# DYNAMICS OF tRNA<sup>val</sup> MEASURED WITH A LONG-LIFETIME METAL-LIGAND COMPLEX

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 $[Ru(bpy)_2(dppz)]^{2+}$  (bpy = 2,2'-bipyridine, dppz = dipyrido[3,2-a:2',3'-c]phenazine) (RuBD), a long-lifetime metal-ligand complex displays favorable photophysical properties including long lifetime, polarized emission, and very little background fluorescence. To further show the usefulness of this luminophore (RuBD) for probing nucleic acid dynamics, its intensity and anisotropy decays when bound to tRNA<sup>val</sup> were examined using frequency-domain fluorometry with a blue light-emitting diode (LED) as the modulated light source. Unexpectedly much longer mean lifetime was obtained at  $4^{\circ}C$  ( $<\tau>=178.3$  ns) as compared to at  $25^{\circ}C$  ( $<\tau>=117.0$  ns), suggesting more favorable conformation of tRNA<sup>val</sup> for RuBD when intercalated at  $4^{\circ}C$ . The anisotropy decay data showed longer rotational correlation times at  $4^{\circ}C$  (52.7 and 13.0 ns) than at  $25^{\circ}C$  (32.9 and 10.3 ns). The presence of two rotational correlation times suggests that RuBD reveals both local and overall rotational motion of tRNA<sup>val</sup>. Due to long lifetime of RuBD and small size of tRNA<sup>val</sup>, very low steady-state anisotropy values were observed, 0.048 and 0.036 at 4 and  $25^{\circ}C$ , respectively. However, a clear difference in the modulated anisotropy values was seen between 4 and  $25^{\circ}C$ . These results indicate that RuBD can be useful for studying hydrodynamics of small nucleic acids such as tRNA<sup>val</sup>.

**key words:** tRNA<sup>val</sup>; long-lifetime metal ligand complex; frequency-domain fluorometry; anisotropy decay; light-emitting diode

### INTRODUCTION

Long-lifetime metal-ligand complexes (MLCs) which display decay times ranging from 100 ns to more than 10 µs have only recently become available [1-3] and have been proven to have favorable chemical, photochemical and photophysical properties. Because of large Stokes' shift, the MLCs do not self-quench [2,3]. In addition, the MLCs show good water solubility and high chemical and photochemical stability [2,3]. The long lifetimes of the MLCs allow the use of gated detection, which can be employed to suppress interfering autofluorescence from biological samples and can thus provide increased sensitivity [4]. And finally, the MLCs display polarized emission, making them useful for microsecond dynamics [2,3].

Barton and coworkers [5-7] reported that the dipyrido[3,2-a:2',3'-c]phenazine (dppz) complexes of ruthenium appear to be a prime candidate for a spectroscopic probe for nucleic acids because of their "molecular light switch" properties for DNA. They are known to be nonfluorescent in water but emit when the dppz ligand is intercalated into the DNA helix. Emission from the dppz complexes is quenched

when the nitrogens on the dppz are exposed to water. These nitrogens are shielded from water when the dppz intercalates into double-helical DNA, and the complex becomes fluorescent. Since the luminescent enhancement upon DNA binding is ≥10⁴, there is essentially no background with the dppz complexes of ruthenium. Lakowicz and colleagues [8,9] introduced the use of the dppz complexes of ruthenium to study the hydrodynamics of nucleic acids. Using the ruthenium complexes, they examined the intensity and anisotropy decays of calf thymus DNA and could observe rotational correlation times up to several hundred nanoseconds.

In the present study, I have examined the intensity and anisotropy decays of  $[Ru(bpy)_2(dppz)]^{2+}$  (bpy = 2,2'-bipyridine) (RuBD) when bound to tRNAval in order to further show the usefulness of this luminophore (RuBD) for probing nucleic acid dynamics. For this study, frequency-domain fluorometry was used with a high-intensity, blue light-emitting diode (LED) as the light source. With this LED I was able to directly modulate the excitation light up to 100 MHz without the need for an external modulator like a Pockels cell which reduces the light amount to a great extent. Due to very low quantum yield (Q = 0.008) of RuBD [10], time-domain fluorometry was used in the previous studies of intensity and anisotropy decays of RuBD [8,9]. However, with a highintensity, blue LED as the light source, I have been able to obtain very reliable time-resolved intensity and anisotropy decays using the frequency-domain method.

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#### MATERIALS AND METHODS

*Materials* tRNA<sup>val</sup> was obtained from Sigma (St. Louis, MO, U.S.A.) and used without further purification. RuBD was synthesized by the method described previously [8,9]. All other chemicals were of the reagent grade, and water was deionized with a Milli-Q system. All measurements were carried out in 100 mM KCl, 10 mM Tris, pH 7.4.

Absorption and steady-state fluorescence measurement The tRNA  $^{\rm val}$  concentration was 600  $\mu M$  while that of RuBD was 15 M. The concentrations of tRNA  $^{\rm val}$  and RuBD were determined using molar extinction coefficients of 8,200  $M^{-1} cm^{-1}$  (expressed as nucleotide) at 260 nm and 13,000  $M^{-1} cm^{-1}$  at 440 nm, respectively. UV-visible absorption spectra were measured with a Hewlett-Packard 8453 diode array spectrophotometer. Steady-state intensity and anisotropy measurements were carried out using an Aminco SLM AB2 spectrofluorometer (Spectronic Instruments, Inc., Rochester, NY, U.S.A.). RuBD was excited at 440 nm.

Frequency-domain intensity and anisotropy decay measurements Measurements were performed using the instruments described previously [11] and modified with a data acquisition card from ISS, Inc. (Urbana, IL, U.S.A.) [12]. The excitation source was a blue LED LNG992CFBW (Panasonic, Japan) with luminous intensity of 1,500 mcd. An LED driver LDX-3412 (ILX Lightwave, Boseman, MO, U.S.A.) provided 30 mA of current at frequencies from 0.4 to 15 MHz. A  $450\pm20$  nm interference filter and a 620 nm cut-off filter were used for isolating excitation and emission, respectively. Rhodamine B in water ( $\tau = 1.68$  ns) was utilized as a lifetime standard.

The intensity decays were recovered from the frequency-domain data in terms of a multiexponential model:

$$I(t) = \sum_{i=1}^{n} \alpha_i e^{-t/\tau_i}$$
 (1)

where the preexponential factors  $\alpha_i$  are the amplitude of each component,  $\Sigma \alpha_i = 1.0$ ,  $\tau_i$  are the decay times, and n are the numbers of exponential components. These values were determined by nonlinear least squares analysis as described previously [13,14]. Mean lifetimes were calculated according to Eq. (2):

$$<\tau> = \frac{\sum_{i} \alpha_{i} \tau_{i}^{2}}{\sum_{i} \alpha_{i} \tau_{i}} = \sum_{i} f_{i} \tau_{i}$$
 (2)

where  $f_i$  is the fractional steady-state contribution of each component to the total emission, and  $\Sigma f_i$  is normalized to unity.  $f_i$  is given by

$$f_i = \frac{\alpha_i \tau_i}{\sum_i \alpha_j \tau_j} \tag{3}$$

The best fits were obtained by a minimum value for the goodness-of-fit parameters  $\chi_R^2$ :

$$\chi_R^2 = \frac{1}{\nu} \sum_{\omega} \left[ \left( \frac{\varphi_{\omega} - \varphi_{c\omega}}{\delta \varphi} \right)^2 + \left( \frac{m_{\omega} - m_{c\omega}}{\delta m} \right)^2 \right] \tag{4}$$

where  $\nu$  is the number of degrees of freedom, and  $\phi_{\omega}$  and  $m_{\omega}$  are the experimental phase and modulation, respectively. The subscript c is used to indicate calculated values for assumed values of  $\alpha_i$  and  $\tau_i$ , and  $\delta \phi$  and  $\delta m$  are the experimental uncertainties. Here, the values of  $\delta \phi$  and  $\delta m$  were set at 0.2° and 0.005, respectively.

The frequency-domain anisotropy decays were also analyzed in terms of the multiexponential model:

$$r(t) = r_0 \sum_{i} g_i e^{-t/\theta_i}$$
 (5)

where  $g_i$  is the amplitude of the anisotropy component with a rotational correlation time  $\theta_i$ ,  $\Sigma g_i = 1.0$ , and  $r_o$  is the anisotropy in the absence of rotational diffusion. The total anisotropy  $r_o$  was a fitted parameter.

#### RESULTS AND DISCUSSION

In this report, I have studied the hydrodynamic properties of tRNA<sup>val</sup> using a long-lifetime MLC, RuBD. The chemical structure of RuBD is shown in Fig. 1. Fig. 2 shows the emission spectrum of RuBD bound to tRNA<sup>val</sup>, which shows an emission peak at about 620 nm. In aqueous solution the luminescence of this probe is undetectable. The frequency-domain intensity decays of RuBD bound to tRNA<sup>val</sup> are shown in Fig. 3. The intensity decays were best fit by a triple exponential decay. Table 1 shows the intensity decay results at 4 and 25°C. Mean lifetime values (<\tau>> ) obtained at 4 and 25°C were 178.3 and 117.0 ns, respectively. As expected, increased decay time was observed at 4°C (about 150%) in comparison to 25°C, probably due to the reduced collisional motions at the lower temperature. However, the important point is the magnitude of the difference. I also have mea-

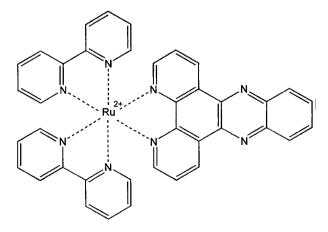


Fig. 1. Chemical structure of [Ru(bpy)<sub>2</sub>(dppz)]<sup>2+</sup> (RuBD).

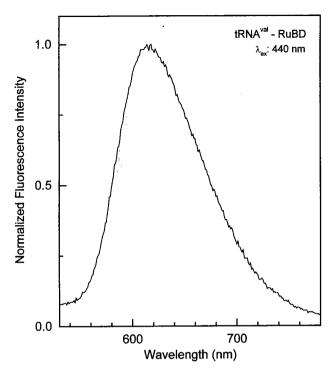


Fig. 2. Emission spectrum of [Ru(bpy)<sub>2</sub>(dppz)]<sup>2+</sup> (RuBD) bound to tRNA<sup>val</sup>. Excitation was at 440 nm.

Table 1. Multiexponential intensity decay analyses of  $[Ru(bpy)_2(dppz)]^{2+}$  bound to  $tRNA^{val}$  at different temperatures

| Temperature(°C) | τi (ns) | $\alpha_i$ | $f_i^{\mathrm{a}}$ | < <b>t&gt;</b> <sup>a</sup> (ns) | $\chi_R^{\rm 2b}$ |
|-----------------|---------|------------|--------------------|----------------------------------|-------------------|
| 4               | 280.3   | 0.12       | 0.55               | 178.3                            | 2.5               |
|                 | 61.0    | 0.37       | 0.37               |                                  |                   |
|                 | 8.9     | 0.51       | 0.08               |                                  |                   |
| 25              | 180.0   | 0.16       | 0.55               | 117.0                            | 2.9               |
|                 | 48.1    | 0.40       | 0.35               |                                  |                   |
|                 | 12.0    | 0.44       | 0.10               |                                  |                   |

<sup>&</sup>lt;sup>a</sup>Fractional intensities  $f_i$  and mean lifetimes  $<\tau>$  were calculated using Eqs. (3) and (2), respectively.

sured the intensity decays of RuBD bound to supercoiled, cut and relaxed pTZ18U plasmids at 4 and 25°C (unpublished data). However, in all forms of pTZ18U plasmids, the differences were much smaller than in tRNA<sup>val</sup>. The mean intensity decay times of RuBD bound to supercoiled, cut and relaxed pTZ18U plasmids at 4°C were each about 120% of those at 25°C. It seems likely that the RuBD MLC was much more efficiently shielded from water at 4°C than at 25°C for tRNA<sup>val</sup>. In other words, RuBD binds much more efficiently to tRNA<sup>val</sup> adopts more favorable conformations for RuBD when bound at 4°C. Still, the mean lifetime at 25°C ( $<\tau>=117.0$  ns) is comparable to the values previously reported for calf thymus DNA [6-10].

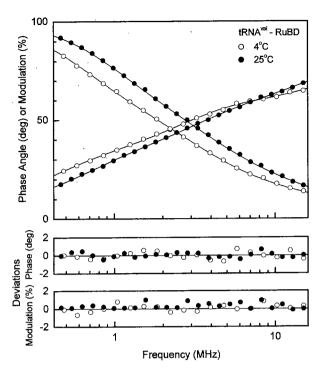


Fig. 3. Frequency-domain intensity decays of [Ru(bpy)<sub>2</sub>(dppz)]<sup>2+</sup> (RuBD) bound to tRNA<sup>val</sup> at 4 and 25°C. The open and closed circles represent the values at 4 and 25°C, respectively. The solid lines in the upper panel show the best triple-exponential fits to the data. The middle and lower panels show the plots of residuals between the experimental data and the fitted curves.

Table 2. Anisotropy decay analyses of [Ru(bpy)<sub>2</sub>(dppz)]<sup>2+</sup> bound to tRNA<sup>val</sup> at different temperatures.

| $\theta_i$ (ns) | $r_o*g_i$            | $\Sigma(r_o^*g_i)$                     | $\chi_R^{2a}$   |
|-----------------|----------------------|--|---|
| 52.7            | 0.087                | 0.129                                  | 4.5   |
| 13.0            | 0.042                |  |   |
| 32.9            | 0.050                | 0.128                                  | 4.2   |
| 10.3            | 0.078                |  |   |
|                 | 52.7<br>13.0<br>32.9 | 52.7 0.087<br>13.0 0.042<br>32.9 0.050 | 52.7    0.087    0.129      13.0    0.042      32.9    0.050    0.128 |

<sup>a</sup>The goodness-of-fit parameters  $\chi_R^2$  were calculated by Eq. (4), and the standard errors of phase angle and modulation were set at 0.2° and 0.005, respectively.

In addition to the intensity decay measurements, the anisotropy decays of RuBD bound to tRNA<sup>val</sup> were also measured at 4 and 25°C (Fig. 4), and the results are summarized in Table 2. The best fits of the anisotropy decay data were obtained using the two exponential model. The slow rotational correlation times (52.7 and 32.9 ns at 4 and 25°C, respectively) appear to be consistent with those expected for overall rotational mobility of tRNA<sup>val</sup>. The fast rotational correlation times (13.0 and 10.3 ns, respectively, at 4 and 25°C) may be due to independent motion of two helical arms of the L-shaped structure of tRNA<sup>val</sup> because the correlation time constant of local internal motion of the bases within tRNA is usually in the range of picoseconds [15]. X-ray crystal struc-

<sup>&</sup>lt;sup>b</sup>The goodness-of-fit paramaeters  $\chi_R^2$  were calculated by Eq. (4), and the standard errors of phase angle and modulation were set at 0.2° and 0.005, respectively.

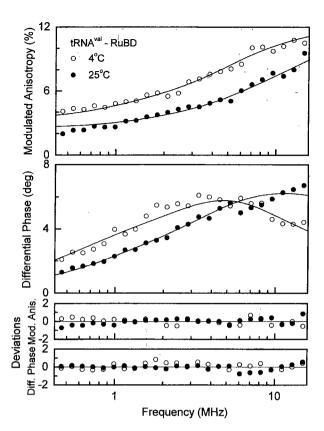


Fig. 4. Frequency-domain anisotropy decays of [Ru(bpy)<sub>2</sub>(dppz)]<sup>2+</sup> (RuBD) bound to tRNA<sup>val</sup> at 4 and 25°C. The open and closed circles represent the values at 4 and 25°C, respectively. The modulated anisotropies and differential phases are shown in the first and second panels, respectively. The solid lines in the first and second panels show the best double-exponential fits to the data. The lower two panels show the plots of residuals between the experimental data and the fitted curves.

ture analyses have revealed that tRNAs have more or less the same L-shaped configuration [16,17]. From the result of the present study, it seems that the two helical arms of the L-shape show identical rotational correlation times (13.0 and 10.3 ns, respectively, at 4 and 25°C). As expected, longer rotational correlation times were observed at 4°C (52.7 and 13.0 ns) in contrast to 25°C (32.9 and 10.3 ns). Thus, it seems that both the overall and local rotational mobility of tRNAval depend on temperature. However, it should be pointed out that the overall rotational correlation times were slightly longer than the values reported previously for tRNAval by fluorescence anisotropy decays [18] and by electrooptical measurements [19]. Hence, the possibility of the presence of somewhat denatured extended form of tRNAval cannot be ruled out. A plausible explanation may be the probe concentration, about 2:1 of RuBD to tRNAval molar ratio. Because of very low quantum yield (Q = 0.008) of RuBD [10], quite high probe concentration was used. The absence of Mg<sup>2+</sup> in the buffer, which plays an important role in tRNA conformation, can be a factor of the possibility as well. Fig. 4 (first panel) also shows the modulated anisotropy values for RuBD

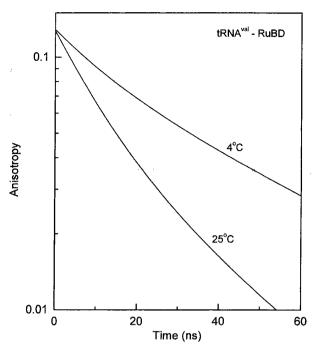


Fig. 5. Time-domain representation of anisotropy decays of  $[Ru(bpy)_2(dppz)]^{2+}$  (RuBD) bound to tRNA<sup>val</sup> at 4 and 25°C.

bound to tRNAval. Of course, lower modulated anisotropy values were observed at 4°C than those at 25°C. I also measured the steady-state anisotropy, however, the steady-state anisotropy values were quite low, 0.048 and 0.036 at 4 and 25°C, respectively. It is thus likely that RuBD, a long-lifetime MLC, is not suitable for measuring the steady-state anisotropy values of small nucleic acids. However, the results of time-resolved anisotropy decay measurements clearly showed that RuBD can be useful for studying hydrodynamics of smaller nucleic acids such as tRNAval. It is informative to examine the anisotropy decays in the time-domain reconstructed from the frequency-domain data (Fig. 5). The timedomain representation of multiexponential analysis data explicitly shows an initial rapid decrease in anisotropy, followed by a slower loss of the remaining anisotropy at longer times, which is attributed to overall rotational motion of tRNAval.

The dppz complexes of ruthenium bind to double helical DNA through intercalation with affinities exceeding 10<sup>7</sup> M<sup>-1</sup>. There are two common dppz complexes of ruthenium, the bpy derivative RuBD and the 1,10-phenanthroline (phen) derivative [Ru(phen)<sub>2</sub>(dppz)]<sup>2+</sup>. The phen derivative has some advantages over the other bpy derivative RuBD because of its longer lifetimes and higher quantum yield [6,7,9]. However, Jenkins *et al.* [6] showed that RuBD exhibited better sensitivity to conformational differences in DNA because of the incomplete shielding of the dppz ligand from water in the presence of bpy in contrast to the other phen derivative. The use of MLCs to measure nucleic acid hydrodynamics

is just beginning, and additional MLCs for nucleic acid dynamics are yet to be developed.

An interesting aspect of this study was the use of a blue LED as the source of excitation. Using it I was able to obtain approximately 3 to 5 mW of excitation power depending on the LED current. The use of LED allowed a simpler "turnkey" operation, i.e., less optical adjustments and easier electronic operations. Direct electronic modulation of up to 100 MHz eliminates the expensive and cumbersome electrooptic modulators found in most frequency-domain instruments [20]. The low cost of LEDs, about \$2 to \$10 a piece, together with its simpler mode of operation will allow laboratories to perform time-resolved fluorescence measurements with less sophisticated fluorescence equipments in an economical way. In general, the use of the new high-intensity UV, blue and green GaN LEDs will facilitate the use of timeresolved measurements for many laboratories and a wide variety of applications.

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