

SUPERCONVERGENCE OF FINITE ELEMENT METHODS FOR LINEAR QUASI-PARABOLIC INTEGRO-DIFFERENTIAL EQUATIONS

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ABSTRACT. We consider finite element methods applied to a class of quasi parabolic integro-differential equations in R^d . Global strong superconvergence, which only requires that partitions are quasi-uniform, is investigated for the error between the approximate solution and the Sobolev-Volterra projection of the exact solution. Two order superconvergence results are demonstrated in $W^{1,p}(\Omega)$ and $L_p(\Omega)$, for $2 \leq p < \infty$.

1. Introduction

Assume that Ω is a bounded domain in R^d with piecewise smooth boundary $\partial\Omega$. Consider the following initial boundary value problem of linear quasi-parabolic integro-differential equation

$$\left\{ \begin{array}{ll} u_t = \nabla \cdot \{a(t)\nabla u_t + b(t)\nabla u + \int_0^t c(t,\tau)\nabla u(\tau)d\tau\} + f(x,t), & \text{in } \Omega \times (0, T], \\ u(x,t) = 0, & \text{on } \partial\Omega \times J, J = [0, T], \\ u(x,0) = u_0(x), & \text{in } \Omega. \end{array} \right. \quad (1.1)$$

Where $a(t) = a(x,t)$, $b(t) = b(x,t)$, $c(t,\tau) = c(t,\tau,x)$, f , u_0 are known functions and bounded together with their derivatives up to certain orders as far as the ensuring analysis requires, and $a(x,t)$ satisfies

$$0 < a_0 \leq a(x,t) \leq a_1, (x,t) \in \Omega \times [0, T], \quad (1.2)$$

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for positive constants a_0, a_1 . The existence and uniqueness of the solution of (1.1) have been considered and analyzed by Cui[1].

Let $L_p(\Omega)$ and $W^{m,p}(\Omega)$, $H^m(\Omega) = W^{m,2}(\Omega)$, for any integer $m \geq 0$ and $1 \leq p \leq \infty$, denote the usual Lebesgue and Sobolev spaces respectively. L_2 and L_p norms are denoted by $\|\cdot\|$ and $\|\cdot\|_{0,p}$, Sobolev norms by $\|\cdot\|_{m,p}$ and $\|\cdot\|_m$. For any integer $s \geq 0$ and $t \in J$, we define

$$\|u(t)\|_{s,m,p} = \sum_{j=0}^s \|D_t^j u(t)\|_{m,p} + \int_0^t \sum_{j=0}^s \|D_t^j u(\tau)\|_{m,p} d\tau. \quad (1.3)$$

Let (\cdot, \cdot) denote the inner product in $L_2(\Omega)$ or $L_2(\Omega)^2$. In addition, we also use the notation $p' = \frac{p}{p-1}$ to denote the conjugate index of p for $2 \leq p < \infty$. We shall denote by C a constant independent of h , not necessarily the same at different occurrences.

The mathematical methods applied to (1.1) have been considered and analyzed by several authors [1], [2]. In Cui [1], optimal order error estimates are obtained in L_2 . Because in the right side of (1.1), there are not only ∇u_t but also ∇u and the integration of t for ∇u , it seems very complex. If we employ the traditional projection of finite element, such as Ritz projection, Sobolev projection and Ritz-Sobolev projection, we can't reflect its eccentric characters. So it's very difficult to get the superconvergence estimate. The paper develops a new projection of finite element (called Sobolev-Volterra projection) and by employing a special method for initial value selection, to study superconvergence of the error between the approximation solution U and the Sobolev-Volterra projection $V_h u$ of the exact solution u of (1.1). Two order superconvergence results in L_p and $W^{1,p}$ for $2 \leq p < \infty$ are demonstrated for the general quasi-uniform partition. The superconvergence estimates for this class of problem have been studied by [3]-[8].

The rest of the paper is organized as follows. In section 2 we lay out the Galerkin approximation formula of the problem, and give the Sobolev-Volterra projection and some necessary preparations. Several Lemmas that are needed for the main results are proved in Section 3. Section 4 derives the main conclusions of the paper, two order superconvergence in $L_p(\Omega)$ and $W^{1,p}(\Omega)$ is given in Theorems 4.1 and 4.3, respectively.

2. Finite-Element Formulations And Sobolev-Volterra Projection

In this section we will consider the semidiscrete Galerkin method with respect to space for (1.1) and introduce the Sobolev-Volterra projection.

Let $\{S_h\}_{0 < h \leq 1}$ be a family of finite-dimensional subspaces of $H_0^1(\Omega)$ satisfying the following approximation property: for some $r \geq 2$, $1 \leq s \leq r$, $1 \leq p \leq \infty$,

$$\inf_{\chi \in S_h} \{\|v - \chi\|_{0,p} + h\|v - \chi\|_{1,p}\} \leq Ch^s \|v\|_{s,p}, \quad v \in W^{s,p}(\Omega) \cap H_0^1(\Omega). \quad (2.1)$$

We also assume that the standard inverse properties holds in subspaces S_h .

We first introduce the Ritz projection operator $R_h = R_h(t) : H_0^1(\Omega) \rightarrow S_h$, for $t \in J$, defined by

$$(a(t)\nabla(R_h w - w), \nabla\chi) = 0, \quad \chi \in S_h. \quad (2.2)$$

The weak form of (1.1) is to find $u \in H_0^1(\Omega)$, $t \in J$, such that

$$\begin{cases} (a) (u_t, v) + (a(t)\nabla u_t + b(t)\nabla u + \int_0^t c(t, \tau)\nabla u(\tau)d\tau, \nabla v) = (f(t), v), & v \in H_0^1(\Omega), \\ (b) u(x, 0) = u_0(x), & x \in \Omega. \end{cases} \quad (2.3)$$

We introduce the Sobolev-Volterra projection operator $V_h = V_h(t) : H_0^1(\Omega) \rightarrow S_h$, for $t \in J$, satisfies

$$(a(t)\nabla(V_h u_t - u_t) + b(t)\nabla(V_h u - u) + \int_0^t c(t, \tau)\nabla(V_h u - u)d\tau, \nabla\chi) = 0, \quad \chi \in S_h. \quad (2.4)$$

Now we define the semidiscrete finite element approximation to the solution u of (1.1) is to find a map $U(t) : J \rightarrow S_h$, such that

$$\begin{cases} (a) (U_t, \chi) + (a(t)\nabla U_t + b(t)\nabla U + \int_0^t c(t, \tau)\nabla U(\tau)d\tau, \nabla\chi) = (f(t), \chi), & \chi \in S_h, t \in J \\ (b) U(0) = V_h u(0), & x \in \Omega. \end{cases} \quad (2.5)$$

Noting that Sobolev-Volterra projection is Ritz-Sobolev projection when $t = 0$, Ritz-Sobolev projection has been studied in Li [8].

Let $\eta = V_h u - u$, then (2.4) becomes

$$(a(t)\nabla\eta_t + b(t)\nabla\eta + \int_0^t c(t, \tau)\nabla\eta(\tau)d\tau, \nabla\chi) = 0, \quad \chi \in S_h. \quad (2.6)$$

3. LEMMAS

It is easy to prove the problem (2.6) is equivalent to the following equation

$$(a(t)\nabla\eta, \nabla\chi) + \left(\int_0^t [(b(\tau)\nabla\eta(\tau) - a_t(\tau)\nabla\eta(\tau) + \int_0^\tau c(\tau, s)\nabla\eta(s)ds]d\tau, \nabla\chi \right) = 0, \quad \chi \in S_h \quad (3.1)$$

From [9], We derive the following lemma.

Lemma 3.1. Let $2 \leq p < \infty$, then there exists a constant C such that for $V_h u$ defined by (2.4)

$$\|V_h u - u\|_{0,p} + h\|V_h u - u\|_{1,p} \leq Ch^r \|u\|_{0,r,p}. \quad (3.2)$$

Proof. We first estimate $\|\eta\|_{1,p}$. Let η_x be any component of $\nabla\eta$, then

$$\|\eta_x\|_{0,p} = \sup\{ |(\eta_x, \varphi)|, \varphi \in L_{p'}(\Omega), \|\varphi\|_{0,p'} = 1 \}.$$

Introducing the dual auxiliary problem, for any $\varphi \in L_{p'}(\Omega)$, $\|\varphi\|_{0,p'} = 1$, let $\psi \in W^{1,p'}(\Omega)$, be the solution of

$$(a(t)\nabla\psi, \nabla v) = -(v, \varphi_x), v \in H_0^1(\Omega), t \in J, \quad (3.3)$$

Thus we have

$$\|\psi\|_{1,p'} \leq C\|\varphi\|_{0,p'} \leq C, t \in J. \quad (3.4)$$

Then by Green's formula, (3.3)

$$\begin{aligned} (\eta_x, \varphi) &= -(\eta, \varphi_x) = (a(t)\nabla\eta, \nabla\psi) \\ &= (a(t)\nabla\eta, \nabla(\psi - R_h\psi)) + (a(t)\nabla\eta, \nabla R_h\psi) \\ &= (a(t)\nabla(V_h u - R_h u + R_h u - u), \nabla(\psi - R_h\psi)) + (a(t)\nabla\eta, \nabla R_h\psi) \\ &= (a(t)\nabla(R_h u - u), \nabla(\psi - R_h\psi)) + (a(t)\nabla\eta, \nabla R_h\psi) \\ &= I_1 + I_2. \end{aligned}$$

By Hölder inequality, (3.4) and (3.1)

$$\begin{aligned} |I_1| &\leq C\|R_h u - u\|_{1,p} \|\psi - R_h\psi\|_{1,p'} \\ &\leq Ch^{r-1}\|u\|_{r,p} \|\psi\|_{1,p'} \\ &\leq Ch^{r-1}\|u\|_{r,p}. \end{aligned}$$

And

$$\begin{aligned} |I_2| &= \left| \left(\int_0^t [(a_t(\tau) - b(\tau))\nabla\eta(\tau) - \int_0^\tau c(\tau, s)\nabla\eta(s)ds]d\tau, \nabla R_h\psi \right) \right| \\ &\leq C \int_0^t [\|\eta\|_{1,p} + \int_0^\tau \|\eta\|_{1,p}ds]d\tau \|R_h\psi\|_{1,p'} \\ &\leq C \int_0^t \|\eta\|_{1,p}d\tau \|\psi\|_{1,p'} \\ &\leq C \int_0^t \|\eta\|_{1,p}d\tau. \end{aligned}$$

Then

$$|(\eta_x, \varphi)| \leq |I_1| + |I_2| \leq Ch^{r-1}\|u\|_{r,p} + C \int_0^t \|\eta\|_{1,p}d\tau.$$

Thus

$$\|\eta_x\|_{0,p} = \sup_{\substack{\varphi \in L_{p'}(\Omega) \\ \|\varphi\|_{0,p'}=1}} |(\eta_x, \varphi)| \leq Ch^{r-1}\|u\|_{r,p} + C \int_0^t \|\eta\|_{1,p}d\tau,$$

We have

$$\|\eta\|_{1,p} \leq Ch^{r-1}\|u\|_{r,p} + C \int_0^t \|\eta\|_{1,p} d\tau,$$

By Gronwall's Lemma, we have

$$\|\eta\|_{1,p} \leq Ch^{r-1}\|u\|_{r,p}. \quad (3.5)$$

Next we turn to estimate $\|\eta\|_{0,p}$. Introducing the other dual auxiliary problem, for any $\phi \in L_{p'}(\Omega)$, $\|\phi\|_{0,p'} = 1$, let $\Phi \in W^{2,p'}(\Omega)$ be the solution of

$$(a(t)\nabla\Phi, \nabla v) = (\phi, v), \quad v \in H_0^1(\Omega), \quad (3.6)$$

Thus we have

$$\|\Phi\|_{2,p'} \leq C\|\phi\|_{0,p'} \leq C. \quad (3.7)$$

Then $(\eta, \phi) = (a(t)\nabla\eta, \nabla\Phi) = (a(t)\nabla\eta, \nabla(\Phi - R_h\Phi)) + (a(t)\nabla\eta, \nabla R_h\Phi) = I_1 + I_2$.
By Hölder inequality, (3.5) and (3.7)

$$|I_1| \leq C\|\eta\|_{1,p} \|\Phi - R_h\Phi\|_{1,p'} \leq Ch^r\|u\|_{r,p} \|\Phi\|_{2,p'} \leq Ch^r\|u\|_{r,p}.$$

Hence by using (3.1),

$$\begin{aligned} |I_2| &= |(a(t)\nabla\eta, \nabla R_h\Phi)| \\ &= \left| \left(\int_0^t [(b(\tau) - a_t(\tau))\nabla\eta(\tau) + \int_0^\tau c(\tau, s)\nabla\eta(s)ds]d\tau, \nabla R_h\Phi \right) \right| \\ &\leq \left| \left(\int_0^t [(b(\tau) - a_t(\tau))\nabla\eta(\tau) + \int_0^\tau c(\tau, s)\nabla\eta(s)ds]d\tau, \nabla(R_h\Phi - \Phi) \right) \right| \\ &\quad + \left| \left(\int_0^t [(b(\tau) - a_t(\tau))\nabla\eta(\tau) + \int_0^\tau c(\tau, s)\nabla\eta(s)ds]d\tau, \nabla\Phi \right) \right| \\ &= J_1 + J_2. \end{aligned}$$

Following from Hölder inequality and (3.5)

$$\begin{aligned} J_1 &\leq C \int_0^t \|\eta\|_{1,p} d\tau \|R_h\Phi - \Phi\|_{1,p'} \\ &\leq Ch^r \int_0^t \|u\|_{r,p} d\tau \|\Phi\|_{2,p'} \\ &\leq Ch^r \int_0^t \|u\|_{r,p} d\tau. \end{aligned}$$

By using integration by parts, we have

$$J_2 \leq C \int_0^t \|\eta\|_{0,p} d\tau \|\Phi\|_{2,p'} \leq C \int_0^t \|\eta\|_{0,p} d\tau.$$

Thus

$$|(\eta, \phi)| \leq Ch^r(\|u\|_{r,p} + \int_0^t \|u\|_{r,p} d\tau) + C \int_0^t \|\eta\|_{0,p} d\tau.$$

It follows from Gronwall's Lemma

$$\|\eta\|_{0,p} = \sup_{\substack{\phi \in L_{p'}(\Omega) \\ \|\phi\|_{0,p'}=1}} |(\eta, \phi)| \leq Ch^r (\|u\|_{r,p} + \int_0^t \|u\|_{r,p} d\tau)$$

The proof is completed.

We have the following estimate for Sobolev-Volterra projection.

Lemma3.2. Let $0 \leq l \leq r - 2$, $p > 1$, $k = 0, 1, 2$, and $t \in J$, we have

$$|(D_t^k \eta, \phi)| \leq Ch^{r+l} \|u(t)\|_{k,r,p} \|\phi\|_{l,p'}, \quad \phi \in W^{l,p'}(\Omega). \quad (3.8)$$

Proof. We only prove Lemma3.2 when $k = 0$, i.e.

$$|(\eta, \phi)| \leq Ch^{r+l} \|u(t)\|_{0,r,p} \|\phi\|_{l,p'}, \quad \phi \in W^{l,p'}(\Omega). \quad (3.9)$$

Introducing the dual auxiliary problem, for any $\phi \in W^{l,p'}(\Omega)$, let $\Phi \in W^{l+2,p'}(\Omega)$ be the solution of

$$(a(t)\nabla\Phi, \nabla v) = (v, \phi), \quad v \in H_0^1(\Omega), \quad (3.10)$$

Then we have

$$\|\Phi\|_{l+2,p'} \leq C\|\phi\|_{l,p'}. \quad (3.11)$$

Let Φ^h be the Galerkin approximation of Φ , then by (3.10), we have

$$\begin{aligned} (\eta, \phi) &= (a(t)\nabla\eta, \nabla\Phi) = (a(t)\nabla\eta, \nabla(\Phi - \Phi^h)) + (a(t)\nabla\eta, \nabla\Phi^h) \\ &= (a(t)\nabla(V_h u - R_h u + R_h u - u), \nabla(\Phi - \Phi^h)) + (a(t)\nabla\eta, \nabla\Phi^h) \\ &= (a(t)\nabla(R_h u - u), \nabla(\Phi - \Phi^h)) + (a(t)\nabla\eta, \nabla\Phi^h) \\ &= I_1 + I_2. \end{aligned}$$

Then by (1.2) and Hölder inequality

$$|I_1| \leq C\|R_h u - u\|_{1,p} \|\Phi - \Phi^h\|_{1,p'} \leq Ch^{r+l} \|u\|_{r,p} \|\Phi\|_{l+2,p'},$$

And by (3.1)

$$\begin{aligned} |I_2| &= |(a(t)\nabla\eta, \nabla\Phi^h)| \\ &= |(\int_0^t [(b(\tau)\nabla\eta(\tau) - a_t(\tau)\nabla\eta(\tau) + \int_0^\tau c(\tau,s)\nabla\eta(s)ds]d\tau, \nabla\Phi^h)| \\ &= |(\int_0^t [(b(\tau)\nabla\eta(\tau) - a_t(\tau)\nabla\eta(\tau) + \int_0^\tau c(\tau,s)\nabla\eta(s)ds]d\tau, \nabla(\Phi^h - \Phi))| \\ &\quad + |(\int_0^t [(b(\tau)\nabla\eta(\tau) - a_t(\tau)\nabla\eta(\tau) + \int_0^\tau c(\tau,s)\nabla\eta(s)ds]d\tau, \nabla\Phi)| \\ &= J_1 + J_2. \end{aligned}$$

By Hölder inequality and (3.2)

$$\begin{aligned}
J_1 &\leq C \left\{ \int_0^t \|\eta\|_{1,p} + \int_0^\tau \|\eta\|_{1,p} ds \right\} \|\Phi^h - \Phi\|_{1,p'} \\
&\leq C \int_0^t \|\eta\|_{1,p} d\tau \|\Phi^h - \Phi\|_{1,p'} \\
&\leq Ch^{r+l} \int_0^t \|u\|_{0,r,p} d\tau \|\Phi\|_{l+2,p'}.
\end{aligned}$$

Following from Green's formula

$$\begin{aligned}
J_2 &\leq \left| \int_0^t ((b(\tau) - a_t(\tau)) \nabla \eta(\tau), \nabla \Phi) d\tau \right| + \left| \int_0^t \int_0^\tau (c(\tau, s) \nabla \eta(s), \nabla \Phi) ds d\tau \right| \\
&= \left| \int_0^t (\eta, \nabla \cdot ((b(\tau) - a_t(\tau)) \nabla \Phi)) d\tau \right| + \left| \int_0^t \int_0^\tau (\eta(s), \nabla \cdot (c(\tau, s) \nabla \Phi)) ds d\tau \right| \\
&= \left| \int_0^t \frac{(\eta, \nabla \cdot ((b(\tau) - a_t(\tau)) \nabla \Phi))}{\|\nabla \cdot ((b(\tau) - a_t(\tau)) \nabla \Phi)\|_{l,p'}} \|\nabla \cdot ((b(\tau) - a_t(\tau)) \nabla \Phi)\|_{l,p'} d\tau \right| \\
&\quad + \left| \int_0^t \int_0^\tau \frac{(\eta(s), \nabla \cdot (c(\tau, s) \nabla \Phi))}{\|\nabla \cdot (c(\tau, s) \nabla \Phi)\|_{l,p'}} \|\nabla \cdot (c(\tau, s) \nabla \Phi)\|_{l,p'} ds d\tau \right| \\
&\leq C \int_0^t \sup_{\psi \in W^{l,p'}(\Omega)} \frac{|(\eta, \psi)|}{\|\psi\|_{l,p'}} \|\Phi\|_{l+2,p'} d\tau + C \int_0^t \sup_{\psi \in W^{l,p'}(\Omega)} \frac{|(\eta, \psi)|}{\|\psi\|_{l,p'}} \|\Phi\|_{l+2,p'} d\tau \\
&\leq C \int_0^t \sup_{\psi \in W^{l,p'}(\Omega)} \frac{|(\eta, \psi)|}{\|\psi\|_{l,p'}} d\tau \|\Phi\|_{l+2,p'}.
\end{aligned}$$

Therefore, by (3.11), combing I_1 , J_1 , J_2 ,

$$|(\eta, \phi)| \leq Ch^{r+l} [\|u\|_{r,p} + \int_0^t \|u\|_{0,r,p} d\tau] \|\phi\|_{l,p'} + C \int_0^t \sup_{\psi \in W^{l,p'}(\Omega)} \frac{|(\eta, \psi)|}{\|\psi\|_{l,p'}} d\tau \|\phi\|_{l,p'}$$

Then

$$\begin{aligned}
\frac{|(\eta, \phi)|}{\|\phi\|_{l,p'}} &\leq Ch^{r+l} \|u\|_{0,r,p} + C \int_0^t \sup_{\psi \in W^{l,p'}(\Omega)} \frac{|(\eta, \psi)|}{\|\psi\|_{l,p'}} d\tau \\
\sup_{\psi \in W^{l,p'}(\Omega)} \frac{|(\eta, \psi)|}{\|\psi\|_{l,p'}} &\leq Ch^{r+l} \|u\|_{0,r,p} + C \int_0^t \sup_{\psi \in W^{l,p'}(\Omega)} \frac{|(\eta, \psi)|}{\|\psi\|_{l,p'}} d\tau
\end{aligned}$$

By Gronwall's Lemma, we have

$$\sup_{\psi \in W^{l,p'}(\Omega)} \frac{|(\eta, \psi)|}{\|\psi\|_{l,p'}} \leq Ch^{r+l} \|u\|_{0,r,p}$$

Hence

$$\frac{|(\eta, \phi)|}{\|\phi\|_{l,p'}} \leq \sup_{\psi \in W^{l,p'}(\Omega)} \frac{|(\eta, \psi)|}{\|\psi\|_{l,p'}} \leq Ch^{r+l} \|u\|_{0,r,p}$$

Then

$$|(\eta, \phi)| \leq \sup_{\psi \in W^{l,p'}(\Omega)} \frac{|(\eta, \psi)|}{\|\psi\|_{l,p'}} \|\phi\|_{l,p'} \leq Ch^{r+l} \|u\|_{0,r,p} \|\phi\|_{l,p'}$$

The proof is completed.

The rest of the paper, u and U will stand for the solution of problems (1.1) and (2.5), respectively. Writing the error $U - u = (U - V_h u) + (V_h u - u) = \xi + \eta$, we then derive superconvergence estimates of the initial value error $\xi_t(0)$.

Lemma 3.3. For $2 \leq p \leq \infty$, we have

$$\begin{cases} (a) \|\xi_t(0)\|_{1,p} \leq Ch^{r+1} [\|u_0\|_{r,p} + \|u_t(0)\|_{r,p}], & r > 2, \\ (b) \|\xi_t(0)\|_{1,p} \leq Ch^2 [\|u_0\|_{2,p} + \|u_t(0)\|_{2,p}], & r = 2, \\ (c) \|\xi_t(0)\|_{0,p} \leq Ch^{r+2} [\|u_0\|_{r,p} + \|u_t(0)\|_{r,p}], & r > 3. \end{cases} \quad (3.12)$$

Proof. We first combine (2.3), (2.5) and (2.6) to obtain the error equation

$$(\xi_t, \chi) + (a(t)\nabla\xi_t + b(t)\nabla\xi + \int_0^t c(t, \tau)\nabla\xi(\tau)d\tau, \nabla\chi) = -(\eta_t, \chi), \quad \chi \in S_h, \quad t \in J. \quad (3.13)$$

To bound $\|\xi_t(0)\|_{1,p}$, we need an auxiliary problem. For any $\phi \in W^{1,p'}(\Omega)$ with $\|\phi\|_{0,p'} = 1$. Let $\Phi \in H_0^1(\Omega)$ be the solution of

$$(a(t)\nabla v, \nabla\Phi) + (v, \Phi) = -(v, \phi_x), \quad v \in H_0^1(\Omega), \quad t \in J, \quad (3.14)$$

where ϕ_x is some component of $\nabla\phi$. Thus we have

$$\|\Phi\|_{1,p'} \leq C_p(t)\|\phi\|_{0,p'} \leq C_p(t), \quad t \in J, \quad (3.15).$$

According to our assumptions, $C_p(t)$ is uniformly bounded on t . Let Φ_h be the Galerkin approximation of Φ in S_h . Setting $t = 0$ in (3.13) and by (2.5), noting $\xi(0) = 0$, we see that

$$(\xi_t(0), \chi) + (a(0)\nabla\xi_t(0), \nabla\chi) = -(\eta_t(0), \chi). \quad (3.16)$$

Then by Green's formula, (3.14), (3.16), (3.8) and (3.15), we know that

$$\begin{aligned} (\xi_{tx}(0), \phi) &= -(\xi_t(0), \phi_x) \\ &= (a(0)\nabla\xi_t(0), \nabla\Phi) + (\xi_t(0), \Phi) \\ &= (a(0)\nabla\xi_t(0), \nabla\Phi^h) + (\xi_t(0), \Phi^h) \\ &= -(\eta_t(0), \Phi^h) \\ &\leq Ch^{r+1}\|u(0)\|_{1,r,p} \|\Phi^h\|_{1,p'} \\ &\leq Ch^{r+1}\|u(0)\|_{1,r,p} \|\Phi\|_{1,p'} \\ &\leq Ch^{r+1}\|u(0)\|_{1,r,p}. \end{aligned}$$

Thus

$$\|\xi_{tx}(0)\|_{0,p} = \sup_{\substack{\phi \in L_{p'}(\Omega) \\ \|\phi\|_{0,p'}=1}} |(\xi_{tx}(0), \phi)| \leq Ch^{r+1}\|u(0)\|_{1,r,p}. \quad (3.17)$$

By summing on $\|\xi_{tx}(0)\|_{0,p}$, we have

$$\|\xi_t(0)\|_{1,p} \leq Ch^{r+1}\|u(0)\|_{1,r,p} \leq Ch^{r+1}[\|u_0\|_{r,p} + \|u_t(0)\|_{r,p}], \quad r > 2.$$

Arguing as before, we can also derive (3.12b).

Finally we turn to the proof of (3.12c). For this end, we introduce the other auxiliary problem. For any $\psi \in L_{p'}(\Omega)$ with $\|\psi\|_{0,p'} = 1$. Let $\Psi \in W^{2,p'}(\Omega)$ be the solution of

$$(a(t)\nabla v, \nabla \Psi) + (v, \Psi) = (v, \psi), \quad v \in H_0^1(\Omega), \quad t \in J. \quad (3.18)$$

Then we have

$$\|\Psi\|_{2,p'} \leq C_p(t)\|\psi\|_{0,p'} \leq C_p(t), \quad t \in J. \quad (3.19).$$

Here $C_p(t)$ is uniformly bounded on t . Let Ψ^h be the Galerkin approximation of Ψ . Using argument similar to those used to get (3.12a), by (3.18), (3.16), (3.8) and (3.19), we have

$$\begin{aligned} (\xi_t(0), \psi) &= (a(0)\nabla \xi_t(0), \nabla \Psi) + (\xi_t(0), \Psi) \\ &= (a(0)\nabla \xi_t(0), \nabla \Psi^h) + (\xi_t(0), \Psi^h) \\ &= -(\eta_t(0), \Psi^h) \\ &= (\eta_t(0), \Psi - \Psi^h) - (\eta_t(0), \Psi) \\ &\leq |(\eta_t(0), \Psi - \Psi^h)| + |(\eta_t(0), \Psi)| \\ &\leq Ch^{r+1}\|u(0)\|_{1,r,p} \|\Psi - \Psi^h\|_{1,p'} + Ch^{r+2}\|u(0)\|_{1,r,p} \|\Psi\|_{2,p'} \\ &\leq Ch^{r+2}[\|u(0)\|_{r,p} + \|u_t(0)\|_{r,p}]. \end{aligned}$$

Thus,

$$\|\xi_t(0)\|_{0,p} = \sup_{\substack{\psi \in L_{p'}(\Omega) \\ \|\psi\|_{0,p'}=1}} |(\xi_t(0), \psi)| \leq Ch^{r+2}[\|u(0)\|_{r,p} + \|u_t(0)\|_{r,p}].$$

The proof is completed.

4. Superconvergence In $W^{1,p}(\Omega)$ And $L_p(\Omega)$

Our object of this section is to demonstrate the main results of the paper. We shall begin by proving two order global superconvergence estimate of ξ and ξ_t in $W^{1,p}(\Omega)$ for $2 \leq p < \infty$.

Theorem4.1. For $r > 2$ and $2 \leq p < \infty$, we have, for $t \in J$,

$$\begin{cases} (a) \|\xi\|_{1,p} \leq Ch^{r+1} \int_0^t \|u\|_{1,r,p} d\tau, \\ (b) \|\xi_t\|_{1,p} \leq Ch^{r+1} [\|u(0)\|_{1,r,p} + \int_0^t \|u\|_{2,r,p} d\tau]. \end{cases} \quad (4.1)$$

Proof. We first estimate $\|\xi\|_{1,p}$. Let Φ, ϕ satisfy (3.14),(3.15), then an analogy of the proof for (3.12a) implies, by Green's formula, (3.14), (3.13) ,(3.8) and (3.15),

$$\begin{aligned} (\xi_x, \phi) &= \int_0^t (\xi_{tx}, \phi) d\tau = - \int_0^t (\xi_t, \phi_x) d\tau \\ &= \int_0^t [(a(\tau) \nabla \xi_t, \nabla \Phi) + (\xi_t, \Phi)] d\tau \\ &= \int_0^t [(a(\tau) \nabla \xi_t, \nabla \Phi^h) + (\xi_t, \Phi^h)] d\tau \\ &= - \int_0^t [(\eta_t(\tau), \Phi^h) + (b(\tau) \nabla \xi + \int_0^\tau c(\tau, s) \nabla \xi(s) ds, \nabla \Phi^h)] d\tau \\ &\leq C \left\{ \int_0^t [h^{r+1} \|u(\tau)\|_{1,r,p} \|\Phi^h\|_{1,p'} + \|\xi\|_{1,p} \|\Phi^h\|_{1,p'} + \int_0^\tau \|\xi\|_{1,p} ds \|\Phi^h\|_{1,p'}] d\tau \right\} \\ &\leq Ch^{r+1} \int_0^t \|u(\tau)\|_{1,r,p} d\tau \|\Phi\|_{1,p'} + C \int_0^t \|\xi\|_{1,p} d\tau \|\Phi\|_{1,p'} \\ &\leq Ch^{r+1} \int_0^t \|u(\tau)\|_{1,r,p} d\tau + C \int_0^t \|\xi\|_{1,p} d\tau, \end{aligned}$$

Then (4.1a) follows from Gronwall's Lemma.

To bound $\|\xi_t\|_{1,p}$, we differentiate (3.13) with respect to t to get

$$\begin{aligned} &(\xi_{tt}, \chi) + (a(t) \nabla \xi_{tt}, \nabla \chi) + ((a_t(t) + b(t)) \nabla \xi_t, \nabla \chi) + ((b_t(t) + c(t, t)) \nabla \xi(t), \nabla \chi) \\ &+ \left(\int_0^t c_t(t, \tau) \nabla \xi(\tau) d\tau, \nabla \chi \right) = -(\eta_{tt}, \chi), \quad \chi \in S_h. \end{aligned} \quad (4.2)$$

By applying Green's formula, (3.14), (4.2), Hölder inequality, Lemma 3.2, (3.15), (3.12a) and (4.1a), we have

$$\begin{aligned}
(\xi_{tx}, \phi) &= (\xi_{tx}(0), \phi) + \int_0^t (\xi_{ttx}, \phi) d\tau \\
&= (\xi_{tx}(0), \phi) - \int_0^t (\xi_{tt}, \phi_x) d\tau \\
&= (\xi_{tx}(0), \phi) + \int_0^t [(a(\tau) \nabla \xi_{tt}, \nabla \Phi^h) + (\xi_{tt}, \Phi^h)] d\tau \\
&= (\xi_{tx}(0), \phi) - \int_0^t [(\eta_{tt}, \Phi^h) + (a_t(\tau) + b(\tau)) \nabla \xi_t, \nabla \Phi^h] \\
&\quad + (b_t(\tau) + c(\tau, \tau)) \nabla \xi(\tau), \nabla \Phi^h + \left(\int_0^\tau c_\tau(\tau, s) \nabla \xi(s) ds, \nabla \Phi^h \right) d\tau \\
&\leq \|\xi_{tx}(0)\|_{0,p} \|\phi\|_{0,p'} + C \left[\int_0^t \{h^{r+1} \|u\|_{2,r,p} \|\Phi^h\|_{1,p'} \right. \\
&\quad \left. + \|\xi_t\|_{1,p} \|\Phi^h\|_{1,p'} + \|\xi\|_{1,p} \|\Phi^h\|_{1,p'} + \int_0^\tau \|\xi\|_{1,p} d\tau \|\Phi^h\|_{1,p'} \} d\tau \right] \\
&\leq \|\xi_t(0)\|_{1,p} + C \left[\int_0^t \{h^{r+1} \|u\|_{2,r,p} + \|\xi_t\|_{1,p} + \|\xi\|_{1,p} + \int_0^\tau \|\xi\|_{1,p} ds \} d\tau \right] \\
&\leq Ch^{r+1} \|u(0)\|_{1,r,p} + Ch^{r+1} \int_0^t \|u\|_{2,r,p} d\tau + C \int_0^t \|\xi\|_{1,p} d\tau + C \int_0^t \|\xi_t\|_{1,p} d\tau \\
&\leq Ch^{r+1} \|u(0)\|_{1,r,p} + Ch^{r+1} \int_0^t \|u\|_{2,r,p} d\tau + C \int_0^t \|\xi_t\|_{1,p} d\tau.
\end{aligned}$$

The desired estimate (4.1b) is now concluded by Gronwall's Lemma. The proof is completed.

For the case of $r = 2$, we also have the following superconvergence.

Theorem 4.2. For $r = 2$ and $2 \leq p < \infty$, we have, for $t \in J$

$$\begin{cases} (a) \|\xi\|_{1,p} \leq Ch^2 \int_0^t \|u\|_{1,2,p} d\tau, \\ (b) \|\xi_t\|_{1,p} \leq Ch^2 [\|u(0)\|_{1,2,p} + \int_0^t \|u\|_{2,2,p} d\tau]. \end{cases} \quad (4.3)$$

We finally derive two order superconvergence of ξ and ξ_t in $L_p(\Omega)$, for $2 \leq p < \infty$.

Theorem4.3. For $r > 3$, $2 \leq p < \infty$, $t \in J$, we have

$$\begin{cases} (a) \|\xi\|_{0,p} \leq Ch^{r+2} \int_0^t \|u\|_{1,r,p} d\tau, \\ (b) \|\xi_t\|_{0,p} \leq Ch^{r+2} [\|u(0)\|_{1,r,p} + \int_0^t \|u\|_{2,r,p} d\tau]. \end{cases} \quad (4.4)$$

Proof. Let ψ , Ψ satisfy (3.18), (3.19), and Ψ^h is the Galerkin approximation of Ψ . In the same way as the proof for (3.12c), it follows from (3.18) and (3.13) that

$$\begin{aligned} (\xi, \psi) &= \int_0^t (\xi_t, \psi) d\tau \\ &= \int_0^t [(a(\tau)\nabla\xi_t, \nabla\Psi^h) + (\xi_t, \Psi^h)] d\tau \\ &= - \int_0^t [(\eta_t, \Psi^h) + (b(\tau)\nabla\xi + \int_0^\tau c(\tau, s)\nabla\xi(s)ds, \nabla\Psi^h)] d\tau \\ &= \int_0^t (\eta_t, \Psi - \Psi^h) d\tau - \int_0^t (\eta_t, \Psi) d\tau + \int_0^t (b(\tau)\nabla\xi \\ &\quad + \int_0^\tau c(\tau, s)\nabla\xi(s)ds, \nabla(\Psi - \Psi^h)) d\tau \\ &\quad - \int_0^t (b(\tau)\nabla\xi + \int_0^\tau c(\tau, s)\nabla\xi(s)ds, \nabla\Psi) d\tau \\ &= I_1 + I_2 + I_3 + I_4. \end{aligned}$$

Next we estimate $I_1 - I_4$ respectively.

By Lemma3.2, we have

$$\begin{aligned} |I_1| + |I_2| &= \left| \int_0^t (\eta_t, \Psi - \Psi^h) d\tau \right| + \left| \int_0^t (\eta_t, \Psi) d\tau \right| \\ &\leq Ch^{r+1} \int_0^t \|u\|_{1,r,p} \|\Psi - \Psi^h\|_{1,p'} d\tau + Ch^{r+2} \int_0^t \|u\|_{1,r,p} d\tau \|\Psi\|_{2,p'} \\ &\leq Ch^{r+2} \int_0^t \|u\|_{1,r,p} d\tau \|\Psi\|_{2,p'}. \end{aligned} \quad (4.5)$$

By Hölder Inequality and (4.1a)

$$\begin{aligned}
|I_3| &= \left| \int_0^t (b(\tau)\nabla\xi + \int_0^\tau c(\tau, s)\nabla\xi(s)ds, \nabla(\Psi - \Psi^h))d\tau \right| \\
&\leq C \int_0^t (\|\xi\|_{1,p} + \int_0^\tau \|\xi\|_{1,p}ds) \|\Psi - \Psi^h\|_{1,p'}d\tau \\
&\leq C \int_0^t \|\xi\|_{1,p}d\tau \|\Psi - \Psi^h\|_{1,p'} \\
&\leq Ch^{r+2} \int_0^t \|u\|_{1,r,p}d\tau \|\Psi\|_{2,p'}
\end{aligned} \tag{4.6}$$

In order to estimate I_4 , let $B(\tau, \xi, \Psi) = (b(\tau)\nabla\xi + \int_0^\tau c(\tau, s)\nabla\xi(s)ds, \nabla\Psi)$, B^* is the adjoint operator of B, then by Green's formula, we have

$$\begin{aligned}
|I_4| &= \left| \int_0^t (b(\tau)\nabla\xi + \int_0^\tau c(\tau, s)\nabla\xi(s)ds, \nabla\Psi)d\tau \right| \\
&= \left| \int_0^t (\xi, B^*(\tau)\Psi)d\tau \right| \\
&\leq C \int_0^t \|\xi\|_{0,p}d\tau \|\Psi\|_{2,p'}.
\end{aligned} \tag{4.7}$$

Combine (4.5)-(4.7) and by (3.19), we have

$$|(\xi, \psi)| \leq Ch^{r+2} \int_0^t \|u\|_{1,r,p}d\tau + C \int_0^t \|\xi\|_{0,p}d\tau.$$

Then (4.4a) follows from Gronwall's Lemma.

We next turn to the proof of (4.4b). An analogue of the proof for(4.1b), let ψ, Ψ

satisfy (3.18), (3.19), by (3.18) and (4.2), we have

$$\begin{aligned}
(\xi_t, \psi) &= (\xi_t(0), \psi) + \int_0^t (\xi_{tt}, \psi) d\tau \\
&= (\xi_t(0), \psi) + \int_0^t [(a(\tau)\nabla\xi_{tt}, \nabla\Psi) + (\xi_{tt}, \Psi)] d\tau \\
&= (\xi_t(0), \psi) + \int_0^t [a(\tau)\nabla\xi_{tt}, \nabla\Psi^h) + (\xi_{tt}, \Psi^h)] d\tau \\
&= (\xi_t(0), \psi) - \int_0^t [(\eta_{tt}, \Psi^h) + ((a_t(\tau) + b(\tau))\nabla\xi_t, \nabla\Psi^h) \\
&\quad + ((b_t(\tau) + c(\tau, \tau))\nabla\xi(\tau), \nabla\Psi^h) + (\int_0^\tau c_\tau(\tau, s)\nabla\xi(s)ds, \nabla\Psi^h)] d\tau \\
&= (\xi_t(0), \psi) + \int_0^t (\eta_{tt}, \Psi - \Psi^h) d\tau - \int_0^t (\eta_{tt}, \Psi) d\tau \\
&\quad + \int_0^t ((a_t(\tau) + b(\tau))\nabla\xi_t, \nabla(\Psi - \Psi^h)) d\tau \\
&\quad - \int_0^t ((a_t(\tau) + b(\tau))\nabla\xi_t, \nabla\Psi) d\tau + \int_0^t ((b_t(\tau) + c(\tau, \tau))\nabla\xi(\tau), \nabla(\Psi - \Psi^h)) d\tau \\
&\quad - \int_0^t ((b_t(\tau) + c(\tau, \tau))\nabla\xi(\tau), \nabla\Psi) d\tau + \int_0^t (\int_0^\tau c_\tau(\tau, s)\nabla\xi(s)ds, \nabla(\Psi - \Psi^h)) d\tau \\
&\quad - \int_0^t (\int_0^\tau c_\tau(\tau, s)\nabla\xi(s)ds, \nabla\Psi) d\tau \\
&= I_1 + \dots + I_9. \tag{4.8}
\end{aligned}$$

Then by Hölder Inequality, (3.12c) and Lemma 3.2, (4.1b) and (3.19), we have

$$\begin{aligned}
|I_1| + |I_2| + |I_3| + |I_4| &\leq Ch^{r+2}[\|u_0\|_{r,p} + \|u_t(0)\|_{r,p}] \\
&\quad + Ch^{r+1} \int_0^t \|u\|_{2,r,p} d\tau \|\Psi - \Psi^h\|_{1,p'} \\
&\quad + Ch^{r+2} \int_0^t \|u\|_{2,r,p} d\tau \|\Psi\|_{2,p'} \\
&\quad + C \int_0^t \|\xi_t\|_{1,p} d\tau \|\Psi - \Psi^h\|_{1,p'} \\
&\leq Ch^{r+2}[\|u_0\|_{1,r,p} + \int_0^t \|u\|_{2,r,p} d\tau] \|\Psi\|_{2,p'}.
\end{aligned}$$

By (4.1a)

$$\begin{aligned} |I_6| + |I_8| &\leq C \int_0^t \|\xi\|_{1,p} d\tau \|\Psi - \Psi^h\|_{1,p'} + C \int_0^t \left(\int_0^\tau \|\xi\|_{1,p} ds \right) d\tau \|\Psi - \Psi^h\|_{1,p'} \\ &\leq C \int_0^t \|\xi\|_{1,p} d\tau \|\Psi - \Psi^h\|_{1,p'} \\ &\leq Ch^{r+2} \int_0^t \|u\|_{1,r,p} d\tau \|\Psi\|_{2,p'} \end{aligned}$$

As in the proof for (4.7), we have, by (4.4a)

$$\begin{aligned} |I_5| + |I_7| &\leq C \int_0^t \|\xi_t\|_{0,p} d\tau \|\Psi\|_{2,p'} + C \int_0^t \|\xi\|_{0,p} d\tau \|\Psi\|_{2,p'} \\ &\leq C[h^{r+2} \int_0^t \|u\|_{1,r,p} d\tau + \int_0^t \|\xi_t\|_{0,p} d\tau] \|\Psi\|_{2,p'}. \end{aligned}$$

And

$$|I_9| \leq C \int_0^t \left(\int_0^\tau \|\xi\|_{0,p} ds \right) d\tau \|\Psi\|_{2,p'} \leq C \int_0^t \|\xi\|_{0,p} d\tau \|\Psi\|_{2,p'}.$$

Combining $I_1 - I_9$ and (4.8), we have

$$|(\xi_t, \psi)| \leq Ch^{r+2} [\|u_0\|_{1,r,p} + \int_0^t \|u\|_{2,r,p} d\tau] \|\Psi\|_{2,p'} + C \int_0^t \|\xi_t\|_{0,p} d\tau \|\Psi\|_{2,p'}.$$

The proof of (4.4b) is now concluded by (3.19) and Gronwall's Lemma.

The proof is completed.

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