears as plateaus in $H\alpha$, $[O\ I]\ \lambda 6300$, $[N\ II]\ \lambda 6584$, and $[S\ II]\ \lambda 6716$. Its intensity in $[N\ II]\ \lambda 6584$ is almost twice than that of $[O\ I]\ \lambda 6300$.

The faint emissions at a distance of $54''$ from DG Tau, detected at $[O\ I]\ \lambda 6300$, $[S\ II]\ \lambda 6716$ and $[S\ II]\ \lambda 6731$ may not come from HH 158 but from a part of HH 159 (see Figure 1), which is ejected from DG Tau B. These features are labeled with ‘x’ symbols.

In Figure 4, we identified the knots A and E of McGroarty & Ray (2004). The knot A of HH 702 is bright only in $H\alpha$ emission, while it is seen in the $[S\ II]$ images of McGroarty & Ray (2004) and Sun et al. (2003). Knot E is detected in all emission lines.

Figure 5 shows the variations of the line profile, peak intensity, peak velocity, and velocity width (FWHM) of each emission line with the distance from the star for HH 158. The data are sampled at $6''.30, 10''.35, 14''.40$, and $18''.45$ along the blueshifted jet. The plot starts from $6''.30$ where the knot B1 is located. The $[O\ I]\ \lambda 6300$ emission at $18''.45$ could not be measured due to faintness. In Figures 5b and 5c, the peak intensities and velocities of all the emission lines decrease with distance from the source. The uncertainty in the peak velocity measurement of $[N\ II]$ and $[S\ II]$ at $18''.45$ is $\pm 170\ \text{km s}^{-1}$. It is large, due to low a signal-to-noise ratio ($S/N < 2$). Even though there are small increases in velocity (Figure 5c), those are not significant due to a large uncertainty. The decreasing rate in the velocity of $H\alpha$ is $|\Delta v_{peak}| \sim 7\ \text{km}''$.

In Figure 5d, the range of velocity widths (FWHM)

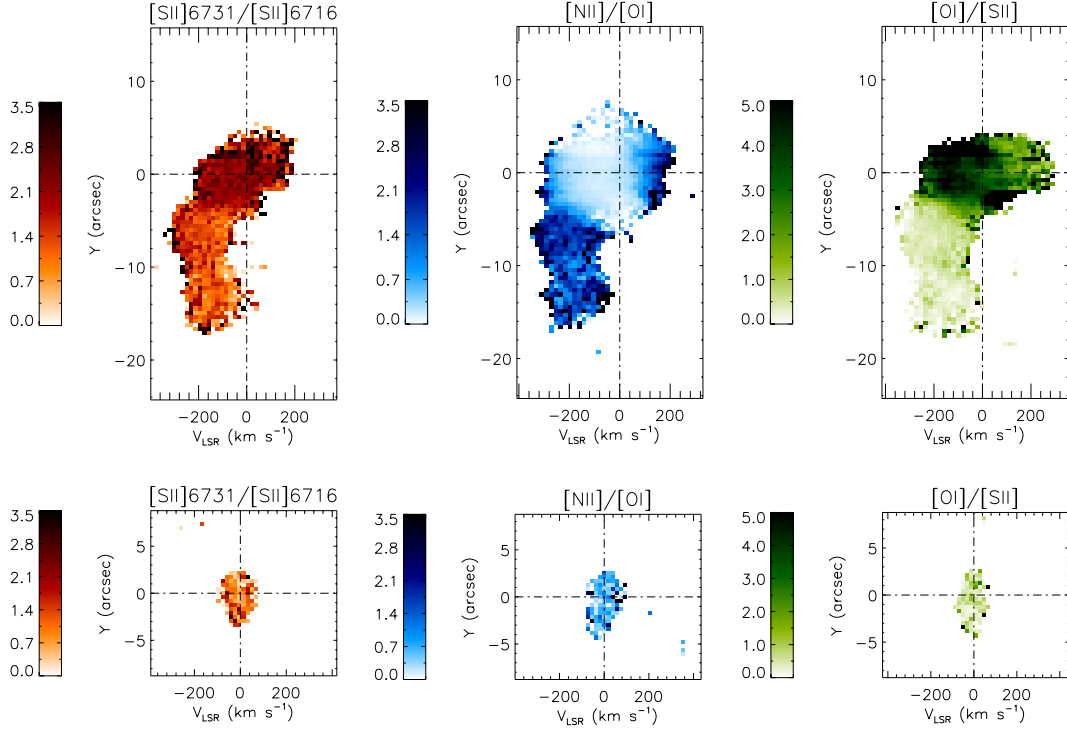


Figure 7. Ratios of forbidden emission lines: (*top*) HH 158 and (*bottom*) knot E of HH 702. (*Left*) [S II] λ 6731 / λ 6716, (*center*) [N II] λ 6548 + λ 6584 / [O I] λ 6300 + λ 6363, and (*right*) [O I] λ 6300 + λ 6363 / [S II] λ 6716 + λ 6731.

varies from 150 to 350 km s⁻¹, which is marginally resolved considering the velocity resolution of 150 km s⁻¹ in our observation. The low velocity component shows a wider width (\sim 340 km s⁻¹ in H α) than those estimated for high velocity components (\sim 160 – 230 km s⁻¹), compatible with results obtained in previous studies (e.g., Hirth et al. 1997; Pyo et al. 2003). Pyo et al. (2002) interpreted the wider velocity width of LVC as follows: if we assume that the outflow has the shape of a diverging wind, the velocity width represents the difference in radial velocities between the nearest and farthest streamlines from us. Thus, the wider velocity width means a wider opening angle.

Figure 6 shows the variations of the line profile along the distance of HH 702. The peak velocities of two knots are \sim -80 km s⁻¹ in the H α emission.

3.2. Line Ratios

The diagnostics by the ratios of forbidden emission lines is a very effective tool to study jet properties (Hartigan et al. 1994; Bacciotti et al. 1995). In Figure 7, we show the flux ratios of the forbidden emission lines of HH 158 and knot E of HH 702 in PVDs : [S II] λ 6731 / λ 6716, [N II] λ 6548 + λ 6584 / [O I] λ 6300 + λ 6363, and [O I] λ 6300 + λ 6363 / [S II] λ 6716 + λ 6731. The ratio of the [S II] doublet can be used to estimate the electron density (Osterbrock 1989), which is larger than 2 within $\pm 5''$ and decreases with the distance. The [N II]/[O I] ratio, which is proportional to the ionization fraction of hydrogen, is less than 0.5 at the center and increases at distances further than $\sim 5''$. The [O I]/[S II] ratio shows a similar trend with the

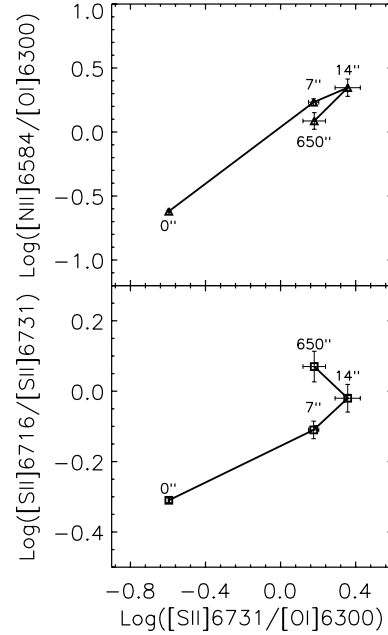


Figure 8. Line ratios estimated from HH 158 and HH 702, for four distances: $d = 0''$, $7''$, $14''$, $650''$.

[S II] doublet, large near the center and less than 1.0 further than $5''$. This ratio is related to the electron temperature. The line ratios of HH 702 are comparable with those obtained from HH 158 at $\sim 14''$ from the source. In Figure 8, we show the relations between three ratios, [S II] λ 6731 / [O I] λ 6300, [N II] λ 6584 / [O I] λ 6300, and [S II] λ 6716 / λ 6731. In the figure, the

Table 2
Locations of the knots with respect to the source position during the last 30 years.

Observation date	Distance (")			Reference
	B1 ^a	B2 ^a	C ^b	
1984 Oct. 9	–	–	8.5	Mundt et al. (1987) ^c
1986 Dec.	–	–	8.7	Eislöffel & Mundt (1998)
1992 Oct.	2.25	3.25	–	Solf & Böhm (1993)
1994 Nov. 3	2.7	3.9	–	Lavalley et al. (1997)
1997 Jan.	3.3	–	–	Dougados et al. (2000)
1998 Jan. 23-26	3.6	–	–	Lavalley-Fouquet et al. (2000)
1998 Feb. 5	–	–	10.65	HST ^d
1999 Jan. 14	3.8	–	–	Maurri et al. (2014)
2002 Sep. 17	4.6	–	12	Whelan et al. (2004)
2012 Nov. 15	6.74	8.75	13.7	This work
2014 Nov. 21	7.65	–	13.5	This work
Proper motion ("/yr ⁻¹)	0.230	0.272	0.180	–
Ejection date	1981	1980	1939	–

^a knot name from Lavalley et al. (1997).

^b knot name from Eislöffel & Mundt (1998).

^c position of knot C measured from Figure 9 of Mundt et al. (1987).

^d data taken from the HST archive, image of WFPC2 with the F675WF filter (proposal ID. 6855).

four points represent the values at distances $d = 0''$, $7''$, $14''$, $650''$ from the source. The distance of $650''$ is the location of the knot E of HH 702. The ratios increase with distance, as in a previous study (Lavalley-Fouquet et al. 2000, Figure 3a–b).

4. DISCUSSION

4.1. Velocity Variation

In Figure 3, the HVC is getting slower with distance. The velocity decreasing rate derived in Section 3 is ~ 7 km $''$. If the jet is continuously slowing, then it should stop at $35''$ from DG Tau. However, the knots of HH 702 located at $650''$ show a radial velocity of ~ -80 km s $^{-1}$.

Radial velocity (v_{rad}) measurements of knot C in HH 158 were reported in several previous studies. The knot was observed at $8''.5$ (Mundt et al. 1987) and $8''.7$ (Eislöffel & Mundt 1998) from the source in 1984 and 1987, respectively. We converted the heliocentric velocities to LSR velocities for the comparison with our data. In Mundt et al. (1987), v_{rad} were -136 km s $^{-1}$ in H α and -123 km s $^{-1}$ in [S II] $\lambda 6731$. In Eislöffel & Mundt (1998), [S II] is ~ -146 km s $^{-1}$. In 2002 (Whelan et al. 2004), [S II] $\lambda 6731$ and [O I] $\lambda 6300$ show v_{rad} of ~ -150 and -200 km s $^{-1}$, respectively. For the three emission lines mentioned above (H α , [S II], [O I]), our data show v_{rad} of -190 , -128 , and -180 km s $^{-1}$, respectively. Since different emission lines originate from different excitation conditions (i.e., temperature, density, etc.), they represent various regions in the outflow and show peaks in different velocities. The temporal variation of the velocity indicates that the velocity of knot is variable not only along the distance but also in time. This may imply that the ejected knot has been decelerated when it passes through the ambient medium. The velocity change of the knot could also occur due to the

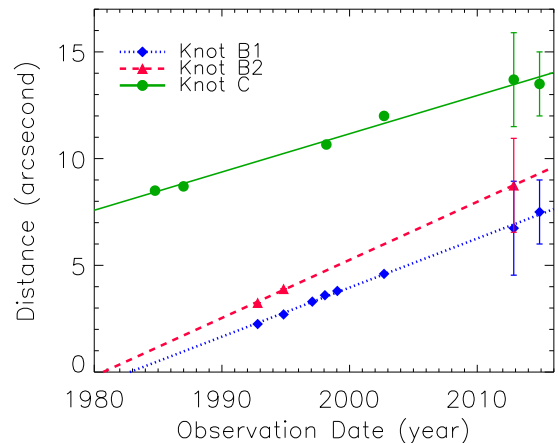


Figure 9. Positional variations of three knots. For knot B1, the data points are expanded from those in Pyo et al. (2003).

shock caused by the interaction between the faster gas stream behind and the slower front knot.

4.2. Proper Motion

In the emission lines of HH 158, we have distinguished three knots at $6''.74$ and $8''.75$, and $13''.7$ from the source in our data taken on November 15, 2012. We identified these knots as B1 and B2 of Lavalley et al. (1997), and knot C of Eislöffel & Mundt (1998). Recently, we could measure the knot positions from our additional data taken on November 21, 2014. For knot B1 and C, it was measured to $7''.65$ and $13''.5$, respectively. We note that knot B2 is not resolved in our spectra in 2014. The distance of knot C from the source is $0''.2$ smaller in 2014 than that in 2012. The seeing size is $> 2''$, so it may be due to the uncertainty of the measurement.

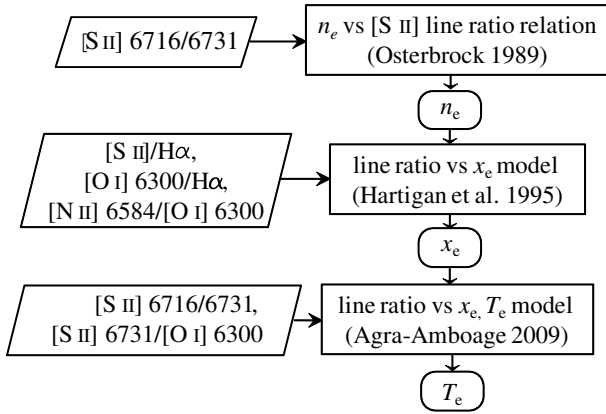


Figure 10. Flow chart showing how we derive the electron density (n_e), ionization fraction (x_e), and electron temperature (T_e).

In Table 2, we list the locations of the knots during the last ~ 30 years including our work. Basically, the positions of knot B1 are extended from Table 1 of Pyo et al. (2003). In the case of knot C, the first detection was reported in 1984 (Mundt et al. 1987) and it was seen at $12''$ in long-slit spectroscopy of Whelan et al. (2004) whose data were taken in 2002. In addition to these data points, we found an Hubble Space Telescope (HST) archival image obtained in 1998. The central peak in intensity contour of the knot in the HST image is located about $10''.65$ from DG Tau.

Figure 9 shows the temporal variation of the distance for the knots B1, B2, and C. We estimated the proper motion by linear regression: $0''.230 \pm 0''.006 \text{ yr}^{-1}$ for B1, $0''.272 \pm 0''.002 \text{ yr}^{-1}$ for B2, $0''.180 \pm 0''.008 \text{ yr}^{-1}$ for C (see Table 2). For knot B1, the estimated value is smaller than those obtained from previous studies, e.g., $0''.272$ in Pyo et al. (2003). Due to the large spatial resolution of our observation, the emission from knot B1 could be combined with that from knot A1 which has identified from previous studies (Bacciotti et al. 2000; Dougados et al. 2000). The smaller proper motion value of knot B1 may be estimated because knot A1 is located $2''.5$ closer (Maurri et al. 2014) to the star than knot B1. From the radial and the tangential velocities of each knots, the inclination angles (i) of B1, B2, and C are calculated as $32^\circ.5 \pm 0^\circ.7$, $37^\circ.0 \pm 0^\circ.2$, $32^\circ.3 \pm 1^\circ.1$ with respect to the line of sight.

For HH 702, McGroarty et al. (2007) estimated the tangential velocities of knots A and E as $129 \pm (10 - 28)$ and $186 \pm (10 - 28) \text{ km s}^{-1}$ in the $\text{H}\alpha$ emission, which correspond to $0''.19 \pm 0''.03$ and $0''.28 \pm 0''.03 \text{ yr}^{-1}$ in proper motion, respectively. From the radial velocity of $\sim -80 \text{ km s}^{-1}$ and the tangential velocities, the i of knot A and E are estimated as $\sim 58^\circ.2 \pm 1^\circ.1$ and $\sim 66^\circ.7 \pm 1^\circ.1$ with respect to the line of sight. The i of HH 702 are $\sim 21-35^\circ$ different from those of HH 158. We cannot avoid the possibility that the two HH objects were ejected from different sources although McGroarty et al. (2007) postulated that DG Tau seems the most probable driving source of HH 702 based on the proper motion direction. The different direction or the curved

shape of a jet may indicate the precession of the jet (Raga et al. 2001).

4.3. Line Ratios and Physical Parameters

Maurri et al. (2014) shows the line ratios and physical parameters within $5''$ of DG Tau. Their $[\text{S II}] \lambda 6731/\lambda 6716$ ratio is similar to our value. Knots B0 and B1 at $3''.3$ and $3''.8$ in Maurri et al. (2014) correspond to B1 at $6''.7$ in our data, after considering the proper motion. For this knot, the $[\text{S II}]$ doublet ratio in our data is slightly smaller than their estimation. For the $[\text{N II}]/[\text{O I}]$ ratio, the value in Maurri et al. (2014) is $0.01 - 2.5$, in a similar range as our data.

In Figure 10, the flow chart showing how we estimate the electron density (n_e), ionization fraction (x_e), and electron temperature (T_e) is displayed. Figure 11 shows the derived physical parameters for HH 158 and HH 702. The estimated properties of each knot show good agreement with the planar shock model (Hartigan et al. 1994; Lavalley-Fouquet et al. 2000). For the variation of the jet with the distance (i.e., variation with time), the model that treats the time variability of the jet (Raga et al. 1990; Raga & Kofman 1992) could explain our result with the internal shock caused by a variable (or periodic) jet ejection.

In our data, it is hard to distinguish the two velocity components within $\pm 5''$ from DG Tau due to the limitations in spatial and velocity resolution. The LVC is usually dominant close to the source, while HVC is extended further as shown in Figure 3. With this point of view, it is noticeable that the x_e of HVC at $|Y| > 4''$ is higher than that of the LVC at $|Y| < 4''$ but n_e is lower in HVC in Figure 11. On the other hand, previous studies (Lavalley-Fouquet et al. 2000; Maurri et al. 2014) show that the n_e and x_e are larger in their high and medium velocity interval than those in low velocity. Hamann (1994) also showed that the two velocity components of 20 TTS and 12 HAeBe stars are different in T_e and n_e . These results indicate that the two velocity components have different physical conditions. They are supposed to originate from different parts of star-disk systems or have different launching mechanisms (Kwan & Tademaru 1995; Pyo et al. 2003, 2006).

4.3.1. Electron density (n_e)

We calculated the electron density using the typical method of $[\text{S II}] \lambda 6716/\lambda 6731$ doublet ratio (Osterbrock 1989). This ratio is known to be insensitive to temperature variation, so we assumed the electron temperature as $\sim 10,000 \text{ K}$ to estimate the electron density (n_e). In the bottom panel of Figure 12, n_e is shown at four different distance positions. At the proximity of the star, n_e is $\sim 10^4 \text{ cm}^{-3}$ and decreases to $\sim 500 \text{ cm}^{-3}$ at $14''$ from the source. In HH 702 at $650''$, n_e is estimated to be $\sim 200 \text{ cm}^{-3}$.

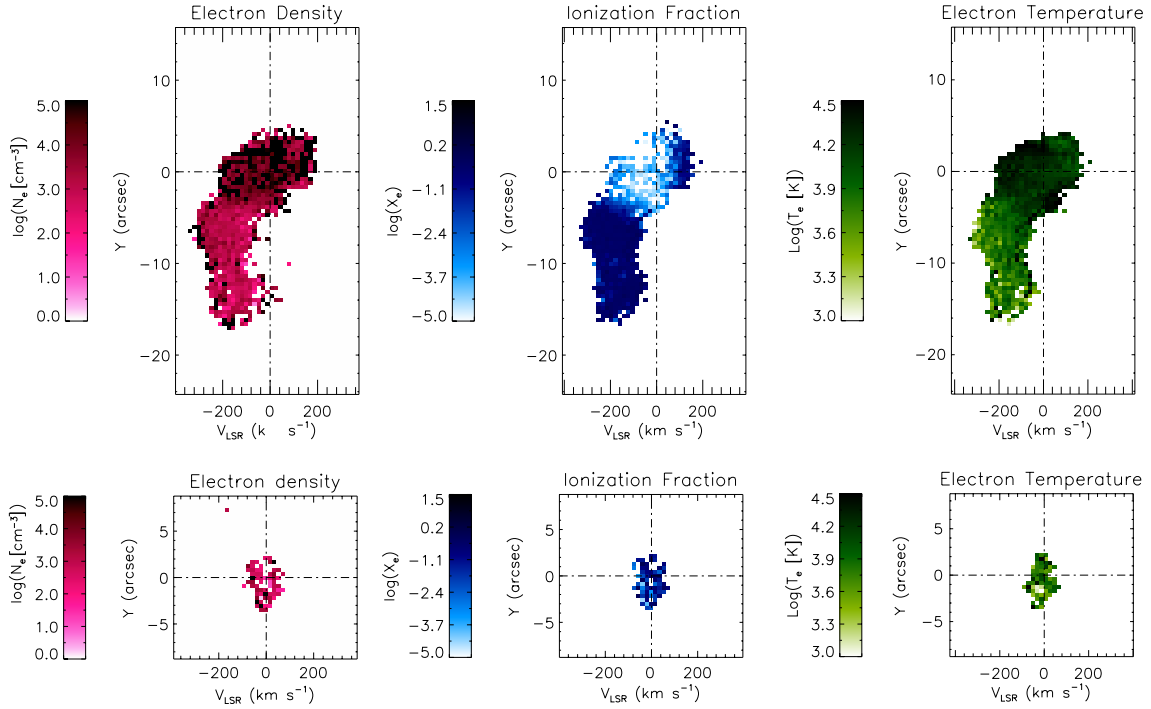


Figure 11. (Left) Electron density (n_e), (center) ionization fraction (x_e), and (right) electron temperature (T_e) of HH 158 and HH 702.

4.3.2. Ionization fraction (x_e)

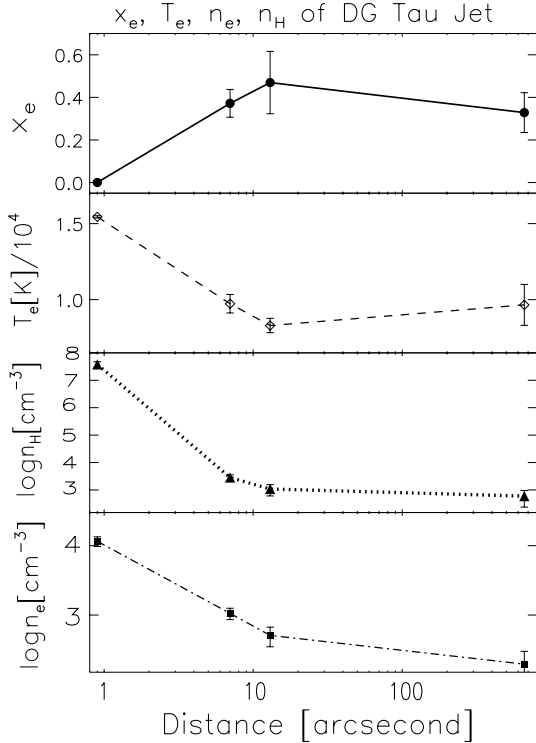


Figure 12. (From top to bottom) Ionization fraction (x_e), electron temperature (T_e), total hydrogen density (n_H), and electron density (n_e) of DG Tau jet as a function of the distance from the source.

We estimated the ionization fraction from the model grid of Hartigan et al. (1994), which shows the relations between the ionization fraction and the line ratio of $[\text{O I}] \lambda 6300/\text{H}\alpha$, $[\text{S II}] \lambda 6716+\lambda 6731/\text{H}\alpha$, $[\text{N II}] \lambda 6548/[\text{O I}] \lambda 6300$ in Figures 3, 4, and 6. We assumed the magnetic strength of $100 \mu\text{G}$, close to the average value in the model. The middle panel of Figure 11 and the top panel of 12 show the ionization fraction obtained from the $[\text{N II}]/[\text{O I}]$ ratio. Figure 13 shows the x_e values derived from three different ratios of $[\text{O I}] \lambda 6300/\text{H}\alpha$, $[\text{S II}] \lambda 6716+\lambda 6731/\text{H}\alpha$, $[\text{N II}] \lambda 6548/[\text{O I}] \lambda 6300$. The x_e values deduced by $[\text{O I}]/\text{H}\alpha$ and $[\text{S II}]/\text{H}\alpha$ ratios are in the range of $0.02 - 0.2$ which are much different from the values deduced from $[\text{N II}]/[\text{O I}]$. Since the $\text{H}\alpha$ emission flux comes not only from hydrogen recombination but also from collision (Bacciotti & Eisloffel 1999; Hartigan 2003), the use of the $\text{H}\alpha$ emission in the estimation of the ionization fraction may cause uncertainty. Thus, we used only the $[\text{N II}]/[\text{O I}]$ ratio in the calculation of x_e .

In Figures 11 and 12, x_e varies from $\sim 10^{-3}$ to ~ 0.5 with distance from the source. According to Figure 16 of Hartigan et al. (1994) which shows the relation between ionization fraction and shock velocity, the shock velocity surpasses 100 km s^{-1} with x_e of ~ 0.5 around knot C, which has $\sim 190 \text{ km s}^{-1}$ radial velocity. Lavalley-Fouquet et al. (2000) showed that x_e was increasing from 0.02 to 0.6 in HVC and from 0.004 to 0.1 in LVC within $4''$ from the source. Recently, Maurri et al. (2014) reported similar results within $5''$ using HST/STIS data obtained in 1999.

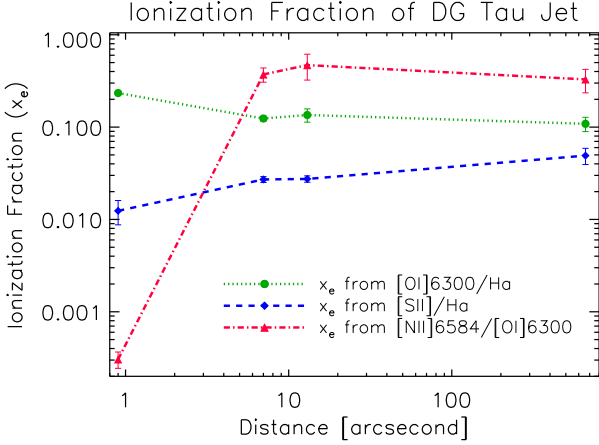


Figure 13. Ionization fraction (x_e) derived from the line ratios of [O I]/H α , [S II]/H α , and [N II]/[O I] based on the model of Hartigan et al. (1994).

For HH 702, x_e is greater than 0.2, which is similar to the value at 14'' from the source.

4.3.3. Electron temperature (T_e)

Electron temperature (T_e) values are shown in Figures 11 and 12, and are estimated from the model grid of Figure 3.6 in Agra-Amboage (2009) which shows the relation between emission line ratios and electron density, temperature, and ionization fraction. Using the x_e values obtained from 4.3.2, we have calculated the temperature from the model. T_e is higher than 15,000 K in the central region and decreases to $\sim 5,000$ K at 14''. At HH 702, the temperature was estimated to be slightly higher than that measured at knot C of HH 158.

4.4. Mass-loss Rate

The mass-loss rate of DG Tau was estimated in several previous studies. Hartigan et al. (1995) estimated the value of $3.2 \times 10^{-7} M_\odot \text{ yr}^{-1}$ from the [O I] $\lambda 6300$ emission line flux. In Lavalley-Fouquet et al. (2000), they have obtained the total value of $1.4 \times 10^{-8} M_\odot \text{ yr}^{-1}$ from the low, intermediate, and high velocity interval. The lower limit of mass-loss rate estimated by Beck et al. (2010) using the high-velocity Br γ emission was $1.2 \times 10^{-8} M_\odot \text{ yr}^{-1}$. Agra-Amboage et al. (2011) has estimated a total mass-loss rate from -300 km s^{-1} to -50 km s^{-1} components of the [Fe II] emission to $3.3 \times 10^{-8} M_\odot \text{ yr}^{-1}$. Maurri et al. (2014) also has estimated a mass-loss rate in the range of 10^{-8} to $10^{-9} M_\odot \text{ yr}^{-1}$ within the region 0''.7 from the source.

To estimate the mass-loss rate, we have adopted the method of Hartigan et al. (1995). We can calculate the mass-loss rate with the equation below:

$$\dot{M}_J = M_{TOT} V_\perp / l_\perp \quad (1)$$

V_\perp is the projected velocity on the plane of the sky and l_\perp is the projected jet size of the aperture. The total

mass of the jet M_{TOT} is estimated from the forbidden [O I] emission line with the following equation:

$$M_{TOT} = 9.61 \times 10^{-6} \left(1 + \frac{n_c}{n_e} \right) \left(\frac{L_{6300}}{L_\odot} \right) M_\odot \quad (2)$$

We estimate the mass-loss rate close to the star in the region within $\pm 5''.4$ (~ 750 AU) from the star. The equivalent width of [O I] $\lambda 6300$ emission line is measured to be 21 Å. We assume $R = 8.74$ magnitude, which is the dereddened R magnitude of DG Tau in Hartigan et al. (1995). The line luminosity of the [O I] $\lambda 6300$ (L_{6300}) converted into solar units is $9.2 \times 10^{-3} L_\odot$. V_\perp is estimated to be 75 km s^{-1} for the measured radial velocity 120 km s^{-1} , applying the inclination angle of 32° with respect to the line of sight. For l_\perp , we adopted $6.07 \times 10^{10} \text{ km}$ as the projected jet aperture size which corresponds to our slit-width of 2''.9 at a 140 pc distance to the Taurus molecular cloud. The FWHM of the jet measured from Dougados et al. (2000) is $\sim 0''.21 - 1''.25$ at a distance of $\sim 0''.4 - 4''$ from the star. Agra-Amboage et al. (2011) obtained the value of $\sim 0''.1 - 0''.8$ within $\sim 1''.4$ from the source, for the medium velocity component of the blueshifted jet. The jet width increases with distance in the vicinity of the source. We assume that the jet width in our estimation would be larger than 1''.25 because we estimate a mass-loss rate within $\pm 5''.4$. l_\perp of 2''.9 should be regarded as an upper limit. The temperature is assumed to be 8200 K, the average temperature for the [O I] $\lambda 6300$ line (Hartigan et al. 1994) and the critical density n_c in this temperature is $1.97 \times 10^6 \text{ cm}^{-3}$. For the electron density n_e , $7 \times 10^4 \text{ cm}^{-3}$ is used. The mass-loss rate, \dot{M}_J , for the [O I] emission line is $\sim 1.0 \times 10^{-7} M_\odot \text{ yr}^{-1}$.

The \dot{M}_J can be also estimated using the [S II] $\lambda 6731$ emission line. M_{TOT} is given by:

$$M_{TOT} = 1.43 \times 10^{-3} \left(\frac{L_{6731}}{L_\odot} \right) M_\odot \quad (3)$$

The equivalent width of the [S II] $\lambda 6731$ line is ~ 4.5 Å. The \dot{M}_J from [S II] is $1.1 \times 10^{-7} M_\odot \text{ yr}^{-1}$, which is almost the same as that of [O I] $\lambda 6300$. The mass-loss rates of [O I] $\lambda 6300$ and [S II] $\lambda 6731$ lines are a factor of 3 lower than the value of Hartigan et al. (1995) and is 3 – 7 times larger than those estimated by Lavalley-Fouquet et al. (2000) and Agra-Amboage et al. (2011). The calculated rate can differ with time because the measured line intensities and velocities of the jet could change due to the variability of the accretion/outflow rate of YSO. The rate becomes about two times larger if we replace the jet aperture size in our calculation with the value in Hartigan et al. (1995), which is 1''.5. The difference in spatial resolution between ours and higher-resolution (e.g., Agra-Amboage et al. 2011) observations cloud cause the difference in the mass-loss rate.

5. SUMMARY

We have investigated the kinematics and physical characteristics of HH 158 and HH 702 from the optical long slit spectroscopy. From the analysis through the PVDs and line profiles depending on the distance from the source, the peak velocity varies in the range of -50 to -250 km s $^{-1}$ and the line width was measured to 150 to 250 km s $^{-1}$. The proper motions were measured for knot B1, B2 and C of HH 158: $0''.230$, $0''.272$, and $0''.180$ yr $^{-1}$ for the knots, respectively.

From our analysis of the line ratios, we find that the physical properties (n_e , x_e , T_e) are changing along the jet from the proximity to the star to knot C at $14''$: (1) n_e varies from an order of 10^4 to 10^2 cm $^{-3}$, (2) x_e varies from ~ 0 to 0.4 at $14''$, (3) T_e varies from 15,000 K to $\sim 5,000$ K. It is noticeable that the line ratios and physical parameters of HH 702 do not show much difference from those obtained at knot C of HH 158. The mass-loss rate (\dot{M}_J) is estimated to be $\sim 10^{-7} M_\odot$ yr $^{-1}$, similar to values obtained in previous studies.

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