# Transmission Line Based Plucked String Model 전송선로 기반 탄현 모델 

Jingeol Lee ${ }^{\dagger}$ and Mark French*<br>(이진걸 ${ }^{\dagger}$, Mark French*)<br>${ }^{\dagger}$ Department of Electronic Engineering, Paichai University<br>*Department of Mechanical Engineering Technology, Purdue University<br>(Received February 18, 2013; accepted March 26, 2013)


#### Abstract

As one way to describe the behavior of a vibrating string, analogies to a transmission line have been made based on the fact that they have oppositely travelling waves on each of them. In such analogies, a rigid end to the string has been represented as an open circuit, and the displacement of the string as the current on the transmission line. However it turns out that the rigid end corresponds to a short circuit, the displacement to the voltage by the theory of the transmission line, and it is confirmed by experiments with circuit simulations. Based on these discoveries, a transmission line based plucked string model comprising a transmission line, two piecewise linear current sources, and switches is proposed. The proposed model is validated by showing that the voltage at the arbitrarily chosen location, and the voltage calculated over an infinitesimal portion at the end of the transmission line are consistent with the displacement at the corresponding location and the force on the rigid end of the string from the well known difference form of a wave equation governing the behavior of the string with its fundamental frequency tuned to that for the proposed model, respectively. Moreover, the applicability of the proposed model to modeling string and wind instruments is presented.


Keywords: Transmission line, Plucked string, Travelling waves, Difference form
PACS numbers: 43.75.Gh

초 록: 진동하는 현의 성질을 나타내는 방법으로 반대 방향으로 진행하는 파가 현과 전송선로에 존재한다는 사실에 기초하여 현은 전송선로에 비유되어왔다. 이러한 비유에서 현의 강역(rigid end)과 변위는 각각 전송선로의 개방회로 와 전류로 나타내어졌다. 그러나, 본 연구에서 강역과 변위는 각각 단락회로와 전압에 해당됨이 전송선로의 이론으로 부터 밝혀졌고 이를 회로시뮬레이션으로 확인하였다. 이러한 발견에 기초하여 전송선로, 구분적 선형 전류원, 스위치 들로 구성된 전송선로 기반 탄현 모델을 제안하였다. 임의로 선택된 지점에서의 전압과 전송선로 끝 극소 부분 양단에 서 계산된 전압이 현의 성질을 지배하는 파동방정식의 차분형식(difference form)으로 구한 해당 지점에서 변위와 강 역에서의 힘과 일치함을 보임으로서 제안한 모델이 정당함을 증명하였다. 또한, 제안된 모델의 현악기 및 관악기 모델 링의 적용성을 제시하였다.
핵심용어: 전송선로, 탄현, 진행파, 차분형식

## I. Introduction

The behavior of strings has been studied and simulated in various ways. Based on the fact that oppositely travelling waves exist on each of a string and an electrical transmission line, the vibrating string was considered as the transmission

[^0]line in continuous time domain. ${ }^{[1-3]}$ The string was modeled in discrete time domain. The representative is known as the structure of the Karplus-Strong algorithm consisting of a delay line followed by a filter in a loop, and the filter was later modified to accomplish the frequencydependent damping of harmonics. ${ }^{[4,5]}$ The other structure in discrete time domain was built in the form of digital waveguides in which two delay lines carry sampled
travelling waves. ${ }^{[6]}$ As a numerical approach, the wave equation governing a flexible string was approximated to a finite difference form. ${ }^{[7,8]}$

In the string models built with an analogy to the transmission line, a rigid end to strings has been represented as an open circuit, and the displacement of a string as the current on the transmission line. ${ }^{[1-3]}$ The usual analogy of a mechanical system to an electrical circuit, that is, the force to the voltage and the velocity to the current, has lead to the replacement of the rigid end with the open circuit. ${ }^{[3]}$ However it turns out that the rigid end corresponds to a short circuit, the displacement to the voltage by the theory of the transmission line, and it is confirmed by experiments with circuit simulations. Based on these discoveries, a transmission line based plucked string model comprising a transmission line, two piecewise linear current sources, and switches is proposed. The proposed model is validated by showing that the voltage at the arbitrarily chosen location, which is the center of the transmission line, and the voltage calculated over an infinitesimal portion at the end of the transmission line are consistent with the displacement at the center of the string and the force on the rigid end from the well known difference form of the wave equation governing the behavior of a string with its fundamental frequency tuned to that for the proposed model, respectively. Moreover, the applicability of the proposed model to modeling string and wind instruments is presented.
The paper is organized as follows; The difference form of the wave equation is reviewed in section II. In section III, the transmission line based plucked string model is proposed and is validated by comparing it with the difference form, and then conclusions are drawn in sectionIV.

## II. Review of the Difference Form

The wave equation for an ideal flexible string with no damping is given by

$$
\begin{equation*}
\frac{\partial^{2} y}{\partial t^{2}}=c^{2} \frac{\partial^{2} y}{\partial x^{2}} \tag{1}
\end{equation*}
$$

where $y(x, t)$ is the displacement of a string, and $x$ and $c$ are the distance and the velocity of a wave along the string, respectively. ${ }^{[9]}$ The velocity is given by $c=\sqrt{T / \mu}$ where $T$ is the tension, and $\mu$ is the mass per unit length of the string. The solution of the wave equation(1) is

$$
\begin{equation*}
y(x, t)=y_{1}(c t-x)+y_{2}(c t+x) \tag{2}
\end{equation*}
$$

where the function $y_{1}(c t-x)$ and $y_{2}(c t+x)$ represent a wave travelling in the $+x$ and $-x$ direction, respectively. By sampling time and space with a time step, $\Delta t$ and a spatial step, $\Delta x$, the discrete form of the wave equation is derived as $y(x, t) \rightarrow y(i \Delta x, n \Delta t) \rightarrow y(i, n) .{ }^{[7]}$ The finite difference form for the wave equation in (1) is

$$
\begin{align*}
& \frac{y(i, n+1)+y(i, n-1)-2 y(i, n)}{(\Delta t)^{2}} \\
& \approx c^{2}\left[\frac{y(i+1, n)+y(i-1, n)-2 y(i, n)}{(\Delta x)^{2}}\right] . \tag{3}
\end{align*}
$$

The displacement of the string at time step $n+1$ is derived from the equation in (3) as

$$
\begin{align*}
y(i, n+1) & =2\left(1-r^{2}\right) y(i, n)-y(i, n-1) \\
& +r^{2}[y(i+1, n)+y(i-1, n)], \tag{4}
\end{align*}
$$

where $r=c \Delta t / \Delta x$. The force at the end of the string, that is, $x=0$ is given by

$$
\begin{equation*}
F_{\text {bridge }}=T \frac{[y(1, n)-y(0, n)]}{\Delta x} \tag{5}
\end{equation*}
$$

The initial displacement for the case of plucking 5 mm at the one fifth of the distance from the end is shown in Fig. 1, and the time evolution of the displacement at the center of the string and the force at the end are shown in Fig. 2 and


Fig. 1. Initial displacement of the string.


Fig. 2. Displacement at the center.


Fig. 3. Force at the end.

Fig. 3, respectively with the length of the string of 650 mm , the spatial step $\Delta x$ of $0.65 \mathrm{~mm}, r=1$, the velocity of 429 $\mathrm{m} / \mathrm{s}$, and the tension of 149 N .

## III. Transmission Line Based Model

The voltage, $v(x, t)$ and the current, $i(x, t)$ on a lossless transmission are given by

$$
\begin{align*}
& -\frac{\partial v}{\partial x}=L \frac{\partial i}{\partial t}  \tag{6}\\
& -\frac{\partial i}{\partial x}=C \frac{\partial v}{\partial t}
\end{align*}
$$

where $L$ is the inductance, and $C$ is the capacitance per unit length of the transmission line. ${ }^{[10]}$ From equation(6), wave equations for the voltage and the current are derived as

$$
\begin{align*}
& \frac{\partial^{2} i}{\partial x^{2}}=L C \frac{\partial^{2} i}{\partial t^{2}}  \tag{7}\\
& \frac{\partial^{2} v}{\partial x^{2}}=L C \frac{\partial^{2} v}{\partial t^{2}}
\end{align*}
$$

The phasor representations for the voltage, $\nu(x, t)$ and the current, $i(x, t)$ are

$$
\begin{align*}
& v(x, t)=\mathfrak{R e}\left[\widetilde{V}(x) e^{j \omega t}\right] \\
& i(x, t)=\mathfrak{R} e\left[\widetilde{I}(x) e^{j \omega t}\right] . \tag{8}
\end{align*}
$$

The phasors, $\widetilde{V}(x)$ and $\widetilde{I}(x)$ on the transmission line are given by

$$
\begin{align*}
& \widetilde{V}(x)=V_{0}^{+} e^{-\gamma x}+V_{0}^{-} e^{\gamma x} \\
& \widetilde{I}(x)=I_{0}^{+} e^{-\gamma x}+I_{0}^{-} e^{\gamma x}=\frac{V_{0}^{+}}{Z_{0}} e^{-\gamma x}-\frac{V_{0}^{-}}{Z_{0}} e^{\imath x}, \tag{9}
\end{align*}
$$

where $V_{0}^{+}, I_{0}^{+}, V_{0}^{-}$, and $I_{0}^{-}$are the amplitude of the voltage and the current waves travelling to the $+x$, and to the $-x$ direction, respectively, and $\gamma$ is a complex propagation constant which is $j \omega \sqrt{L C}$ for the lossless
transmission line, and $Z_{0}$ is the characteristic impedance of the transmission line.

On the other hand, the reflection coefficient between an incident and a reflected wave at the end of the transmission line is given by

$$
\begin{equation*}
\Gamma=\frac{V_{0}^{-}}{V_{0}^{+}}=-\frac{I_{0}^{-}}{I_{0}^{+}}=\frac{Z_{L}-Z_{0}}{Z_{L}+Z_{0}}, \tag{10}
\end{equation*}
$$

where $Z_{L}$ is the load impedance. As represented in the equation(9), both voltage and current wave travel oppositely along the transmission line, and they are added to make a standing wave as in the string. The temporal solution of the wave equation given by equation(2) is consistent with the voltage and the current in phasor form given by equation(9). It was stated that only the current obeyed the wave equation, however so does the voltage as shown in the equation(7). ${ }^{[1-3]}$ Moreover, the displacement of the string had an analogy to the current, and the momentum of unit length of the string to the voltage, and a rigid end to an open circuit. In the plucked string, two initial waves, which are given by plucking, and take same forms spatially, travel oppositely. From the equation(9), it is evident that the voltage has a direct correspondence to the displacement if $V_{0}^{+}$and $V_{0}^{-}$take the same form spatially whereas the polarity of $V_{0}^{-}$has to be reversed in order to replace the current with the displacement, which is not the case in the string. With the voltage replaced with the displacement, the rigid end corresponds to $Z_{L}=0$, which gives rise to $\Gamma=-1$, and a free end to $Z_{L}=\infty$,
which produces $\Gamma=1$. Since the polarity of the reflected wave at the rigid end is reversed to the incident wave, the rigid end is equivalent to the short circuit. The open circuit which has been replaced with the rigid end is equivalent to not the rigid but the free end. By making the polarity of $V_{0}^{-}$to be reversed, and the both ends terminated with the open circuits, the current can be replaced with the displacement. However, the frequency response functions of the bodies of the string instruments generally carry a series of resonances characterized by the lower values of mechanical impedance, by which the frequency components in the vicinity of the those resonances sustain longer and are made to be dominant over others, and thus a specific timber of the instrument is accomplished. This reinforces the replacements of the displacement with the voltage and the rigid end with the short circuit.

An initial voltage distribution along the transmission line corresponding to the initial displacement of the string can be obtained by having two voltage waves travelled from the both ends as shown in Fig. 4. Two voltage waves last as long as the time delay, $t_{D}$ of the transmission line, and their shapes are spatially same, and are temporally symmetric to each other with respect to the half of their duration, $t_{D} / 2$. At $t_{D}$, two voltage wave coincide each other, and are summed together to be doubled to make the




Fig. 4. Two oppositely travelling waves from both ends of the transmission line.


Fig. 5. String model.
initial voltage distribution.
Based on the above mentioned ideas, the circuit for the string with the rigid ends is implemented with PSpice as shown in Fig. 5 consisting of the current sources, I1 and I2 having piecewise linear currents and three cascaded transmission lines, $\mathrm{T} 1, \mathrm{~T} 2$, and T 3 . The time delays of the transmission line $\mathrm{T} 1, \mathrm{~T} 2$, and T 3 are set to 0.7575 ms , 0.7565 ms , and 0.001 ms , respectively summing up to 1.515 ms , which produces the same fundamental frequency as in the difference form for comparison, and its characteristic impedance to $8 \Omega$ arbitrarily. The time delay of each of the cascaded transmission lines is set to monitor the voltage at the center of the transmission line and the voltage difference across T3 multiplied by 347,300 which correspond to the displacement and the force in the difference form, respectively, and thus the cascaded transmission lines can be replaced by a transmission line with the time delay of 1.515 ms and the same characteristic impedance of $8 \Omega$. The scale factor of 347,300 is obtained by 149 (tension)/ 0.00065 m (length of the string element in the difference form) $\times 1.515 / 1000$ (number of the string elements)/ 0.001. The initial voltage distribution is produced by the two piecewise linear currents as shown in Fig. 6. The current provided by current source I1 linearly increases from 0 mA at 0 ms to 0.625 mA at 0.303 ms which is one fifth of the total time delay of the transmission lines, and decreases back to 0 mA at 1.515 ms , while the current by I 2 is symmetric to that by I 1 with the respect to


Fig. 6. Currents from I1 and I2.
the half of the total time delay. The both current sources provide the voltage waves travelling to the opposite direction. At 1.515 ms , the voltage waves, which are obtained by the current flown into the transmission line multiplied by the input impedance of the transmission line from the end, coincide along the transmission line. At the same time, the switches, u3 and u4 open, and u1 and u2 close, by which the current sources and resistor, R2 and R4 are disconnected from the circuit, and the ends of transmission line, T 1 and T 3 are shorted, by which the voltage waves are ready to be reflected back and forth between the both ends.

In the transmission line, the input impedance is given by

$$
\begin{equation*}
Z_{i n}=Z_{0} \frac{Z_{L}+j Z_{0} \tan \beta l}{Z_{0}+j Z_{L} \tan \beta l} \tag{11}
\end{equation*}
$$

where $\beta=2 \pi / \lambda$, and $l$ is the length of the transmission line. Since the term of $\beta l$ is the multiples of $\pi$ for the fundamental frequency and its harmonics, the input impedance simply ends up with the load impedance of $8 \Omega$ assuming that the ideal current source having the source impedance of infinity. The current from the current source I1 and I2 divided equally flows through the transmission line and resistor, R2 and R4, which produces the peak value of 2.5 mV . When two voltage waves coincide at 1.515 ms , the voltage level doubles to 5 mV which corresponds to the initial displacement of the string in the difference form given by plucking 5 mm at the one fifth of the distance from the end. The voltage waveforms at the junction between the transmission line T 1 and T 2 and at the output of the gain in Fig. 5 are shown in Fig. 7 and Fig. 8, which are consistent with those from the difference form as presented in Fig. 2 and 3. The waveforms in Fig. 7 and Fig. 8 are recorded from 1.515 ms because it takes 1.515 ms to build the voltage distribution along the transmission line which corresponds to the initial displacement along the string given by plucking it.

In order to show the extensibility of the proposed string


Fig. 7. Voltage at the center of the transmission line.


Fig. 8. Voltage at the output of the gain.


Fig. 9. Voltage across the resistor of $2 \Omega$.
model to modeling of string and wind instruments, the short circuit to the right of T 3 in Fig. 5 is replaced with a resistor of $2 \Omega$, and the voltage across the resistor, which corresponds to the displacement at the resistive load in the electrical equivalence, is presented in Fig. 9. In this case, the voltage reflection coefficient is given by

$$
\begin{equation*}
\Gamma=\frac{Z_{L}-Z_{0}}{Z_{L}+Z_{0}}=\frac{2-8}{2+8}=-0.6 \tag{12}
\end{equation*}
$$

In Fig. 9, the voltage waveform between 1.515 ms and 3.03 ms is resulted from the sum of an incident voltage wave to the resistor of $2 \Omega$ and its reflected wave, which produces the peak value of $2.5+2.5 \times(-0.6)=1 \mathrm{mV}$, the waveform between 3.03 ms and 4.545 ms from the sum of a reflected from the short circuit to the left of T 1 and travelled wave to the resistor and its reflected wave with the peak value of -1 mV , the waveform between 4.545 ms and 6.06 ms from the sum of a reflected at the resistor between 1.515 ms and 3.03 ms , and reflected from the short circuit, and travelled back to the resistor, and its reflected wave with the peak value of $2.5 \times(-0.6)$ $\times(-1)(1-0.6)=0.6 \mathrm{mV}$, the waveform between 6.06 ms and 7.575 ms from the sum of a reflected between 3.03 ms and 4.545 ms , and reflected from the short circuit, and travelled back to the resistor, and its reflected wave with the peak value of -0.6 mV , and so forth. From this experiment, it is expected that the proposed string model can be applied to modeling plucked string instruments such as guitar by connecting a circuit for a top plate coupled with air in a sound box in place of the resistor of $2 \Omega$, and to struck and bowed string instruments by modifying the way the voltage wave is applied and connecting circuits for their bodies. It is well known that standing waves in an enclosed air column are analogous to those on the string, and thus the proposed model is also expected to be applicable to even wind instruments. ${ }^{[11]}$

## IV. Conclusions

In existing transmission line based string models, the displacement of a string had an analogy to the current on a transmission line, and a rigid end to an open circuit. In this paper, it is shown with the theory of the transmission line that the displacement corresponds to the voltage on the transmission line, and the rigid end to a short circuit. By applying these correspondences and the presented method of the excitation of the transmission line, a novel transmission line based plucked string model is proposed. The proposed model is validated with circuit simulations by demonstrating that the displacement of the string and the force at the end from the difference form are consistent with those from the proposed model. Moreover, the applicability of the proposed string model to modeling string and wind instruments is presented. Researches on modeling string and wind instruments using the proposed string model will succeed.

## References

1. Winston E. Kock, "The vibrating string considered as an electrical transmission line," J. Acoust Soc Am. 8, 227-233 (1937).
2. John C. Schelleng, "The violin as a circuit," J. Acoust Soc Am. 35, 326-338 (1963).
3. R. J. Clarke, "The analysis of mutiple resonance in a vibrating mechanical system by the use of the electrical transmission line analogy," ACUSTICA 40, 34-39 (1978).
4. Kevin Karplus and Alex Strong, "Digital synthesis of plucked-string and drum timbres," Comput Music J. 7, 43-55 (1983)
5. Sangjin Cho, "Development of loop filter design of plucked string instruments" (in Korean), J. Acoust. Soc. kr. 30, 107-113 (2011).
6. Julius O. Smith III, "Physical modeling using digital waveguides," Comput Music J. 16, 74-91 (1992).
7. Nicholas J. Giordano and Hisao Nakanishi, Computational Physics, Second Edition (Prentice Hall, New Jersey, 2005).
8. Antoine Chaigne and Anders Askenfelt, 'Numerical simulations of piano strings. I. A physical model for a struck string using finite difference methods," J. Acoust Soc Am. 95, 1112-1118 (1994).
9. Lawrence E. Kinsler, Austin R. Frey, Alan B. Coppens, and James V. Sanders, Fundamentals of Acoustics, Fourth

Edition (John Wiley \& Sons, New York, 2000).
10. Fawwaz T. Ulaby, Fundamentals of Applied Electromagnetics, Second Edtion (Prentice Hall,New Jersey, 2001).
11. Donald E. Hall, Musical Acoustics, 3rd. Ed. (Brooks/Cole, Pacific Grove, 2002).

## 저자 약력

- Jingeol Lee(이 진 걸)


Jingeol Lee received the Bachelor of Engineering degree; the Master of Science degree from Korea University in 1981 and 1985, respectively; and the Doctor of Philosophy degree from University of Forida in 1994. From 1982 to 1990, he was with Agency for Defense Development, and from 1995 to 1996, he had worked for Samsung Eectronics, where he was involved in research and development of military electronics throughout his job experiences. Since 1997 he has been with the Department of Eectronic Engineering, Paichai University. His professional interest includes musical acoustics.

- Mark French


Mark French received a BS in Aerospace and Ocean Engineering from Virginia Tech in 1985. He began work as a civilian engineer for the US Air Force while also pursuing further education. He received an MS in Aerospace Engineering in 1988 and a PhD in Aerospace Engineering in 1993, both from the University of Dayton. He moved to the auto industry in 1995, where he worked as a noise and vibration engineer and a lab manager. He moved to Purdue University in 2004, where is now an associate professor in the Department of Mechanical Engineering Technology. He works in experimental mecharics and in musical instrument design.


[^0]:    $\dagger$ Corresponding author: Jingeol Lee (jingeol@pcu.ac.kr)
    Department of Electronic Engineering, Paichai University, 439-6 Domadong Seogu Daejeon 302-735, Republic of Korea
    (Tel: 82-42-520-5707; Fax: 82-70-4362-6311)

