

## NOTE ABOUT THE ISOCHRONICITY OF HAMILTONIAN SYSTEMS AND THE CURVATURE OF THE ENERGY

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ABSTRACT. In this note we show the relationship between the period of a isochronous center of a planar Hamiltonian system and the Gauss curvature of the surface  $S = (x, y, H(x, y))$  where  $H$  is the energy function of the system.

### 1. Introduction

It is well known that the planar analytic Hamiltonian systems are the planar differential systems of the form

$$(1.1) \quad \begin{cases} \dot{x} = -\frac{\partial H}{\partial y}(x, y) \\ \dot{y} = \frac{\partial H}{\partial x}(x, y), \end{cases}$$

where  $H$  is an analytic function on  $\mathbb{R}^2$ . The solutions of these systems are contained in the level curves  $\{H(x, y) = h, h \in \mathbb{R}\}$ . A point  $p$  is called center if it has a neighborhood formed by periodic orbits. The largest neighborhood of  $p$  which is entirely covered by periodic orbits is called the *period annulus* of  $p$  and we will denote it by  $\mathcal{P}$ . The function which associates to any periodic orbit  $\gamma$  in  $\mathcal{P}$  its period is called the *period function*. The center is called *isochronous center* when the period function is a constant  $\omega$ . It is well known that only nondegenerate centers can be isochronous. And from now on we will assume that  $H(0, 0) = 0$  and the system (1.1) has a nondegenerate center at the origin.

Cima, Mañosas and Villadelprat in [1] study the isochronous centers of Hamiltonians systems of the form

$$(1.2) \quad H(x, y) = A(x) + B(x)y + C(x)y^2$$

with  $A, B, C$  analytical functions. They prove that if the center is isochronous of period  $\omega$  then

$$(1.3) \quad \frac{d^2 (4AC - B^2)}{dx^2}(0) = \frac{8\pi^2}{\omega^2}.$$

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Clearly, this allows to calculate  $\omega$  in an analytic way and without carrying out any integration.

However, in fact, this is a particular case of the following general result.

## 2. Result

**Theorem 1.** *Let*

$$S \equiv \vec{x} : U \subset \mathbb{R}^2 \rightarrow \mathbb{R}^3$$

$$(x, y) \rightarrow \vec{x}(x, y) = (x, y, H(x, y))$$

*be the surface generated by the energy function  $H$ . Let  $K = K_{(0,0,0)}(S)$  be the Gauss curvature of  $S$  at the point  $(0, 0, 0)$ . Then:*

*If the center is isochronous of period  $\omega$  we have*

$$K\omega^2 = 4\pi^2.$$

*Proof.* We know that exist two analytical functions  $f, g$  such that

$$H = \frac{f^2 + g^2}{2},$$

with  $f(0, 0) = g(0, 0) = 0$  and

$$(2.1) \quad \det \begin{pmatrix} \frac{\partial f}{\partial x} & \frac{\partial f}{\partial y} \\ \frac{\partial g}{\partial x} & \frac{\partial g}{\partial y} \end{pmatrix} = \frac{2\pi}{\omega}$$

in every point of a neighborhood  $U$  of the center point  $(0, 0)$ , (see [2], [3]). We consider the surface:

$$S \equiv \vec{x} : U \subset \mathbb{R}^2 \rightarrow \mathbb{R}^3$$

$$(x, y) \rightarrow \vec{x}(x, y) = (x, y, H(x, y)) = (x, y, \frac{f^2(x, y) + g^2(x, y)}{2}).$$

Then, the Gauss application is

$$N : S \rightarrow S^2$$

$$\vec{x}(x, y) \rightarrow N(x, y) = \frac{\frac{\partial \vec{x}}{\partial x} \times \frac{\partial \vec{x}}{\partial y}}{\left\| \frac{\partial \vec{x}}{\partial x} \times \frac{\partial \vec{x}}{\partial y} \right\|} = (N_1(x, y), N_2(x, y), N_3(x, y))$$

with

$$N_1(x, y) = -\frac{f \frac{\partial f}{\partial x} + g \frac{\partial g}{\partial x}}{\sqrt{(f \frac{\partial f}{\partial x} + g \frac{\partial g}{\partial x})^2 + (f \frac{\partial f}{\partial y} + g \frac{\partial g}{\partial y})^2 + 1}}$$

$$N_2(x, y) = -\frac{f \frac{\partial f}{\partial y} + g \frac{\partial g}{\partial y}}{\sqrt{(f \frac{\partial f}{\partial x} + g \frac{\partial g}{\partial x})^2 + (f \frac{\partial f}{\partial y} + g \frac{\partial g}{\partial y})^2 + 1}}$$

$$N_3(x, y) = \frac{1}{\sqrt{(f \frac{\partial f}{\partial x} + g \frac{\partial g}{\partial x})^2 + (f \frac{\partial f}{\partial y} + g \frac{\partial g}{\partial y})^2 + 1}}.$$

We calculate and obtain that the Weingarten endomorphism

$$dN_{(0,0,0)}(S) : T_{(0,0,0)}(S) \rightarrow T_{(0,0,0)}(S)$$

at the point  $(0, 0, 0)$  in the basis  $\left\{ \frac{\partial \vec{x}}{\partial x} = \vec{x}_1, \frac{\partial \vec{x}}{\partial y} = \vec{x}_2 \right\}$  has the associated matrix

$$dN_{(0,0,0)}(S) |_{\vec{x}_i} = \begin{pmatrix} -\left(\frac{\partial f}{\partial x}\right)^2 - \left(\frac{\partial g}{\partial x}\right)^2 & -\frac{\partial f}{\partial x} \frac{\partial f}{\partial y} - \frac{\partial g}{\partial x} \frac{\partial g}{\partial y} \\ -\frac{\partial f}{\partial x} \frac{\partial f}{\partial y} - \frac{\partial g}{\partial x} \frac{\partial g}{\partial y} & -\left(\frac{\partial f}{\partial y}\right)^2 - \left(\frac{\partial g}{\partial y}\right)^2 \end{pmatrix} |_{(0,0)}.$$

We also calculate and obtain:

$$K = K_{(0,0,0)}(S) = \det(dN_{(0,0,0)}(S) |_{\vec{x}_i}) = \left( \frac{\partial f}{\partial x} \frac{\partial g}{\partial y} - \frac{\partial f}{\partial y} \frac{\partial g}{\partial x} \right)^2 |_{(0,0)}.$$

Put in other way

$$K = \left( \det \begin{pmatrix} \frac{\partial f}{\partial x} & \frac{\partial f}{\partial y} \\ \frac{\partial g}{\partial x} & \frac{\partial g}{\partial y} \end{pmatrix} |_{(0,0)} \right)^2,$$

and, using (2.1), this implies

$$K\omega^2 = 4\pi^2.$$

This complete the proof.  $\square$

### 3. Consequence

We have said in the introduction that (1.3) is a particular case because: if the system has the form (1.2) then

$$\begin{aligned} \frac{\partial \vec{x}}{\partial x} &= \vec{x}_1 = (1, 0, \frac{dA}{dx}(x) + \frac{dB}{dx}(x)y + \frac{dC}{dx}(x)y^2) \\ \frac{\partial \vec{x}}{\partial y} &= \vec{x}_2 = (0, 1, B(x) + 2C(x)y). \end{aligned}$$

We calculate and obtain:

$$\begin{aligned} dN_{(0,0,0)}(S) : T_{(0,0,0)}(S) &\rightarrow T_{(0,0,0)}(S) \\ \vec{x}_1 |_{(0,0)} \rightarrow dN_{(0,0,0)}(\vec{x}_1 |_{(0,0)}) &= -\frac{d^2 A}{dx^2}(0)\vec{x}_1 |_{(0,0)} - \frac{dB}{dx}(0)\vec{x}_2 |_{(0,0)} \\ \vec{x}_2 |_{(0,0)} \rightarrow dN_{(0,0,0)}(\vec{x}_2 |_{(0,0)}) &= -\frac{dB}{dx}(0)\vec{x}_1 |_{(0,0)} - 2C(0)\vec{x}_2 |_{(0,0)}. \end{aligned}$$

and using the Theorem 1 we have

$$K = 2C(0) \frac{d^2 A}{dx^2}(0) - \frac{dB}{dx}(0)^2 = \frac{4\pi^2}{\omega^2}.$$

But if the center is isochronous then  $\frac{dA}{dx}(0) = A(0) = B(0) = 0$  (see [1]); and in consequence

$$\frac{d^2 (4AC - B^2)}{dx^2}(0) = 4C(0) \frac{d^2 A}{dx^2}(0) - 2 \frac{dB}{dx}(0)^2 = \frac{8\pi^2}{\omega^2}.$$

### References

- [1] A. Cima, F. Mañosas, and J. Villadelprat, *Isochronicity for Several Classes of Hamiltonian Systems*, J. Differential Equations **157** (1999), no. 2, 373–413.
- [2] H. Ito, *Convergence of Birkhoff Normal Forms for Integrable Systems*, Comment. Math. Helv. **64** (1989), no. 3, 412–461.
- [3] J. Villadelprat, *Index of Vector Fields on Manifolds and Isochronicity for Planar Hamiltonian Differential Systems*, Ph. D. Thesis, Pub. Univ. Autònoma de Barcelona, 1999.

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