

## MOVING FRAMES ON GENERALIZED FINSLER STRUCTURES

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*In memory of Professor Cornel Urs*

ABSTRACT. We study the relation between an  $R$ -Cartan structure  $\alpha$  and an  $(I, J, K)$ -generalized Finsler structure  $\omega$  on a 3-manifold  $\Sigma$  showing the difficulty in finding a general transformation that maps  $\alpha$  to  $\omega$ . In some particular cases, the mapping can be uniquely determined by geometrical conditions. Moreover, we are led in this way to a negative answer to our conjecture in [12].

### 1. Introduction

Let us start by recalling that a Finsler norm on a real smooth,  $n$ -dimensional manifold  $M$  is a function  $F : TM \rightarrow [0, \infty)$  that is positive and smooth on  $\widetilde{TM} = TM \setminus \{0\}$ , has the *homogeneity property*  $F(x, \lambda v) = \lambda F(x, v)$ , for all  $\lambda > 0$  and all  $v \in T_x M$ , having also the *strong convexity* property that the Hessian matrix

$$g_{ij} = \frac{1}{2} \frac{\partial^2 F^2}{\partial y^i \partial y^j}$$

is positive definite at any point  $u = (x^i, y^i) \in \widetilde{TM}$ .

The fundamental function  $F$  of a Finsler structure  $(M, F)$  determines and it is determined by the (tangent) *indicatrix*, or the total space of the unit tangent bundle of  $F$

$$\Sigma_F := \{u \in TM : F(u) = 1\},$$

which is a smooth hypersurface of  $TM$  such that at each  $x \in M$  the *indicatrix at  $x$*

$$\Sigma_x := \{v \in T_x M \mid F(x, v) = 1\} = \Sigma_F \cap T_x M$$

is a smooth, closed, strictly convex hypersurface in  $T_x M$ .

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A Finsler structure  $(M, F)$  can be therefore regarded as a smooth hypersurface  $\Sigma \subset TM$  for which the canonical projection  $\pi : \Sigma \rightarrow M$  is a surjective submersion and having the property that for each  $x \in M$ , the  $\pi$ -fiber  $\Sigma_x = \pi^{-1}(x)$  is strictly convex including the origin  $O_x \in T_xM$ .

A generalization of this notion is the *generalized Finsler structure* introduced by R. Bryant (see [4], [5]). We will consider in the following only the low dimensional case where an  $(I, J, K)$ -*generalized Finsler structure* is a coframing  $\omega = (\omega^1, \omega^2, \omega^3)$  on a three dimensional manifold  $\Sigma$  that satisfies the structure equations

$$(1) \quad \begin{aligned} d\omega^1 &= -I\omega^1 \wedge \omega^3 + \omega^2 \wedge \omega^3, \\ d\omega^2 &= -\omega^1 \wedge \omega^3, \\ d\omega^3 &= K\omega^1 \wedge \omega^2 - J\omega^1 \wedge \omega^3. \end{aligned}$$

Here  $I, J, K$  are smooth functions on  $\Sigma$ , called the invariants (in the sense of equivalence problem) of the generalized Finsler structure  $(\Sigma, \omega)$ . In order to put in evidence the invariants, we will use the naming  $(I, J, K)$ -generalized Finsler structure.

The exterior differentials of these equations give the Bianchi type equations

$$(2) \quad J = I_{\omega^2}, \quad K_{\omega^3} + KI + J_{\omega^2} = 0,$$

where the subscripts denote the directional derivatives with respect to the coframing  $\omega$ , i.e.,  $df = f_{\omega^1}\omega^1 + f_{\omega^2}\omega^2 + f_{\omega^3}\omega^3$ , for any smooth function  $f$  on  $\Sigma$ .

By extension, one can study the generalized Finsler structure  $(\Sigma, \omega)$  defined in this way ignoring even the existence of an underlying surface  $M$ .

Let us recall that, for an  $(I, J, K)$ -generalized Finsler structure  $(\Sigma, \omega^1, \omega^2, \omega^3)$ , we can always consider

1.  $\mathcal{P} : \{\omega^1 = 0, \omega^3 = 0\}$  the “*geodesic*” foliation of  $\Sigma$ , i.e., the leaves are curves on  $\Sigma$  tangent to  $\hat{e}_2$ ;
2.  $\mathcal{Q} : \{\omega^1 = 0, \omega^2 = 0\}$  the “*indicatrix*” foliation of  $\Sigma$ , i.e., the leaves are curves on  $\Sigma$  tangent to  $\hat{e}_3$ ,

where  $\{\hat{e}_1, \hat{e}_2, \hat{e}_3\}$  is the dual frame of  $\{\omega^1, \omega^2, \omega^3\}$ .

Any of the codimension two foliations above is called *amenable* if the leaf space  $\Lambda$  of the foliation is a smooth surface such that the natural projection  $\pi : \Sigma \rightarrow \Lambda$  is a smooth submersion.

By extension, an  $(I, J, K)$ -generalized Finsler structure  $(\Sigma, \omega)$  whose “*indicatrix*” foliation  $\mathcal{Q}$ , “*geodesic*” foliation  $\mathcal{P}$ , and “*normal*” foliation  $\mathcal{R}$  is amenable, is called *amenable*, *geodesically amenable* and *normally amenable*, respectively. Remark that in the case of a Riemannian surface, the foliations  $\mathcal{P}$  and  $\mathcal{R}$  coincide up to the orientation of the Riemannian volume form, so the notions of *geodesically amenable* and *normally amenable* also coincide. Of course this is not true anymore in the case Finslerian case ([12] for details).

The difference between a classical Finsler structure and a generalized one is global in nature, in the sense that every generalized Finsler surface structure is locally diffeomorphic to a classical Finsler surface structure (see [4]).

The following fundamental result gives the necessary and sufficient conditions for a generalized Finsler structure to become a classical one.

**Theorem 1.1** ([4]). *The necessary and sufficient conditions for an  $(I, J, K)$ -generalized Finsler structure  $(\Sigma, \omega)$  to be realizable as a classical Finsler structure on a smooth surface  $M$  are*

1. *the leaves of the indicatrix foliation  $\{\omega^1 = 0, \omega^2 = 0\}$  are compact;*
2. *it is amenable;*
3. *the canonical immersion  $\iota : \Sigma \rightarrow TM$ , given by  $\iota(u) = \pi_{*,u}(\hat{e}_2)$ , is one-to-one on each  $\pi$ -fiber  $\Sigma_x$ , where  $M$  is the leaf space of the indicatrix foliation.*

In order to study the differential geometry of a classical Finsler structure  $(M, F)$ , it is customary to construct the pull-back bundle  $(\pi^*TM, \pi, \Sigma)$  with the  $\pi$ -fibers  $\pi^{-1}(u)$  diffeomorphic to  $T_xM$ , where  $u = (x, v) \in \Sigma$  (see [1]). In general this is not a principal bundle.

By defining an orthonormal moving coframing on  $\pi^*TM$  with respect to the Riemannian metric on  $\Sigma$  induced by the Finslerian metric  $F$ , the moving equations on this frame lead to the so-called Chern connection. This is an almost metric compatible, torsion free connection of the vector bundle  $(\pi^*TM, \pi, \Sigma)$ .

Let us remark that the condition  $I = 0$  imposes  $J = 0$  from Bianchi identities and therefore the simplest generalized Finsler structure is a  $(0, 0, K)$ -generalized Finsler structure with structure equations

$$(3) \quad \begin{aligned} d\omega^1 &= \omega^2 \wedge \omega^3, \\ d\omega^2 &= -\omega^1 \wedge \omega^3, \\ d\omega^3 &= K\omega^1 \wedge \omega^2. \end{aligned}$$

This kind of structure is called a  $K$ -Cartan structure on the 3-manifold  $\Sigma$  (see [7]).

We can conclude that the necessary and sufficient condition for an  $(I, J, K)$ -generalized Finsler structure to be a  $K$ -Cartan structure is  $I = 0$ .

A  $K$ -Cartan structure that satisfies the conditions in Theorem 1.1 gives a Riemannian structure on the surface  $M$ .

A generalization of the geodesic foliation of a  $K$ -Cartan structure is also known (see [4], [5]). Let us suppose that on the 3-manifold  $\Sigma$  endowed with an amenable  $K$ -Cartan structure  $\alpha$ , we have a basic one-form  $\beta$  for the fibration  $\Sigma \rightarrow \Lambda$ , where  $\Lambda$  is the leaf space of the foliation  $\{\alpha^1 = 0, \alpha^2 = 0\}$ . Then,  $\mathcal{P}_\beta : \{\omega^1 = 0, \omega^3 = \beta\}$  is called the " $\beta$ -geodesic" foliation of  $\Sigma$ . Obviously, in the case  $\beta = 0$  we get the usual geodesic foliation. The definition can be easily extended to the non-amenable case as well.

The geometrical meaning of  $\beta$  geodesics can be thought as follows. Let us assume that the smooth surface  $\Lambda$  is endowed with a Riemannian structure whose orthonormal frame bundle is  $\Sigma$ , and whose tautological one-forms  $\alpha^1, \alpha^2$  together with the Levi-Civita connection form  $\alpha^3$  give a coframe on  $\Sigma$ . Then, a smooth curve  $\gamma : [a, b] \rightarrow \Lambda$  is called a  $\beta$ -geodesic if and only if  $k_g = \beta(\dot{\gamma})$ , where  $k_g$  is the geodesic curvature of  $\gamma$  and  $\dot{\gamma}$  is the tangent vector along  $\gamma$ .

Another particular class of special Finsler structures is the one of positive constant flag curvature  $K = 1$ . This leads to the notion of generalized Finsler structure with  $K = 1$ , or, an  $(I, J, 1)$ -generalized Finsler structure on the 3-manifold  $\Sigma$ . The structure equations of a such structure are obtained immediately by simply substituting the curvature condition  $K = 1$  in (1) and (2), respectively. Many examples of classical Finsler metrics of constant flag curvature are nowadays known, see for example [2] for a complete classification of constant flag curvature Randers structures, or [13] for other examples.

However, the existence of generalized and classical Finsler structures of constant positive flag curvature  $K = 1$  on surfaces was proved for the first time by R. Bryant in [4] using a path geometry approach.

We will point out another interesting class of Finsler structures. In order to do this, let us first recall that, in the case of a classical Finsler structure  $(M, F)$ , the canonical parallel transport  $\Phi_t : T_x M \setminus 0 \rightarrow T_{\sigma(t)} M \setminus 0$ , defined by the Chern connection along a curve  $\sigma$  on  $M$ , is a diffeomorphism that preserves the Finslerian length of vectors. Unlike the parallel transport on a Riemannian manifold,  $\Phi_t$  is not a linear isometry in general.

This unexpected fact leads to some classes of special Finsler metrics. A Finsler metric whose parallel transport is a linear isometry is called a *Berwald structure*, and one whose parallel transport is only a Riemannian isometry is called a *Landsberg structure*.

Equivalently, a Berwald metric is a Finsler metric whose Chern connection coincides with the Levi Civita connection of a certain Riemannian metric on  $M$ , in other words it is "Riemannian-metrizable". These are the closest Finslerian metric to the Riemannian ones. The connection is Riemannian, while the metric is not. However, in the two dimensional case, any Berwald structure is Riemannian or flat locally Minkowski, i.e., there are no geometrically interesting Berwald surfaces.

Landsberg structures have the property that the Riemannian volume of the Finslerian unit ball is a constant. This remarkable property leads to a proof of Gauss-Bonnet theorem on surfaces [1] and other interesting results.

Obviously, any Berwald structure is a Landsberg one. However, there are no examples of global Landsberg structures that are not Berwald. The existence of Landsberg structures that are not Berwald is one of the main open problems in modern Finsler geometry.

These notions can be easily generalized as follows. A *generalized Landsberg structure* on  $\Sigma$ , or an  $(I, 0, K)$ -generalized Finsler structure is a generalized Finsler structure  $(\Sigma, \omega)$  such that  $J = 0$ , or equivalently,  $I_{\omega 2} = 0$ . Such a

generalized structure is characterized by the structure equations

$$(4) \quad \begin{aligned} d\omega^1 &= -I\omega^1 \wedge \omega^3 + \omega^2 \wedge \omega^3, \\ d\omega^2 &= -\omega^1 \wedge \omega^3, \\ d\omega^3 &= K\omega^1 \wedge \omega^2, \end{aligned}$$

and Bianchi identities

$$(5) \quad \begin{aligned} dI &= I_{\omega^1}\omega^1 + I_{\omega^3}\omega^3, \\ dK &= K_{\omega^1}\omega^1 + K_{\omega^2}\omega^2 - KI\omega^3, \end{aligned}$$

where  $I$  and  $K$  are smooth functions defined on  $\Sigma$ .

A *generalized Berwald structure* is a generalized Finsler structure characterized by the structure equations (4), and

$$dI \equiv 0 \pmod{\omega^3},$$

or, equivalently,

$$I_{\omega^1} = I_{\omega^2} = 0, \quad I_{\omega^3} \neq 0.$$

By means of Cartan-Kähler theory we have shown in [12] that there exists non-trivial generalized Landsberg structures on a 3-manifold  $\Sigma$ . Moreover, using a path geometry approach we have constructed locally a generalized Landsberg structure by means of a Riemannian structure on a surface.

The problems we consider in the present paper are:

1. *Clarify the relation between Cartan structures and generalized Finsler structures;*
2. *For given Cartan and generalized Finsler structures  $\alpha$  and  $\omega$ , respectively, how many types of regular matrices  $A$  exist such that  $\omega = A\alpha$ ?*
3. *Give a final answer to the conjecture in [12] on the existence of genuine Landsberg structures on surfaces via a certain construction.*

The essence of the method is to start with an  $R$ -Cartan structure  $\alpha$  on the 3-manifold  $\Sigma$  and to construct by a coframe changing an  $(I, J, K)$ -generalized Finsler structure  $\{\omega^1, \omega^2, \omega^3\}$  by means of a matrix  $A = (a_{ij})$ , namely  $\omega = A\alpha$ . By taking the exterior derivative of this formula and imposing conditions such that the structure equations of the  $R$ -Cartan structure and  $(I, J, K)$ -generalized Finsler structure are satisfied, respectively, we obtain a set of differential conditions for the functions  $a_{ij}$  in terms of its directional derivatives with respect to the coframe  $\omega$ . The functions satisfying these conditions can be regarded as the integral manifolds of a linear Pfaffian system.

Using now Cartan-Kähler theory we can study the existence of the integral manifolds of this linear Pfaffian system. In the most general case, of course this problem is quite difficult, but in the some particular cases, for example an  $(I, J, 1)$ -generalized Finsler structure, or, an  $(I, 0, K)$ -generalized Finsler structure (see [4] and [12]) one can construct explicitly the matrix  $A$ .

We will show in the following the difficulty of constructing such a matrix  $A$  in the general case (§2), and the simplifications appearing when one imposes

conditions on the invariants  $I, J, K$  and on the indicatrix foliation (§3, §4). Our conclusion is that subject to some geometrical conditions there are very many possibilities for the matrix  $A$  in the cases  $K = 1$  and  $J = 0$ , respectively.

Moreover, a Cartan-Kähler analysis shows that foliation conditions imposed in §4 determine the form of the Riemannian structure on the manifold  $\Lambda$  (see §4, §5 for notations). Namely, we ended up with a surface of revolution  $(\Lambda, g)$  whose unit sphere bundle admits an  $(I, 0, K)$ -generalized Finsler structure. This is the same structure as the one constructed by us in [12]. The present study (§6, §7) shows that the manifold of geodesics of the Riemannian structure  $(\Lambda, g)$  is not a smooth manifold, therefore a classical Landsberg structure on a smooth surface cannot be obtained by this construction, giving in this way a negative answer to our previous conjecture in [12].

Our findings do not rule out the existence of classical Landsberg structures on smooth surfaces, but only suggest that a different approach should be chosen.

## 2. The most general case

Let us assume that we have an  $R$ -Cartan structure  $\{\alpha^1, \alpha^2, \alpha^3\}$  on a 3-manifold  $\Sigma$  and an  $(I, J, K)$ -structure  $\{\omega^1, \omega^2, \omega^3\}$  on a 3-manifold  $\Sigma'$ . We are concerned with the problem of the existence of the linear transformation

$$A : T^*\Sigma \rightarrow T^*\Sigma'$$

given by  $\omega = A\alpha$ . Without loss of generality we can assume that there is a diffeomorphism  $\varphi : \Sigma' \rightarrow \Sigma$  and that  $A$  is the matrix of the cotangent map  $\varphi^*$ . Therefore, formally, we identify the manifolds  $\Sigma'$  and  $\Sigma$ .

In components, we have

$$(6) \quad \begin{pmatrix} \omega^1 \\ \omega^2 \\ \omega^3 \end{pmatrix} = \begin{pmatrix} a_{11} & a_{12} & a_{13} \\ a_{21} & a_{22} & a_{23} \\ a_{31} & a_{32} & a_{33} \end{pmatrix} \begin{pmatrix} \alpha^1 \\ \alpha^2 \\ \alpha^3 \end{pmatrix},$$

where the elements  $a_{ij}$  of the matrix  $A$  are smooth functions on the 3-manifold  $\Sigma$ . Recall that an  $R$ -Cartan structure has the structure equations (3) with the structure function  $R$ .

One can easily see that  $\omega^1 \wedge \omega^2 \wedge \omega^3 = \det(A) \alpha^1 \wedge \alpha^2 \wedge \alpha^3$ .

By taking the exterior derivative of (6), it follows

$$(7) \quad d \begin{pmatrix} \omega^1 \\ \omega^2 \\ \omega^3 \end{pmatrix} = dA \wedge \begin{pmatrix} \alpha^1 \\ \alpha^2 \\ \alpha^3 \end{pmatrix} + Ad \begin{pmatrix} \alpha^1 \\ \alpha^2 \\ \alpha^3 \end{pmatrix}.$$

Using now the structure equations of the  $(I, J, K)$ -structure  $\omega$  and relation (6), we can express the left hand side of (7), with respect to the coframe  $\alpha$ , as follows

$$(8) \quad \begin{aligned} d\omega^1 &= (-IA_{23} + A_{13})\alpha^1 \wedge \alpha^2 + (-IA_{22} + A_{12})\alpha^1 \wedge \alpha^3 \\ &\quad + (-IA_{21} + A_{11})\alpha^2 \wedge \alpha^3, \end{aligned}$$

$$\begin{aligned} d\omega^2 &= -A_{23}\alpha^1 \wedge \alpha^2 - A_{22}\alpha^1 \wedge \alpha^3 - A_{21}\alpha^2 \wedge \alpha^3, \\ d\omega^3 &= (KA_{33} - JA_{23})\alpha^1 \wedge \alpha^2 + (KA_{32} - JA_{22})\alpha^1 \wedge \alpha^3 \\ &\quad + (KA_{31} - JA_{21})\alpha^2 \wedge \alpha^3, \end{aligned}$$

where  $A_{ij}$  is the minor of the element  $a_{ij}$  in the matrix  $A$ , for all  $i, j = 1, 2, 3$ , namely  $A_{11} = a_{22}a_{33} - a_{23}a_{32}$ ,  $A_{12} = a_{21}a_{33} - a_{23}a_{31}$ , etc.

If we denote the directional derivatives with respect to the coframe  $\alpha$  by subscripts, namely  $df = f_1\alpha^1 + f_2\alpha^2 + f_3\alpha^3$ , for any smooth function  $f$  on  $\Sigma$ , and  $da_{ij} = a_{ij,1}\alpha^1 + a_{ij,2}\alpha^2 + a_{ij,3}\alpha^3$  for matrices, then from (7) and (8), we obtain the relations

$$(9) \quad \begin{cases} \begin{cases} a_{11,2} = a_{12,1} + a_{13}R + IA_{23} - A_{13}, \\ a_{21,2} = a_{22,1} + a_{23}R + A_{23}, \\ a_{31,2} = a_{32,1} + a_{33}R - KA_{33} + JA_{23}, \end{cases} \\ \begin{cases} a_{13,1} = a_{11,3} + a_{12} - IA_{22} + A_{12}, \\ a_{23,1} = a_{21,3} + a_{22} - A_{22}, \\ a_{33,1} = a_{31,3} + a_{32} + KA_{32} - JA_{22}, \end{cases} \\ \begin{cases} a_{12,3} = a_{13,2} + a_{11} + IA_{21} - A_{11}, \\ a_{22,3} = a_{23,2} + a_{21} + A_{21}, \\ a_{32,3} = a_{33,2} + a_{31} - KA_{31} + JA_{21}. \end{cases} \end{cases}$$

Next, we will express the system of directional PDEs (9), with respect to the coframe  $\alpha$ , using exterior differential systems (see for example [11] for details). Let us formally construct the following one forms:

$$(10) \quad \begin{cases} \begin{cases} \theta_{11} = d(a_{11}) - a_{11,1}\alpha^1 - (a_{12,1} + B_{112})\alpha^2 - a_{11,3}\alpha^3, \\ \theta_{21} = d(a_{21}) - a_{21,1}\alpha^1 - (a_{22,1} + B_{212})\alpha^2 - a_{21,3}\alpha^3, \\ \theta_{31} = d(a_{31}) - a_{31,1}\alpha^1 - (a_{32,1} + B_{312})\alpha^2 - a_{31,3}\alpha^3, \end{cases} \\ \begin{cases} \theta_{12} = d(a_{12}) - a_{12,1}\alpha^1 - a_{12,2}\alpha^2 - (a_{13,2} + B_{123})\alpha^3, \\ \theta_{22} = d(a_{22}) - a_{22,1}\alpha^1 - a_{22,2}\alpha^2 - (a_{23,2} + B_{223})\alpha^3, \\ \theta_{32} = d(a_{32}) - a_{32,1}\alpha^1 - a_{32,2}\alpha^2 - (a_{33,2} + B_{323})\alpha^3, \end{cases} \\ \begin{cases} \theta_{13} = d(a_{13}) - (a_{11,3} + B_{131})\alpha^1 - a_{13,2}\alpha^2 - a_{13,3}\alpha^3, \\ \theta_{23} = d(a_{23}) - (a_{21,3} + B_{231})\alpha^1 - a_{23,2}\alpha^2 - a_{23,3}\alpha^3, \\ \theta_{33} = d(a_{33}) - (a_{31,3} + B_{331})\alpha^1 - a_{33,2}\alpha^2 - a_{33,3}\alpha^3, \end{cases} \end{cases}$$

where we use the notations

$$\begin{aligned} B_{112} &:= a_{13}R + IA_{23} - A_{13}, & B_{212} &:= a_{23}R + A_{23}, & B_{312} &:= a_{33}R - KA_{33} + JA_{23}, \\ B_{123} &:= a_{11} + IA_{21} - A_{11}, & B_{223} &:= a_{21} + A_{21}, & B_{323} &:= a_{31} - KA_{31} + JA_{21}, \\ B_{131} &:= a_{12} - IA_{22} + A_{12}, & B_{231} &:= a_{22} - A_{22}, & B_{331} &:= a_{32} + KA_{32} - JA_{22}. \end{aligned}$$

We can now formulate our setting as follows. Consider the jet space

$$(x, y, p; a_{ij}; a_{ij \cdot k})_{i,j,k \in \{1,2,3\}} \in J^1(3, 9) = \Sigma^3 \times \mathbb{R}^9 \times \mathbb{R}^{27},$$

where  $(x, y, p)$  are some local coordinates on  $\Sigma$ . Next, we consider the codimension 9 submanifold  $\tilde{\Sigma}^{30}$  of  $J^1(3, 9)$  defined by the equations (9), then we can take the local coordinates

$$(x, y, p; a_{ij}; \bar{a}_{ij \cdot k})_{i,j,k \in \{1,2,3\}} \in \tilde{\Sigma}^{30} = \Sigma^3 \times \mathbb{R}^9 \times \mathbb{R}^{18},$$

where  $\bar{a}_{ij \cdot k}$  means all  $a_{ij \cdot k}$  for  $i, j, k \in \{1, 2, 3\}$  less the terms in the left hand side of (9).

Let us consider the Pfaffian system

$$(11) \quad \mathcal{I} := \{\theta_{ij}\}_{i,j,k \in \{1,2,3\}} \subset \mathcal{J} := \{\theta_{ij}; \alpha^1, \alpha^2, \alpha^3\}_{i,j,k \in \{1,2,3\}} \subset T^*\tilde{\Sigma}$$

with independence condition  $\alpha^1 \wedge \alpha^2 \wedge \alpha^3 \neq 0$ , where the one forms  $\theta_{ij}$  are given in (10).

The integral manifolds of the linear Pfaffian system  $(\mathcal{I}, \mathcal{J})$  are the elements of the matrix  $A$ .

Therefore, we are going to study the existence of integral manifolds of the linear Pfaffian system  $(\mathcal{I}, \mathcal{J})$  defined above.

We start by computing exterior differentials of the one forms  $\theta_{ij}$ :

$$\begin{aligned} d\theta_{11} = & -d(a_{11 \cdot 1}) \wedge \alpha^1 - d(a_{12 \cdot 1}) \wedge \alpha^2 - d(a_{11 \cdot 3}) \wedge \alpha^3 \\ & - (a_{11 \cdot 1} - B_{112 \cdot 3})\alpha^2 \wedge \alpha^3 - (a_{12 \cdot 1} + B_{112})\alpha^3 \wedge \alpha^1 \\ & - (a_{11 \cdot 3}R + B_{112 \cdot 1})\alpha^1 \wedge \alpha^2. \end{aligned}$$

It follows

$$(12) \quad d\theta_{11} = \Pi_{111} \wedge \alpha^1 + \Pi_{121} \wedge \alpha^2 + \Pi_{113} \wedge \alpha^3,$$

where

$$(13) \quad \begin{aligned} \Pi_{111} & := -d(a_{11 \cdot 1}) - (a_{12 \cdot 1} + B_{112})\alpha^3 + (a_{11 \cdot 3}R + B_{112 \cdot 1})\alpha^2, \\ \Pi_{121} & := -d(a_{12 \cdot 1}), \\ \Pi_{113} & := -d(a_{11 \cdot 3}) - (a_{11 \cdot 1} - B_{112 \cdot 3})\alpha^2. \end{aligned}$$

In the same way, we obtain

$$(14) \quad d\theta_{21} = \Pi_{211} \wedge \alpha^1 + \Pi_{221} \wedge \alpha^2 + \Pi_{213} \wedge \alpha^3,$$

where

$$(15) \quad \begin{aligned} \Pi_{211} & := -d(a_{21 \cdot 1}) - (a_{22 \cdot 1} + B_{212})\alpha^3 + (a_{21 \cdot 3}R + B_{212 \cdot 1})\alpha^2, \\ \Pi_{221} & := -d(a_{22 \cdot 1}), \\ \Pi_{213} & := -d(a_{21 \cdot 3}) - (a_{21 \cdot 1} - B_{212 \cdot 3})\alpha^2, \end{aligned}$$

and

$$(16) \quad d\theta_{31} = \Pi_{311} \wedge \alpha^1 + \Pi_{321} \wedge \alpha^2 + \Pi_{313} \wedge \alpha^3,$$



where

$$(17) \quad \begin{aligned} \Pi_{311} &:= -d(a_{31.1}) - (a_{32.1} + B_{312})\alpha^3 + (a_{31.3}R + B_{312.1})\alpha^2, \\ \Pi_{321} &:= -d(a_{32.1}), \\ \Pi_{313} &:= -d(a_{31.3}) - (a_{31.1} - B_{312.3})\alpha^2. \end{aligned}$$

Therefore, we conclude

$$(18) \quad \begin{cases} d\theta_{11} = \Pi_{111} \wedge \alpha^1 + \Pi_{121} \wedge \alpha^2 + \Pi_{113} \wedge \alpha^3, \\ d\theta_{21} = \Pi_{211} \wedge \alpha^1 + \Pi_{221} \wedge \alpha^2 + \Pi_{213} \wedge \alpha^3, \\ d\theta_{31} = \Pi_{311} \wedge \alpha^1 + \Pi_{321} \wedge \alpha^2 + \Pi_{313} \wedge \alpha^3. \end{cases}$$

We move now to the second group of  $\theta$ 's and by similar computations we conclude

$$(19) \quad \begin{cases} d\theta_{12} = \Pi_{121} \wedge \alpha^1 + \Pi_{122} \wedge \alpha^2 + \Pi_{132} \wedge \alpha^3, \\ d\theta_{22} = \Pi_{221} \wedge \alpha^1 + \Pi_{222} \wedge \alpha^2 + \Pi_{232} \wedge \alpha^3, \\ d\theta_{32} = \Pi_{321} \wedge \alpha^1 + \Pi_{322} \wedge \alpha^2 + \Pi_{332} \wedge \alpha^3, \end{cases}$$

where  $\Pi_{121}, \Pi_{221}, \Pi_{321}$  are given above, and

$$(20) \quad \begin{aligned} \Pi_{122} &:= -d(a_{12.2}) + (a_{12.1} + B_{123.2})\alpha^3 - (a_{13.2} + B_{123})R\alpha^1, \\ \Pi_{132} &:= -d(a_{13.2}) + (a_{12.2} - B_{123.1})\alpha^1, \\ \Pi_{222} &:= -d(a_{22.2}) + (a_{22.1} + B_{223.2})\alpha^3 - (a_{23.2} + B_{223})R\alpha^1, \\ \Pi_{232} &:= -d(a_{23.2}) + (a_{22.2} - B_{223.1})\alpha^1, \\ \Pi_{322} &:= -d(a_{32.2}) + (a_{32.1} + B_{323.2})\alpha^3 - (a_{33.2} + B_{323})R\alpha^1, \\ \Pi_{332} &:= -d(a_{33.2}) + (a_{32.2} - B_{323.1})\alpha^1. \end{aligned}$$

Finally, we consider the third group of  $\theta$ 's and obtain:

$$(21) \quad \begin{cases} d\theta_{13} = \Pi_{113}\alpha^1 + \Pi_{132}\alpha^2 + \Pi_{133}\alpha^3 + T_{13}\alpha^1 \wedge \alpha^2, \\ d\theta_{23} = \Pi_{213}\alpha^1 + \Pi_{232}\alpha^2 + \Pi_{233}\alpha^3 + T_{23}\alpha^1 \wedge \alpha^2, \\ d\theta_{33} = \Pi_{313}\alpha^1 + \Pi_{332}\alpha^2 + \Pi_{333}\alpha^3 + T_{33}\alpha^1 \wedge \alpha^2, \end{cases}$$

where  $\Pi_{113}, \Pi_{132}, \Pi_{213}, \Pi_{232}, \Pi_{313}, \Pi_{332}$  are given above, and

$$(22) \quad \begin{aligned} \Pi_{133} &:= -d(a_{13.3}) - (a_{11.3} + B_{131})\alpha^2 + (a_{13.2} + B_{131.3})\alpha^1, \\ \Pi_{233} &:= -d(a_{23.3}) - (a_{21.3} + B_{231})\alpha^2 + (a_{23.2} + B_{231.3})\alpha^1, \\ \Pi_{333} &:= -d(a_{33.3}) - (a_{31.3} + B_{331})\alpha^2 + (a_{33.2} + B_{331.3})\alpha^1, \\ T_{13} &:= -(a_{11.1} - B_{112.3}) - (a_{12.2} - B_{123.1}) - (a_{13.3}R - B_{131.2}), \\ T_{23} &:= -(a_{21.1} - B_{212.3}) - (a_{22.2} - B_{223.1}) - (a_{23.3}R - B_{231.2}), \\ T_{33} &:= -(a_{31.1} - B_{312.3}) - (a_{32.2} - B_{323.1}) - (a_{33.3}R - B_{331.2}). \end{aligned}$$

Putting all these together, we get the  $\theta$ 's structure equations

$$(23) \quad \begin{cases} d\theta_{11} = \Pi_{111} \wedge \alpha^1 + \Pi_{121} \wedge \alpha^2 + \Pi_{113} \wedge \alpha^3, \\ d\theta_{12} = \Pi_{121} \wedge \alpha^1 + \Pi_{122} \wedge \alpha^2 + \Pi_{132} \wedge \alpha^3, \\ d\theta_{13} = \Pi_{113} \wedge \alpha^1 + \Pi_{132} \wedge \alpha^2 + \Pi_{133} \wedge \alpha^3 + T_{13}\alpha^1 \wedge \alpha^2, \\ \\ d\theta_{21} = \Pi_{211} \wedge \alpha^1 + \Pi_{221} \wedge \alpha^2 + \Pi_{213} \wedge \alpha^3, \\ d\theta_{22} = \Pi_{221} \wedge \alpha^1 + \Pi_{222} \wedge \alpha^2 + \Pi_{232} \wedge \alpha^3, \\ d\theta_{23} = \Pi_{213} \wedge \alpha^1 + \Pi_{232} \wedge \alpha^2 + \Pi_{233} \wedge \alpha^3 + T_{23}\alpha^1 \wedge \alpha^2, \\ \\ d\theta_{31} = \Pi_{311} \wedge \alpha^1 + \Pi_{321} \wedge \alpha^2 + \Pi_{313} \wedge \alpha^3, \\ d\theta_{32} = \Pi_{321} \wedge \alpha^1 + \Pi_{322} \wedge \alpha^2 + \Pi_{332} \wedge \alpha^3, \\ d\theta_{33} = \Pi_{313} \wedge \alpha^1 + \Pi_{332} \wedge \alpha^2 + \Pi_{333} \wedge \alpha^3 + T_{33}\alpha^1 \wedge \alpha^2. \end{cases}$$

The torsion terms are

$$(24) \quad \begin{aligned} T_{13} &= -(A_{13.3} + A_{11.1} - A_{12.2}) + I(A_{23.3} + A_{21.1} - A_{22.2}) \\ &\quad + (I_1 A_{21} - I_2 A_{22} + I_3 A_{23}), \\ T_{23} &= A_{23.3} + A_{21.1} - A_{22.2}, \\ T_{33} &= -K(A_{33.3} + A_{31.1} - A_{32.2}) + J(A_{23.3} + A_{21.1} - A_{22.2}) \\ &\quad - (K_1 A_{31} - K_2 A_{32} + K_3 A_{33}) + (J_1 A_{21} - J_2 A_{22} + J_3 A_{23}), \end{aligned}$$

where we have used that  $R_3 = 0$ .

*Remark 2.1.* Indeed, since  $\{\alpha^1, \alpha^2, \alpha^3\}$  is an  $R$ -Cartan structure, from the last structure equation  $d\alpha^3 = R\alpha^1 \wedge \alpha^2$ , we get  $d^2\alpha^3 = R_3\alpha^1 \wedge \alpha^2 \wedge \alpha^3 = 0$ . Since the volume form  $\alpha^1 \wedge \alpha^2 \wedge \alpha^3$  does not vanish it follows  $R_3 = 0$ .

We can therefore conclude:

**Theorem 2.1.** *For a given  $R$ -Cartan structure  $(\Sigma, \alpha)$  and an  $(I, J, K)$ -structure  $(\Sigma, \omega)$ , let  $A : T^*\Sigma \rightarrow T^*\Sigma$  be the linear transformation whose matrix elements  $(a_{ij})$  are solutions of the PDE (9). Let  $(\mathcal{I}, \mathcal{J}) = (\theta^{ij}, \alpha^i)$  be the corresponding Pfaffian system on the 30 dimensional manifold  $\tilde{\Sigma}$ . Then,  $(\mathcal{I}, \mathcal{J})$  has non absorbable torsion, namely  $T_{13}, T_{23}, T_{33}$  given by (24).*

We remark here that

$$(25) \quad T_{13} = 0, \quad T_{23} = 0, \quad T_{33} = 0$$

is the necessary condition for the Pfaffian system  $(\mathcal{I}, \mathcal{J})$  to have solution. Following Cartan-Kähler algorithm (see for example [11]), we should restrict the Pfaffian system to the submanifold

$$\Xi := \{u \in \tilde{\Sigma} \mid T_{13}(u) = 0, T_{23}(u) = 0, T_{33}(u) = 0\} \subset \tilde{\Sigma}$$

and compute again the torsion. However, the (25) is a very complicated nonlinear directional PDE, and even after restriction to  $\Xi$ , new non-vanishing torsions are expected to appear, so the problem is not tractable in this general case.

Therefore, we will consider in the following some particular cases when by putting some geometrical conditions we can actually decide the existence of such a linear transformation between cotangent spaces.

### 3. The case $K = 1$

We consider now the similar problem of an *a-priori* given  $R$ -Cartan structure  $\alpha$  and an  $(I, J, 1)$ -structure  $\omega$  on the same 3-manifold  $\Sigma$ . We are again concerned with the problem of the existence of the linear transformation  $A : T^*\Sigma \rightarrow T^*\Sigma$  such that  $\omega = A\alpha$ . Recall that in this case  $(I, J, 1)$ -generalized Finsler structure has the structure equations

$$(26) \quad \begin{aligned} d\omega^1 &= -I\omega^1 \wedge \omega^3 + \omega^2 \wedge \omega^3, \\ d\omega^2 &= -\omega^1 \wedge \omega^3, \\ d\omega^3 &= \omega^1 \wedge \omega^2 - J\omega^1 \wedge \omega^3, \end{aligned}$$

and the Bianchi identities

$$(27) \quad J = I\omega_2, \quad I + J\omega_2 = 0,$$

where the covariant derivatives are with respect to the  $(I, J, 1)$ -structure  $(\Sigma, \omega)$ .

Let us remark that the most important feature of the case  $K = 1$  is the invariance of the two forms  $\omega^1 \wedge \omega^3$  and  $(\omega^1)^2 + (\omega^3)^2$  under the flow of  $\hat{e}_2$ .

Indeed, a straightforward computation shows that

$$(28) \quad \begin{aligned} \mathcal{L}_{\hat{e}_2}[(\omega^1)^2 + (\omega^3)^2] &= 0, & \mathcal{L}_{\hat{e}_2}(\omega^1 \wedge \omega^3) &= 0, \\ \mathcal{L}_{\hat{e}_2}(I\omega^1 + J\omega^3) &= 0, & \mathcal{L}_{\hat{e}_2}(I^2 + J^2) &= 0. \end{aligned}$$

The geometrical meaning of these relations is quite clear.

If  $U \subset \Sigma$  is an open set where the geodesic foliation  $\{\omega^1 = 0, \omega^3 = 0\}$  is amenable (this is always possible locally), i.e.,  $\Lambda_U := U/\{\omega^1=0, \omega^3=0\}$  is a differential manifold and  $l_U : U \rightarrow \Lambda_U$  is a smooth submersion, then there exist

- (1) a quadratic form  $d\sigma^2$  on  $\Lambda_U$  such that  $l_U^*(d\sigma^2) = (\omega^1)^2 + (\omega^3)^2$ ;
- (2) a two form  $\Omega$  on  $\Lambda_U$  such that  $l_U^*(\Omega) = \omega^1 \wedge \omega^3$ ;
- (3) a one form  $\beta$  on  $\Lambda_U$  such that  $l_U^*(\beta) = I\omega^1 + J\omega^3$ .

The indicatrix foliation of this Riemannian structure coincides with the geodesic foliation of the  $(I, J, 1)$ -structure.

In particular, if the  $(I, J, 1)$ -structure  $(\Sigma, \omega)$  is geodesically amenable, i.e.,  $\Lambda := \Sigma/\{\omega^1=0, \omega^3=0\}$  is a differential manifold and  $l : \Sigma \rightarrow \Lambda$  is a smooth submersion, then on  $\Lambda$  there exist a canonical Riemannian metric  $g$ , with area form  $\Omega$  and a one form  $\beta$  whose norm with respect to the Riemannian metric  $g$  is  $I^2 + J^2$ .

On the other hand, we have given an  $R$ -Cartan structure  $(\Sigma, \alpha)$  on the same manifold  $\Sigma$  whose indicatrix foliation is  $\{\alpha^1 = 0, \alpha^2 = 0\}$ . Locally, if  $V$  is a small neighborhood on  $\Sigma$ , where this foliation is amenable, then  $\Lambda_V := V/\{\alpha^1=0, \alpha^2=0\}$  is a differential manifold and  $l_V : V \rightarrow \Lambda_V$  is a smooth submersion. Moreover,  $\Lambda_V$  is naturally endowed with

- (1) a quadratic form  $d\bar{\sigma}^2$ , such that  $l_V^*(d\bar{\sigma}^2) = (\alpha^1)^2 + (\alpha^2)^2$ ;
- (2) a two form  $\bar{\Omega}$ , such that  $l_V^*(\bar{\Omega}) = \alpha^1 \wedge \alpha^2$ .

In the case when the Cartan structure  $(\Sigma, \alpha)$  is amenable,  $\bar{\Lambda} := \Sigma / \{\alpha^1=0, \alpha^2=0\}$  is a differential manifold and  $\bar{l} : \Sigma \rightarrow \bar{\Lambda}$  is a smooth submersion and on  $\bar{\Lambda}$  there exists a canonical Riemannian metric  $\bar{g}$ , with area form  $\bar{\Omega}$ .

Due to all these it is natural to impose the following:

**Foliation Condition:** *The  $(I, J, 1)$ -structure's geodesic foliation  $\{\omega^1 = 0, \omega^3 = 0\}$  coincides to the indicatrix foliation of the given  $R$ -Cartan structure  $\{\alpha^1 = 0, \alpha^2 = 0\}$ .*

This condition has the following natural:

**Corollary 3.1.** *The  $(I, J, 1)$ -structure  $(\Sigma, \omega)$  is geodesically amenable if and only if the  $R$ -Cartan structure  $(\Sigma, \alpha)$  is amenable.*

Hereafter we consider the case of an  $R$ -Cartan structure  $(\Sigma, \omega)$  and an  $(I, J, 1)$ -structure  $(\Sigma, \alpha)$  under the above foliation condition. In terms of the transformation  $A : T^*\Sigma \rightarrow T^*\Sigma$ , this implies  $a_{13} = 0, a_{33} = 0$ , therefore we have

$$(29) \quad \begin{pmatrix} \omega^1 \\ \omega^2 \\ \omega^3 \end{pmatrix} = \begin{pmatrix} a_{11} & a_{12} & 0 \\ a_{21} & a_{22} & a_{23} \\ a_{31} & a_{32} & 0 \end{pmatrix} \begin{pmatrix} \alpha^1 \\ \alpha^2 \\ \alpha^3 \end{pmatrix}.$$

Another useful geometrical remark is that, in the case when  $(\Sigma, \omega)$  is geodesically amenable, or, equivalently,  $(\Sigma, \alpha)$  is amenable, on the 2-manifold  $\Lambda = \Sigma / \{\omega^1=0, \omega^3=0\} = \Sigma / \{\alpha^1=0, \alpha^2=0\}$  we have actually two different quadratic forms  $d\sigma^2$  and  $d\bar{\sigma}^2$  such that  $l^*(d\sigma^2) = (\omega^1)^2 + (\omega^3)^2$  and  $l^*(d\bar{\sigma}^2) = (\alpha^1)^2 + (\alpha^2)^2$ , respectively. In the case of amenability, they induce two different Riemannian metrics on the 2-manifold  $\Lambda$ .

A natural choice is given by the following choice:

**Identification Condition:** *The Cartan structure  $\{\omega^1, \omega^3\}$  coincides with the Cartan structure  $\{\alpha^1, \alpha^2\}$ , namely,*

$$(30) \quad \omega^1 = \alpha^1 \quad \omega^3 = \alpha^2.$$

In this case, we obtain the coframe changing

$$\begin{pmatrix} \omega^1 \\ \omega^2 \\ \omega^3 \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ a_{21} & a_{22} & a_{23} \\ 0 & 1 & 0 \end{pmatrix} \begin{pmatrix} \alpha^1 \\ \alpha^2 \\ \alpha^3 \end{pmatrix},$$

and from the linear independence condition of the one forms  $\alpha$  and  $\omega$ , namely,

$$\omega^1 \wedge \omega^2 \wedge \omega^3 = -a_{23} \alpha^1 \wedge \alpha^2 \wedge \alpha^3 \neq 0$$

it follows  $a_{23} \neq 0$ .

The first structure equation of (26) implies now

$$a_{23} = -1, \quad a_{21} = I.$$

The third structure equation of (26) implies further

$$a_{22} = J,$$

so we obtain the coframe changing

$$(31) \quad \begin{pmatrix} \omega^1 \\ \omega^2 \\ \omega^3 \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ I & J & -1 \\ 0 & 1 & 0 \end{pmatrix} \begin{pmatrix} \alpha^1 \\ \alpha^2 \\ \alpha^3 \end{pmatrix}.$$

Similarly, the second structure equation in (26) leads us to the condition

$$(32) \quad R = -I_2 + J_1 + 1,$$

where the subscripts denote directional derivative with respect to the coframe  $\alpha$ , or, equivalently,

$$R = -(J^2 + I_{\omega^3}) + J_{\omega^1} - I^2 + 1.$$

Hence, we can conclude:

**Proposition 3.2.** *Let  $\alpha$  and  $\omega$  be an  $R$ -Cartan and an  $(I, J, 1)$ -generalized Finsler structure on the 3-manifold  $\Sigma$ , respectively. Then, there is a unique linear transformation  $A : T^*\Sigma \rightarrow T^*\Sigma$ , such that  $\omega = A\alpha$ , subject to the foliation and identification conditions above. The matrix of  $A$  in the coframes  $\omega$  and  $\alpha$  is given in (31). In this case, the  $R$ -Cartan structure must obey the condition (32).*

*Remark.* Since we are in two dimensions and amenability assumed, these Riemannian metrics must be conformal, and therefore their tautological forms  $\{\alpha^1, \alpha^2\}$  and  $\{\omega^1, \omega^3\}$  must be conformal as well, respectively. Here we regard the one forms  $\{\omega^1, \omega^3\}$  and  $\{\alpha^1, \alpha^2\}$  as the tautological one forms of the corresponding Riemannian structures on the same 3-manifold  $\Sigma$ , respectively, up to a local diffeomorphism. Indeed, the orthogonal frame bundles of the Riemannian structures  $(\Lambda, g)$  and  $(\Lambda, \bar{g})$  are different, but it is always possible to define local diffeomorphisms between  $\Sigma$  and the total spaces of these orthonormal frame bundles (see [12] for details).

Therefore, without loss of generality, we can always assume the following:

**Conformal Equivalence:**

$$(33) \quad \begin{pmatrix} \omega^1 \\ \omega^3 \end{pmatrix} = m \begin{pmatrix} \alpha^1 \\ \alpha^2 \end{pmatrix},$$

where  $m$  is a positive function defined on  $\Sigma$ . From the structure equations of the  $R$ -Cartan structure it follows immediately that  $m_3 = 0$ .

Let us point out this conformal equivalence holds in the non-amenable case as well. Indeed, it can be seen from the structure equations of the  $(I, J, 1)$ -structure that  $\{\omega^1, \omega^3\}$  is a Cartan structure as well (pay attention to the fact

that the third one form of this Cartan structure is different from  $\omega^2$ ). In this case, under the foliation condition, the induced quadratic forms live on the 2-orbifold  $\Lambda$  defined similarly as in the amenable case. In this case relation (33) simply says that the two Cartan structures  $\{\omega^1, \omega^3\}$  and  $\{\alpha^1, \alpha^2\}$  belong to the same conformal class of Cartan structures (see [7] for details).

*Remark.* One can also easily see that for a smooth function  $v : \Sigma \rightarrow \mathbb{R}$ , if  $\{\alpha^1, \alpha^2, \alpha^3\}$  is an  $R$ -Cartan structure  $\Sigma$ , and  $v$  a nowhere zero function, then  $\{v\alpha^1, v\alpha^2, \tilde{\alpha}^3\}$  is an  $\tilde{R}$ -Cartan structure, where

$$(34) \quad \tilde{\alpha}^3 := \alpha^3 - *d(\log v), \quad v^2 \tilde{R} = R - \Delta_\alpha(\log v),$$

provided  $v_3 = 0$ . Here  $*$  and  $\Delta_\alpha$  are the Hodge operator and the Laplacian of the Cartan structure  $(\alpha^1, \alpha^2)$ , namely

$$(35) \quad *d(\log v) = -\frac{v_2}{v}\alpha^1 + \frac{v_1}{v}\alpha^2, \quad \Delta_\alpha(\log v) = \frac{1}{v^2}[(v_{11} + v_{22})v - (v_1^2 + v_2^2)].$$

The conformal class of an  $R$ -Cartan structure  $\{\alpha^1, \alpha^2, \alpha^3\}$  is the collection of all coframes  $\{v\alpha^1, v\alpha^2, \tilde{\alpha}^3\}$ , where  $v$  is a smooth positive function on  $\Sigma$  provided  $v_3 = 0$ .

Therefore, we obtain

$$(36) \quad \begin{pmatrix} \omega^1 \\ \omega^2 \\ \omega^3 \end{pmatrix} = \begin{pmatrix} m & 0 & 0 \\ a_{21} & a_{22} & a_{23} \\ 0 & m & 0 \end{pmatrix} \begin{pmatrix} \alpha^1 \\ \alpha^2 \\ \alpha^3 \end{pmatrix},$$

and hence

$$\omega^1 \wedge \omega^2 \wedge \omega^3 = -m^2 a_{23} \alpha^1 \wedge \alpha^2 \wedge \alpha^3.$$

The first structure equation of the  $(I, J, 1)$ -structure gives

$$(37) \quad \begin{aligned} m_3 &= 0, \\ a_{23} &= -1, \\ a_{21} &= (-m_2 + Im^2)/m. \end{aligned}$$

Remark that first relation was also obtained from geometrical considerations above as well. In this case, we obtain

$$\begin{pmatrix} \omega^1 \\ \omega^2 \\ \omega^3 \end{pmatrix} = \begin{pmatrix} m & 0 & 0 \\ -\frac{m_2}{m} + Im & a_{22} & -1 \\ 0 & m & 0 \end{pmatrix} \begin{pmatrix} \alpha^1 \\ \alpha^2 \\ \alpha^3 \end{pmatrix}.$$

The third structure equation of  $\omega$  implies

$$(38) \quad a_{22} = Jm + \frac{m_1}{m},$$

and therefore we obtain

$$(39) \quad \begin{pmatrix} \omega^1 \\ \omega^2 \\ \omega^3 \end{pmatrix} = \begin{pmatrix} m & 0 & 0 \\ -\frac{m_2}{m} + Im & \frac{m_1}{m} + Jm & -1 \\ 0 & m & 0 \end{pmatrix} \begin{pmatrix} \alpha^1 \\ \alpha^2 \\ \alpha^3 \end{pmatrix}.$$

Remark that we can also write formally

$$\omega^2 = \varphi + *d \log m - \alpha^3,$$

where we put  $\varphi := m(I\alpha^1 + J\alpha^2)$ . The Hodge operator has the usual form  $*d \log m = -\frac{m_2}{m}\alpha^1 + \frac{m_1}{m}\alpha^2$ .

The second structure equation (26) leads us to the following structure equation

$$(40) \quad d\varphi = (R - m^2)\alpha^1 \wedge \alpha^2 - d(*d \log m).$$

We can conclude:

**Theorem 3.3.** *Let  $\alpha$  and  $\omega$  be an  $R$ -Cartan and an  $(I, J, 1)$ -generalized Finsler structure on the 3-manifold  $\Sigma$ , respectively, such that the contact one-forms  $\omega^1, \omega^3$  belong to the conformal class of Cartan structure  $\alpha$ . Then, the linear transformation  $A : T^*\Sigma \rightarrow T^*\Sigma$ , such that  $\omega = A\alpha$ , subject to the Foliation condition, is given by (39), where  $m$  is the conformal factor.*

*Moreover, the structure function  $R$  has to satisfy*

$$(41) \quad R\alpha^1 \wedge \alpha^2 = d\varphi + m^2\alpha^1 \wedge \alpha^2 + d(*d \log m)$$

for the one form  $\varphi = m(I\alpha^1 + J\alpha^2)$ .

Let us remark that the indicatrix foliation of the  $(I, J, 1)$ -structure  $(\Sigma, \omega)$  coincides with the  $(\varphi + *d \log m)$ -geodesic foliation of the Cartan structure  $(\Sigma, \alpha)$ .

We can ask:

**Projective Condition:** *The indicatrix foliation of the  $(I, J, 1)$ -structure  $(\Sigma, \omega)$  coincides with the geodesic foliation of the Cartan structure  $(\Sigma, \alpha)$ .*

This implies that we must have

$$\varphi = - * d \log m,$$

i.e., the invariants of the  $(I, J, 1)$ -structure  $(\Sigma, \omega)$  are uniquely determined by the conformal factor  $m$ , namely,

$$I = \frac{m_2}{m^2}, \quad J = -\frac{m_1}{m^2}.$$

With this extra condition, the structure equation (40) reads

$$d\varphi = -m^2\alpha^1 \wedge \alpha^2 + d\alpha^3 - d * d \log m,$$

or, equivalently

$$d\alpha^3 = m^2\alpha^1 \wedge \alpha^2,$$

in other words, we must have  $m^2 = R$ .

Therefore we have:

**Proposition 3.4.** *Let  $\alpha$  be an  $R$ -Cartan structure,  $R > 0$  on the 3-manifold  $\Sigma$ . Then there is only one  $(I, J, 1)$ -generalized Finsler structure on  $\Sigma$  that satisfies the Foliation and the Projective conditions above, namely the one given in (39), where  $m = \pm\sqrt{R}$  is the conformal factor. Its invariants are uniquely determined only by the structure function  $R$ , namely*

$$(42) \quad I = \frac{1}{2} \frac{R_2}{R^{3/2}}, \quad J = -\frac{1}{2} \frac{R_1}{R^{3/2}}.$$

Let us recall that a Riemannian metric  $g$  on a smooth manifold  $\Lambda$  is called a *Zoll metric* if all its geodesics are simple closed curves of equal length. See [3] for basics of Zoll metrics as well as [10] for the abundance of Zoll metrics on  $S^2$ .

One is led in this way to the following construction method of classical Finsler structures with  $K = 1$  in  $S^2$  found in [5].

**Proposition 3.5.** *Let  $(S^2, g)$  be a Zoll metric of positive curvature  $R$  and let  $\alpha^1, \alpha^2$  be the tautological one-forms,  $\alpha^3$  the Levi-Civita connection one-form on the orthonormal sphere bundle  $\Sigma$  of this Riemannian metric. Then the  $(I, J, 1)$ -generalized Finsler structure given in Proposition 3.4 above gives rise to a classical Finsler structure on the manifold of geodesics of  $g$ .*

*Proof.* The proof is straightforward by construction. □

#### 4. The case $J = 0$

We are concerned in this section with the problem of the existence of the linear transformation  $A : T^*\Sigma \rightarrow T^*\Sigma$  such that  $\omega = A\alpha$ , where  $\alpha$  and  $\omega$  are an  $R$ -Cartan structure and an  $(I, 0, K)$ -generalized Finsler structure on the same 3-manifold  $\Sigma$ , respectively. Recall that structure equations and Bianchi identities for  $\omega$  were already given in Introduction by (4) and (5), respectively.

One can see that it is impossible to extract any two one forms from the coframe  $\omega$  to obtain a Cartan structure as we did in the previous case.

Since all forms in the coframe  $\omega$  are contact forms in this case, the best thing to do is to choose an equivalent contact form from the same contact forms class. Namely, we can choose any of the following pairs of contact forms

$$(43) \quad \begin{pmatrix} m\omega^2 \\ \omega^3 \end{pmatrix} \quad \text{or} \quad \begin{pmatrix} \omega^2 \\ m\omega^3 \end{pmatrix},$$

where  $m$  is some smooth nowhere vanishing function on  $\Sigma$ . More conditions will come up later.

Any of these work, but here we choose the first one.

One can easily see that with a supplementary condition on  $m$ , the quadratic form  $(m\omega^2)^2 + (\omega^3)^2$  is invariant under the flow of  $\hat{e}_1$ . Indeed,

$$\mathcal{L}_{\hat{e}_1}((m\omega^2)^2 + (\omega^3)^2) = 2[mm_{\omega^1}(\omega^2)^2 + (m^2 - K)\omega^2\omega^3].$$

Therefore, we get:



**Lemma 4.1.** *Let  $m$  be a smooth non-vanishing function on  $\Sigma$  satisfying*

$$(44) \quad m_{\omega^1} = 0 \quad m^2 = K.$$

*Then  $\mathcal{L}_{\hat{e}_1}((m\omega^2)^2 + (\omega^3)^2) = 0$ .*

Moreover, we have:

**Lemma 4.2.** *Under the conditions (44) the pair of contact one forms*

$$\eta^1 := m\omega^2, \quad \eta^2 := \omega^3, \quad \eta^3 := m\omega^1 + m_{\omega^3}\omega^2$$

*is an R-Cartan structure on  $\Sigma$  with the structure function  $R = 1 - \frac{m_{\omega^3\omega^3}}{m}$ .*

Let us return to our initial setting and denote again the coframe changing by

$$(45) \quad \begin{pmatrix} \omega^1 \\ m\omega^2 \\ \omega^3 \end{pmatrix} = \begin{pmatrix} a_{11} & a_{12} & a_{13} \\ a_{21} & a_{22} & a_{23} \\ a_{31} & a_{32} & a_{33} \end{pmatrix} \begin{pmatrix} \alpha^1 \\ \alpha^2 \\ \alpha^3 \end{pmatrix},$$

where  $\alpha$  and  $\omega$  are arbitrary *a priori* given.

Similar to the case  $K = 1$  all these suggest that we can put the following:

**Foliation Condition:** *Finslerian indicatrix foliation  $\{\omega^1 = 0, \omega^2 = 0\}$  coincides with the geodesic foliation  $\{\alpha^1 = 0, \alpha^3 = 0\}$  of the R-Cartan structure  $(\Sigma, \alpha)$  (see also [12]).*

This foliation condition implies that the matrix  $A$  reads now

$$(46) \quad A = \begin{pmatrix} a_{11} & a_{12} & a_{13} \\ a_{21} & a_{22} & 0 \\ a_{31} & a_{32} & 0 \end{pmatrix}.$$

We can now consider again:

**Identification Condition:** *The Cartan structure  $\{m\omega^2, \omega^3\}$  coincides with the R-Cartan structure  $\{\alpha^1, \alpha^2\}$ , i.e.,*

$$\alpha^1 = m\omega^2, \quad \alpha^2 = \omega^3,$$

*provided (44) is satisfied.*

In this case, we have

$$(47) \quad \begin{pmatrix} \omega^1 \\ \omega^2 \\ \omega^3 \end{pmatrix} = \begin{pmatrix} a_{11} & a_{12} & a_{13} \\ \frac{1}{m} & 0 & 0 \\ 0 & 1 & 0 \end{pmatrix} \begin{pmatrix} \alpha^1 \\ \alpha^2 \\ \alpha^3 \end{pmatrix}.$$

Writing down the second and third structure equations (4) we are immediately lead to the supplementary conditions

$$a_{12} = 0, \quad a_{13} = \frac{1}{m},$$

and

$$a_{11} = \left(\frac{1}{m}\right)_2.$$

We therefore obtain the coframe change

$$(48) \quad \begin{pmatrix} \omega^1 \\ \omega^2 \\ \omega^3 \end{pmatrix} = \begin{pmatrix} \left(\frac{1}{m}\right)_2 & 0 & \frac{1}{m} \\ \frac{1}{m} & 0 & 0 \\ 0 & 1 & 0 \end{pmatrix} \begin{pmatrix} \alpha^1 \\ \alpha^2 \\ \alpha^3 \end{pmatrix}.$$

This is the same coframe changing as in [12]. The first structure equation in (4) leads to

$$(49) \quad I = -2\frac{m_2}{m}, \quad R = 1 - \frac{m_{22}}{m}.$$

Let us point out that the Landsberg condition  $I_{\omega^2} = 0$  is equivalent to  $m_{12} = 0$ .

We also remark that the indicatrix foliation  $\{\omega^1 = 0, \omega^2 = 0\}$  of the  $(I, 0, K)$ -structure  $(\Sigma, \omega)$  coincides with the geodesic foliation  $\{\alpha^1 = 0, \alpha^3 = 0\}$  of the  $R$ -Cartan structure  $(\Sigma, \alpha)$  by construction.

Therefore, we can conclude:

**Proposition 4.3.** *Let  $\alpha$  and  $\omega$  be an  $R$ -Cartan and an  $(I, 0, K)$ -generalized Finsler structure on the 3-manifold  $\Sigma$ , respectively. Then, there is a unique linear transformation  $A : T^*\Sigma \rightarrow T^*\Sigma$ , such that  $\omega = A\alpha$ , subject to the Foliation and Identification conditions above. The matrix of  $A$  in the coframes  $\omega$  and  $\alpha$  is given in (48). In this case, the  $R$ -Cartan structure must obey the condition*

$$(50) \quad R = 1 - \frac{m_{22}}{m},$$

where  $m$  is a non-vanishing smooth function on  $\Sigma$  such that

$$(51) \quad m_3 = 0, \quad m_{12} = 0.$$

In this case, the invariants of  $\omega$  are

$$(52) \quad I = -2\frac{m_2}{m}, \quad K = m^2.$$

We remark that under the foliation and identification conditions, the structure functions  $I, K$  of  $(\Sigma, \omega)$  as well as the structure function  $R$  of  $(\Sigma, \alpha)$  are uniquely determined by the function  $m$  that has to satisfy the PDE (51). In other words, we may say that for given  $m$  solution of (51), the  $(1 - \frac{m_{22}}{m})$ -Cartan structure  $\alpha$  and  $(-2\frac{m_2}{m}, 0, m^2)$ -structure  $\omega$  structures are uniquely determined by  $m$ .

The involutivity of the PDE made of relations in Proposition 4.3 can now be investigated by means of Cartan-Kähler Theorem.

Let us point out that our formula  $m_{12} = 0$  shows that the function  $m_1$  is invariant under the “geodesic” flow of  $(\Sigma, \alpha)$ . In other words, the function  $m_1$  descends on the leaf space  $\Sigma/\{\alpha^1=0, \alpha^3=0\}$ .

We remark also that locally we can relax a little the identification condition by asking the conformal equivalence of these two Cartan structures, namely:

**Conformal Equivalence:** *There exists a positive smooth function  $f$  on  $\Sigma$ , satisfying  $f_3 = 0$ , such that*

$$\begin{pmatrix} m\omega^2 \\ \omega^3 \end{pmatrix} = f \begin{pmatrix} \alpha^1 \\ \alpha^2 \end{pmatrix}.$$

It follows

$$(53) \quad \begin{pmatrix} \omega^1 \\ \omega^2 \\ \omega^3 \end{pmatrix} = \begin{pmatrix} a_{11} & a_{12} & a_{13} \\ \frac{f}{m} & 0 & 0 \\ 0 & f & 0 \end{pmatrix} \begin{pmatrix} \alpha^1 \\ \alpha^2 \\ \alpha^3 \end{pmatrix}.$$

We will investigate now supplementary conditions for the remaining  $a_{ij}$  such that  $(\Sigma, \alpha)$  and  $(\Sigma, \omega)$  are Cartan and  $(I, 0, K)$ -structure, respectively.

The structure equation  $d\omega^2 = \omega^3 \wedge \omega^1$  leads to

$$(54) \quad \begin{aligned} fa_{11} &= \left(\frac{f}{m}\right)_2, \\ fa_{13} &= \frac{f}{m}, \\ \left(\frac{f}{m}\right)_3 &= 0. \end{aligned}$$

From here it follows that  $a_{11}$  and  $a_{13}$  can be written in terms of  $f$  and  $m$  only, namely

$$(55) \quad a_{11} = \frac{1}{f} \left(\frac{f}{m}\right)_2 = \frac{f_2 m - f m_2}{f m^2}, \quad a_{13} = \frac{1}{m}.$$

The structure equation  $d\omega^3 = K\omega^1 \wedge \omega^2$  leads to

$$(56) \quad a_{12} f m = -f_1, \quad f_3 = 0,$$

and from here we get  $a_{12}$  in terms of  $f$  and  $m$ . Also, by combining with (54) we obtain  $m_3 = 0$ . It can be seen that this condition is consistent with  $m_1 = 0$  obtained already.

With all these we obtain the final form of the coframing change

$$(57) \quad \begin{pmatrix} \omega^1 \\ \omega^2 \\ \omega^3 \end{pmatrix} = \begin{pmatrix} \frac{1}{f} \left(\frac{f}{m}\right)_2 & -\frac{f_1}{f m} & \frac{1}{m} \\ \frac{f}{m} & 0 & 0 \\ 0 & f & 0 \end{pmatrix} \begin{pmatrix} \alpha^1 \\ \alpha^2 \\ \alpha^3 \end{pmatrix}.$$

It follows that the Bianchi equation  $K_{\omega 3} + KI = 0$  implies

$$I = -\frac{K_{\omega 3}}{K} = -2\frac{m_2}{f m}.$$

After some long but not complicated computations, we get that  $I_{\omega 2} = 0$  is equivalent to the following **Landsberg Condition:**

$$(58) \quad m_{21} = \frac{f_1 m_2 + f_2 m_1}{f}.$$

The remaining structure equation  $d\omega^1 = -I\omega^1 \wedge \omega^3 + \omega^2 \wedge \omega^3$  leads now to

$$\begin{aligned}
 & -\left(\frac{1}{f}\left(\frac{f}{m}\right)_2\right)_2 - \left(\frac{f_1}{fm}\right)_1 + \frac{1}{m}R = -I\left(\frac{f}{m}\right)_2 + \frac{f^2}{m}, \\
 (59) \quad & \frac{1}{f}\left(\frac{f}{m}\right)_2 + \left(\frac{f_1}{fm}\right)_3 + \left(\frac{1}{m}\right)_2 = I\frac{f}{m}, \\
 & \left(\frac{1}{m}\right)_1 + \frac{f_1}{fm} - \left(\frac{1}{f}\left(\frac{f}{m}\right)_2\right)_3 = 0.
 \end{aligned}$$

It can be seen that the second and third relations are identities. However first relation leads to second order directional PDE for  $f$  in terms of the curvature  $R$  that should be considered together with the Landsberg condition (58). Indeed, by straightforward computation we get the **curvature condition**:

$$(60) \quad R - f^2 = \frac{f_{11} + f_{22}}{f} - \frac{f_1^2 + f_2^2}{f^2} - \frac{f_1m_1 - f_2m_2}{fm} - \frac{m_{22}}{m}.$$

We can conclude:

**Proposition 4.4.** *Let  $(\Sigma, \alpha)$  and  $(\Sigma, \omega)$  be an  $R$ -Cartan structure and an  $(I, 0, K)$ -structure on the 3-manifold  $\Sigma$ , respectively. Let also  $m$  and  $f$  be two smooth functions on  $\Sigma$  such that they are invariant under the flow of  $\hat{e}_1$  and the conditions (58) and (60) are verified. Then, the matrix  $A$  given in (57) gives a coframing change between the  $R$ -Cartan structure  $\alpha$  and the  $(I, 0, K)$ -structure  $\omega$ .*

It can be shown by the means of Cartan-Kähler theorem that the differential system (58), (60) it is involutive (see [12] for similar discussions), but the computations are too complicated to be given here.

### 5. Cartan-Kähler Theory for an $(I, 0, K)$ -generalized Finsler structure

Let us consider here the problem of involutivity of the PDEs (50), (51) for an unknown function  $m$  on the 3-dimensional manifold  $\Sigma$ .

In order to keep things simple, our setting is a 3-manifold  $\Sigma$  endowed with an  $R$ -Cartan structure  $\{\alpha^1, \alpha^2, \alpha^3\}$ . All the subscripts of the scalars are again with respect to the coframe  $\alpha$ .

Let us observe that for any function  $f$  on  $\Sigma$  equipped with an  $R$ -Cartan structure, the following Ricci identities hold:

$$\begin{aligned}
 (61) \quad & f_1 + f_{32} - f_{23} = 0, \\
 & f_2 + f_{13} - f_{31} = 0, \\
 & Rf_3 + f_{21} - f_{12} = 0.
 \end{aligned}$$

In particular, from Bianchi identity  $R_3 = 0$  (see Remark 2.1) we have

$$(62) \quad \begin{aligned} R_1 - R_{23} &= 0, \\ R_2 + R_{13} &= 0, \\ R_{21} - R_{12} &= 0. \end{aligned}$$

We start by considering the following one forms

$$(63) \quad \begin{aligned} \theta_1 &= dm - m_1\alpha^1 - m_2\alpha^2, \\ \theta_2 &= dm_1 - m_{11}\alpha^1 + m_2\alpha^3, \\ \theta_3 &= dm_2 - (1 - R)m\alpha^2 - m_1\alpha^3, \end{aligned}$$

namely, for a given  $R$ -Cartan structure  $\{\alpha^1, \alpha^2, \alpha^3\}$  on the 3-manifold  $\Sigma$ , we consider PDEs (50), (51) and Ricci identities  $m_{13} = -m_2, m_{21} = 0, m_{23} = m_1$ .

Next, let us consider the linear Pfaffian  $(\mathcal{I}, \mathcal{J})$  with independence condition

$$(64) \quad \begin{aligned} \mathcal{I} &= \{\theta_1, \theta_2, \theta_3\}, \\ \mathcal{J} &= \{\alpha^1, \alpha^2, \alpha^3, \theta_1, \theta_2, \theta_3\}, \\ \alpha^1 \wedge \alpha^2 \wedge \alpha^3 &\neq 0 \end{aligned}$$

on the 7-dimensional manifold  $\Omega$  with the coframing

$$(65) \quad \{\alpha^1, \alpha^2, \alpha^3, dm, dm_1, dm_2, dm_{11}\}.$$

The integral manifolds of  $(\mathcal{I}, \mathcal{J})$  are functions  $m$  with the desired proprieties for given  $R$ -Cartan structure.

A straightforward computation shows that  $(\mathcal{I}, \mathcal{J})$  has non-absorbable torsion given by

$$(66) \quad \begin{aligned} d\theta_1 &\equiv 0, \\ d\theta_2 &\equiv (-dm_{11} - m_2R\alpha_2) \wedge \alpha^1 + t_2\alpha^2 \wedge \alpha^3, \quad \text{mod } \theta_1, \theta_2, \theta_3, \\ d\theta_3 &\equiv t_2\alpha^1 \wedge \alpha^3 + t_1\alpha^1 \wedge \alpha^2, \end{aligned}$$

where

$$(67) \quad t_1 := mR_1 - m_1, \quad t_2 := -m_{11} + m(1 - R).$$

Further, we restrict to the 5-dimensional submanifold  $\tilde{\Omega} := \{t_1 = 0, t_2 = 0\} \subset \Omega$  with coframe

$$(68) \quad \{\alpha^1, \alpha^2, \alpha^3, dm, dm_2\}.$$

Then, by a straightforward computation we have

$$\begin{aligned} \theta_2|_{\tilde{\Omega}} &= d(mR_1) - m(1 - R)\alpha^1 + m_2\alpha^3 \\ &\equiv m\{(R_{11} + R_1^2 + R - 1)\alpha^1 + (R_{12} + \frac{R_1m_2}{m})\alpha^2 + (-R_2 + \frac{m_2}{m})\alpha^3\} \\ &\quad \text{mod } \theta_1|_{\tilde{\Omega}}, \theta_2|_{\tilde{\Omega}}, \theta_3|_{\tilde{\Omega}}. \end{aligned}$$

From  $\theta_2|_{\tilde{\Omega}} = 0$ , we obtain the equation

$$-R_2 + \frac{m_2}{m} = 0 \iff m_2 = mR_2$$

and the following conditions on the given function  $R$ :

$$R_{11} + R_1^2 + R - 1 = 0, \quad R_{12} + R_1R_2 = 0.$$

From now on, we assume given  $R$  satisfies above two equations. We restrict again to the 4-dimensional submanifold  $\tilde{\tilde{\Omega}} := \{m_2 = mR_2\} \subset \tilde{\Omega}$  with coframe

$$(69) \quad \{\alpha^1, \alpha^2, \alpha^3, dm\}.$$

It follows

$$\begin{aligned} \theta_3|_{\tilde{\tilde{\Omega}}} &= d(mR_2) - (1 - R)m\alpha^2 - mR_1\alpha^3 \\ &\equiv m(R_{22} + R_2^2 + R - 1)\alpha^2 \quad \text{mod } \theta_1|_{\tilde{\tilde{\Omega}}}, \theta_2|_{\tilde{\tilde{\Omega}}}, \theta_3|_{\tilde{\tilde{\Omega}}}. \end{aligned}$$

Hence, in order to solve the PDEs (50), (51), the function  $R$  must satisfy:

$$(70) \quad R_{22} + R_2^2 + R - 1 = 0.$$

From now on, we assume that the  $R$ -Cartan structure  $\alpha$  satisfies

$$(71) \quad \begin{cases} dR - R_1\alpha^1 - R_2\alpha^2 = 0, \\ dR_1 + (R_1^2 + R - 1)\alpha^1 + R_1R_2\alpha^2 + R_2\alpha^3 = 0, \\ dR_2 + R_1R_2\alpha^1 + (R_2^2 + R - 1)\alpha^2 - R_1\alpha^3 = 0. \end{cases} \iff \begin{cases} R_3 = 0, \\ R_{11} = 1 - R - R_1^2, \\ R_{12} = R_{21} = -R_1R_2, \\ R_{22} = 1 - R - R_2^2. \end{cases}$$

Hence, on the 4-dimensional submanifold  $\tilde{\tilde{\Omega}} := \{m_2 = mR_2\} \subset \tilde{\Omega}$  with the coframe

$$(72) \quad \{\alpha^1, \alpha^2, \alpha^3, dm\}$$

we have  $\theta_2|_{\tilde{\tilde{\Omega}}} = \theta_3|_{\tilde{\tilde{\Omega}}} = 0$  and

$$(73) \quad d\theta_1|_{\tilde{\tilde{\Omega}}} = m(d(\log m - R)).$$

Therefore, the Pfaffian system  $\{\theta_1|_{\tilde{\tilde{\Omega}}} = 0\}$  on  $\tilde{\tilde{\Omega}} := \{m_2 = mR_2\} \subset \tilde{\Omega}$  is a Frobenius system.

From the above computations, we remark that the following relation holds

$$(74) \quad dm - mdR = 0 \iff m = Ce^R,$$

where  $C$  is arbitrary constant, in other words,  $m$  is uniquely determined by the structure function  $R$  of the Cartan structure  $\alpha$ .

We conclude:

**Proposition 5.1.** *For a given  $R$ -Cartan structure which satisfies (71), we have one parameter family of solutions of the PDEs (50), (51).*

Taking into account (74) we can also formulate:

**Proposition 5.2.** *Let us consider an  $R$ -Cartan structure  $\alpha$  on the 3-manifold  $\Sigma$  which satisfies (71). Then, the coframe*

$$(75) \quad \begin{aligned} \omega^1 &= \frac{1}{C} e^{-R} (\alpha^3 - R_2 \alpha^1), \\ \omega^2 &= \frac{1}{C} e^{-R} \alpha^1, \\ \omega^3 &= \alpha^2 \end{aligned}$$

*gives an  $(I, 0, K)$  generalized Finsler structure (i.e., a generalized Landsberg structure) on the 3-manifold  $\Sigma$  with the invariants*

$$(76) \quad I = -2R_2, \quad K = C^2 e^{2R}.$$

*Remark 5.1.* The existence of an  $R$ -Cartan structure which satisfies (71) is guaranteed by the existence of Generalized Landsberg structures (see [12]) or one can prove the existence of the  $R$ -Cartan structures by using the same method (Cartan-Kähler theorem for the frame bundle) as in [12].

## 6. The geometry of the leaf space $\Lambda$

We will study in the present section the geometry of the quotient space  $\Lambda$  that is the leaf space of the codimension two foliation  $\{\alpha^1 = 0, \alpha^2 = 0\}$  on the 3-manifold  $\Sigma$ . This section is based on a private communication with R. Bryant ([6]).

Let us recall from the previous section that we are given an  $R$ -Cartan structure  $\{\alpha^1, \alpha^2, \alpha^3\}$  on  $\Sigma$  whose structure function  $R$  satisfies (71).

**Theorem 6.1.** *The  $R$ -Cartan structure  $(\Sigma, \alpha)$  subject to the conditions (71) induces on the leaf space  $\Lambda = \Sigma / \{\alpha^1 = 0, \alpha^2 = 0\}$  a family of smooth rotationally symmetric Riemannian metrics*

$$(77) \quad g = \eta^2(r) dr^2 + \varphi^2(r) d\theta^2,$$

where

$$(78) \quad \eta(r) = \frac{r e^{-\frac{1}{4}r^2}}{2\Phi(r)}, \quad \varphi(r) = \frac{\Phi(r)}{R_0 - 1},$$

and

$$(79) \quad \Phi(r) = \left[ \left( R_0 - \frac{3}{2} \right) - \left( R_0 - \frac{3}{2} - \frac{1}{4}r^2 \right) e^{-\frac{1}{2}r^2} \right]^{\frac{1}{2}}.$$

Here  $(r, \theta) \in I \times S^1$  are the induced local coordinates on  $\Lambda$ , and  $R_0 > 1$  is a real constant.

*Proof.* It can be seen that from equations (71), one obtains

$$(80) \quad d \left( \left( R_1^2 + R_2^2 + R - \frac{3}{2} \right) e^{2(R-1)} \right) = 0,$$

therefore, it must exist a constant  $C$  such that

$$(81) \quad (R_1^2 + R_2^2 + R - \frac{3}{2})e^{2(R-1)} = C.$$

Observe that the last formula can be rewritten as

$$(82) \quad (R_1^2 + R_2^2)e^{2(R-1)} + f(R) = C,$$

where

$$(83) \quad f(R) = (R - \frac{3}{2})e^{2(R-1)}$$

is a smooth function of one variable. Since  $(R_1^2 + R_2^2)e^{2(R-1)} \geq 0$  it follows that we must have  $f(R) \leq C$ . A simple computation shows that  $f(1) = -\frac{1}{2}$  is the minimum of  $f$ .

Therefore, one gets  $C \geq -\frac{1}{2}$  with equality in the case  $R = 1$ , which is not good for us and it will be therefore eliminated.

For any fixed  $C > -\frac{1}{2}$ , if we denote by  $(f(R_0), R_0)$  the intersection point of the graph of  $f$  with the straight line  $f = C$ , it means that always there exists a unique constant  $R_0 > 1$  such that  $C = f(R_0) = (R_0 - \frac{3}{2})e^{2(R_0-1)}$ , therefore (81) reads now

$$(84) \quad (R_1^2 + R_2^2 + R - \frac{3}{2})e^{2(R-1)} = (R_0 - \frac{3}{2})e^{2(R_0-1)}$$

for any constant  $R_0 > 1$ . Obviously we must have  $R \leq R_0$ .

On the other hand, one can see that

$$(85) \quad d\left[\frac{e^{(R_0-R)}}{R_1^2 + R_2^2}(R_1\alpha^1 + R_2\alpha^2)\right] = d\left[\frac{e^{(R_0-R)}}{R_1^2 + R_2^2}(R_2\alpha^1 - R_1\alpha^2)\right] = 0.$$

Therefore, if we denote

$$(86) \quad dz^1 = \frac{e^{(R_0-R)}}{R_1^2 + R_2^2}dR, \quad dz^2 = \frac{e^{(R_0-R)}}{R_1^2 + R_2^2}dR$$

it follows

$$(87) \quad (dz^1)^2 + (dz^2)^2 = \frac{e^{2(R_0-R)}}{R_1^2 + R_2^2}[(\alpha^1)^2 + (\alpha^2)^2].$$

Therefore, at least locally, one obtains a metric on  $\Lambda$

$$(88) \quad g = (\alpha^1)^2 + (\alpha^2)^2 = \frac{R_1^2 + R_2^2}{e^{2(R_0-R)}}[(dz^1)^2 + (dz^2)^2],$$

by using now (84) we get the equivalent form

$$(89) \quad \begin{aligned} g &= \left((R_0 - \frac{3}{2}) - (R - \frac{3}{2})e^{2(R-R_0)}\right)[(dz^1)^2 + (dz^2)^2] \\ &= \frac{e^{2(R-R_0)}}{(R_0 - \frac{3}{2}) - (R - \frac{3}{2})e^{2(R-R_0)}}dR^2 + \left((R_0 - \frac{3}{2}) - (R - \frac{3}{2})e^{2(R-R_0)}\right)(dz^2)^2. \end{aligned}$$



One could now regard  $(R, z^2)$  as the local coordinates on  $\Lambda$ , but it is better to use the coordinate change  $(r, \theta) \mapsto (R, z^2)$ , where

$$(90) \quad R = R_0 - \frac{r^2}{4}, \quad z^2 = \frac{\theta}{R_0 - 1}.$$

This puts the metric  $g$  in the “polar coordinates” defined on a domain  $I \times S^1$

$$(91) \quad g = \eta^2(r)dr^2 + \varphi^2(r)d\theta^2,$$

where  $I$  is an interval in  $\mathbb{R}$  and

$$(92) \quad \eta(r) = \frac{1}{2} \frac{re^{-\frac{1}{4}r^2}}{\left[ \left( R_0 - \frac{3}{2} \right) - \left( R_0 - \frac{3}{2} - \frac{1}{4}r^2 \right) e^{-\frac{1}{2}r^2} \right]^{\frac{1}{2}}},$$

$$\varphi(r) = \frac{\left[ \left( R_0 - \frac{3}{2} \right) - \left( R_0 - \frac{3}{2} - \frac{1}{4}r^2 \right) e^{-\frac{1}{2}r^2} \right]^{\frac{1}{2}}}{R_0 - 1},$$

or, equivalently,

$$(93) \quad \eta(r) = \frac{re^{-\frac{1}{4}r^2}}{2\Phi(r)}, \quad \varphi(r) = \frac{\Phi(r)}{R_0 - 1},$$

where we put

$$(94) \quad \Phi(r) = \left[ \left( R_0 - \frac{3}{2} \right) - \left( R_0 - \frac{3}{2} - \frac{1}{4}r^2 \right) e^{-\frac{1}{2}r^2} \right]^{\frac{1}{2}}.$$

One can easily see that this is a *smooth rotational metric* even at the origin  $r = 0$  and that the curvature  $R$  reaches its maximum there.  $\square$

*Remark 6.1.* 1. Conversely, for any  $R_0 > 1$ , formula (77) defines a metric on a region of the plane where  $\left( R_0 - \frac{3}{2} \right) - \left( R_0 - \frac{3}{2} - \frac{1}{4}r^2 \right) e^{-\frac{1}{2}r^2} > 0$  that satisfies the desired conditions.

2. The value of the parameter  $R_0 > 1$  influences the geometry of the Riemannian metric  $g$ , indeed

- (1) for  $R_0 \geq \frac{3}{2}$ ,  $(r, \theta) \in \mathbb{R} \times S^1$ , i.e.,  $(r, \theta)$  are global coordinates, in other words  $(\Lambda, g)$  is a smooth Riemannian manifold;
  - (a) when  $R_0 > \frac{3}{2}$ , the metric (77) is incomplete (because the  $r$ -curves have finite length), and can not be extended because

$$\lim_{t \rightarrow \infty} k = -\infty;$$

- (b) when  $R_0 = \frac{3}{2}$ , the metric gets the form

$$(95) \quad g = dr^2 + r^2 e^{-\frac{r^2}{2}} d\theta^2$$

which is a complete Riemannian metric;

- (2) for  $1 < R_0 < \frac{3}{2}$  the condition  $(R_0 - \frac{3}{2}) - (R_0 - \frac{3}{2} - \frac{1}{4}r^2)e^{-\frac{1}{2}r^2} \geq 0$  holds for  $|r| \leq T(R_0)$  for a certain function  $T(R_0) > 0$ .

This metric is not complete at  $t = \pm T(R_0)$ , but the length of the parallels goes to 0 as  $t \rightarrow \pm T(R_0)$ . Thus, the metric “completes” to a singular 2-sphere, with one singularity at the “minimum curvature pole”. This will be an actual singularity if the cone angle is not  $2\pi$ , and this will depend on the value of  $R_0$ .

Let us point out that the prime integral  $m_1$  of the geodesic flow of  $(\Lambda, g)$  constructed in Section 4 reads now  $m_1 = R_1 e^R$  and therefore it gives a first coordinate on the geodesic space  $M = \Sigma / \{\alpha^1=0, \alpha^3=0\}$ . Since the space  $M$  is the leaf space of a codimension two foliation, it should have another coordinate, but to find it explicitly on  $\Sigma$ , we need a second prime integral of the geodesic flow of  $(\Lambda, g)$ , and this is a difficult problem to be studied in the future. Instead of doing this, we will induce local coordinates on  $M$  from the tangent bundle of  $\Lambda$  in the next section.

### 7. The geometry of the surface of revolution $(\Lambda, g)$

Let us consider the complete surface of revolution with the metric (95). We recall that the geodesic flow of a surface of revolution is completely integrable.

In polar coordinates, we have

$$(96) \quad g_{11} = 1, \quad g_{12} = g_{21} = 0, \quad g_{22} = h^2(r) = r^2 e^{-\frac{r^2}{2}}.$$

The sectional curvature of this metric is

$$(97) \quad R = \frac{3}{2} - \frac{1}{4}r^2 \leq \frac{3}{2},$$

and non-vanishing Christoffel symbols

$$(98) \quad \gamma_{22}^1 = -hh' = \frac{r(r^2 - 2)}{2} e^{-\frac{1}{2}r^2}, \quad \gamma_{12}^2 = \frac{h'}{h} = \gamma_{21}^2 = -\frac{1}{2} \frac{r^2 - 2}{r}.$$

It follows that a curve  $\gamma : [a, b] \rightarrow \Lambda$ , with  $\gamma(t) = (r(t), \theta(t))$  is a geodesic on  $\Lambda$  if and only if it satisfies the following geodesic equations:

$$(99) \quad \begin{aligned} r'' + \frac{r(r^2 - 2)}{2} e^{-\frac{1}{2}r^2} (\theta')^2 &= 0, \\ \theta'' - \frac{r^2 - 2}{r} r' \theta' &= 0. \end{aligned}$$

A straightforward computation shows that

$$(100) \quad \frac{d}{dt} [h^2(r)\theta']_{\gamma(t)} = 0,$$

i.e., the function

$$(101) \quad F = h^2(r)\theta' = r^2 e^{-\frac{r^2}{2}} \theta'$$

is constant on the geodesic lines, and thus provides a prime integral of the geodesic flow of  $(\Lambda, g)$ . This is called the *Clairaut integral* (see for example [14]).

A natural question to ask is if our prime integral  $m_1$  is different from the usual Clairaut integral. With the notations above, we have:

**Proposition 7.1.** *The prime integral  $R_1 e^R$  coincides with the Clairaut integral  $F$  up to multiplication with a negative constant.*

*Proof.* Let us denote by  $f_1, f_2$  the  $g$ -orthonormal frame on  $\Lambda$ , then the relation with the natural frame  $\frac{\partial}{\partial r}, \frac{\partial}{\partial \theta}$  can be written as

$$(102) \quad \begin{aligned} f_1 &= \frac{1}{\sqrt{|g|}} \left( \frac{\partial \sqrt{E}}{\partial \zeta^2} \frac{\partial}{\partial r} - \frac{\partial \sqrt{E}}{\partial \zeta^1} \frac{\partial}{\partial \theta} \right) =: m^1 \frac{\partial}{\partial r} + m^2 \frac{\partial}{\partial \theta}, \\ f_2 &= \frac{1}{\sqrt{E}} \left( \zeta^1 \frac{\partial}{\partial r} + \zeta^2 \frac{\partial}{\partial \theta} \right) =: l^1 \frac{\partial}{\partial r} + l^2 \frac{\partial}{\partial \theta}, \end{aligned}$$

where  $E = (\zeta^1)^2 + h^2(r)(\zeta^2)^2$  is the energy function of the Riemannian structure  $(\Lambda, g)$  and  $|g| = h^2$  is the determinant of the matrix  $(g_{ij})$ . Here we use  $(r, \theta, \zeta^1, \zeta^2) \in T\Lambda$  to denote the natural coordinates on the tangent space of  $(\Lambda, g)$ .

Their horizontal lifts to the unit sphere bundle  $\Sigma$  of  $(\Lambda, g)$  are

$$(103) \quad \hat{f}_1 = m^1 \frac{\delta}{\delta r} + m^2 \frac{\delta}{\delta \theta}, \quad \hat{f}_2 = l^1 \frac{\delta}{\delta r} + l^2 \frac{\delta}{\delta \theta},$$

where  $\frac{\delta}{\delta r}, \frac{\delta}{\delta \theta}$  is the adapted basis of the horizontal subspace of  $T\Lambda$  with respect to the Levi-Civita connection, i.e.,

$$(104) \quad \begin{aligned} \frac{\delta}{\delta r} &= \frac{\partial}{\partial r} - \gamma_{12}^2 \zeta^2 \frac{\partial}{\partial \zeta^2} = \frac{\partial}{\partial r} - \frac{h'}{h} \zeta^2 \frac{\partial}{\partial \zeta^2}, \\ \frac{\delta}{\delta \theta} &= \frac{\partial}{\partial \theta} - \gamma_{22}^1 \zeta^2 \frac{\partial}{\partial \zeta^1} - \gamma_{21}^2 \zeta^1 \frac{\partial}{\partial \zeta^2} = \frac{\partial}{\partial \theta} + h h' \zeta^2 \frac{\partial}{\partial \zeta^1} - \frac{h'}{h} \zeta^1 \frac{\partial}{\partial \zeta^2}. \end{aligned}$$

With these notations, we have

$$(105) \quad R_1 = m^1 \frac{\delta R}{\delta r} + m^2 \frac{\delta R}{\delta \theta} = m^1 \frac{\partial R}{\partial r}.$$

We will compute now  $m^1$  as follows

$$(106) \quad m^1 = \frac{1}{\sqrt{|g|}} \frac{\partial \sqrt{E}}{\partial \zeta^2} = \frac{1}{2E \sqrt{|g|}} \frac{\partial E}{\partial \zeta^2} = h(r) \frac{\zeta^2}{E},$$

and taking into account of (97) it follows that on  $\Sigma$  we have

$$(107) \quad R_1 = -r h(r) \zeta^2.$$

On the other hand, let us recall that the canonical lift of a geodesic  $\gamma$  of  $(\Lambda, g)$  to the unit sphere bundle  $\Sigma = \mathcal{F}(\Lambda, g)$  gives the geodesic flow of  $(\Lambda, g)$

on  $\Sigma$ . Its tangent vector is

$$(108) \quad \hat{f}_2 = \hat{T}(t) = \dot{\gamma}^i(t) \frac{\delta}{\delta x^i} = r' \frac{\delta}{\delta r} + \theta' \frac{\delta}{\delta \theta},$$

and therefore, along the lines of the geodesic flow we have  $\theta' = d\theta(\hat{f}_2) = l^2 \equiv \zeta^2$  on  $\Sigma$ .

Finally, observing that

$$(109) \quad e^R = e^{\frac{3}{2} - \frac{1}{4}r^2} = e^{\frac{3}{2}} \frac{h}{r}$$

it follows

$$(110) \quad R_1 e^R = -e^{\frac{3}{2}} h^2 \zeta^2,$$

i.e., on the geodesic flow of  $(\Lambda, g)$  we have  $R_1 e^R = -e^{\frac{3}{2}} F$ , and the proposition is proved.  $\square$

Finally, let us remark that we can regard the space  $M$  as a submanifold of the tangent space  $T\Lambda$ . On  $T\Lambda$  we have already two functionally independent prime integrals of the geodesic flow of  $(\Lambda, g)$ , namely, the Clairaut integral  $F = -e^{-\frac{3}{2}} R_1 e^R$ , regarded now as function on  $T\Lambda$ , and the energy function  $E$  of the Riemannian structure  $(\Lambda, g)$ .

Therefore, we obtain in this way induced local coordinates  $(F, E)$  on  $M$  from  $T\Lambda$ , i.e.,  $\pi_M : T\Lambda \rightarrow M$ ,  $(r, \theta, \zeta^1, \zeta^2) \mapsto (F, E)$ . Pay attention to the fact that  $M$  is not a smooth manifold because  $(\Lambda, g)$  cannot have all geodesics closed and of same length.

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