

CLIFFORD L^2 -COHOMOLOGY ON THE COMPLETE KÄHLER MANIFOLDS

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0. Introduction

In the study of a manifold M , the exterior algebra Λ^*M plays an important role. In fact, the de Rham cohomology theory gives many informations of a manifold. Another important object in the study of a manifold is its Clifford algebra $Cl(M)$, generated by the tangent space. It carries an intrinsic first order elliptic operator D , the Dirac operator. There is a canonical vector (but not algebra) bundle isomorphism $\Lambda^*(M) \rightarrow Cl(M)$. In $\Lambda^*(M)$, the Dirac operator D is $D \cong d + \delta$, where d is the exterior differential and δ is the adjoint operator of d . Therefore many results of the Clifford theory yield the results of the de Rham theory([8]). Moreover the calculus of the pair $Cl(M)$, D carries over formally to bundles of modules over $Cl(M)$. On Kähler manifolds, we obtain operators \mathcal{D} and $\bar{\mathcal{D}}$ such that $\mathcal{D}^2 = \bar{\mathcal{D}}^2 = 0$, $\mathcal{D} + \bar{\mathcal{D}} = \frac{1}{2}D$ and $\bar{\mathcal{D}}$ is the formal adjoint of \mathcal{D} . Using these operators, M. L. Michelsohn([10]) studied the Clifford and spinor cohomology theory and proved some vanishing theorems on compact Kähler manifold. In this paper, we study the Clifford L^2 -cohomology theory, the decomposition theorem for the L^2 -Clifford algebra $L^2(Cl^{p,q}(M))$ and prove some vanishing theorems on complete Kähler manifold.

1. Preliminaries

Let M be a $2n$ -dimensional Kähler manifold with almost complex

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structure J and with connection ∇ . Let $Cl(M)$ be the Clifford bundle generated by the tangent bundle TM . Now we define a derivation $\mathcal{J}_0 : Cl(M) \rightarrow Cl(M)$ induced by J as follows:

$$(1.1) \quad \mathcal{J}_0(v_1 \cdots v_k) = \sum_{j=1}^k v_1 \cdots Jv_j \cdots v_k$$

for $v_1, \dots, v_k \in TM$, where “ \cdot ” is the Clifford multiplication. If it is clear from the context which multiplication is meant, we omit the Clifford multiplication “ \cdot ”. To study \mathcal{J}_0 effectively we consider the complexification $\mathbb{C}l(M) = Cl(M) \otimes_{\mathbb{R}} \mathbb{C}$. This algebra has a natural basis given as follows: Let $e_1, \dots, e_n, Je_1, \dots, Je_n$ be an orthonormal basis of $T_x M$. Let $T_x^{1,0}$ (resp. $T_x^{0,1}$) be the i eigenspace (resp. $-i$ eigenspace) of J in $T_x M \otimes \mathbb{C}$. Put

$$\xi_k = \frac{1}{2}\{e_k - iJe_k\}, \quad \bar{\xi}_k = \frac{1}{2}\{e_k + iJe_k\}.$$

Then ξ_1, \dots, ξ_n (resp. $\bar{\xi}_1, \dots, \bar{\xi}_n$) is the basis of $T_x^{1,0}$ (resp. $T_x^{0,1}$). And $\{\xi_k, \bar{\xi}_k\}$ has the following properties;

$$(1.2) \quad \xi_k \bar{\xi}_\ell + \bar{\xi}_k \xi_\ell = \xi_k \bar{\xi}_\ell + \bar{\xi}_\ell \xi_k = -\delta_{k\ell}, \quad \xi_k \xi_\ell = -\xi_\ell \xi_k, \quad \bar{\xi}_k \bar{\xi}_\ell = -\bar{\xi}_\ell \bar{\xi}_k.$$

Denote $\xi_K \bar{\xi}_I = \xi_{k_1} \cdots \xi_{k_r} \bar{\xi}_{i_1} \cdots \bar{\xi}_{i_s}$, where K and I range over all strictly ascending multiindices from $\{1, \dots, n\}$. For convenience we set $\mathcal{J} = \frac{1}{i} \mathcal{J}_0$. Then by the derivation property, we have

$$(1.3) \quad \mathcal{J}(\xi_K \bar{\xi}_I) = (|K| - |I|)\xi_K \bar{\xi}_I,$$

where $|K|, |I|$ denote the lengths of K and I . This gives a decomposition

$$\mathbb{C}l(M) = \bigoplus_{p=-n}^n \mathbb{C}l^p(M),$$

where $\mathbb{C}l^p(M) = \{\phi \in \mathbb{C}l(M) \mid \mathcal{J}\phi = p\phi\}$.

We now introduce two intrinsically defined linear maps $\mathcal{L}, \bar{\mathcal{L}} : \mathbb{C}l(M) \rightarrow \mathbb{C}l(M)$ as follows; For any $\varphi \in \mathbb{C}l(M)$, set

$$(1.4) \quad \mathcal{L}(\varphi) = -\sum_{k=1}^n \xi_k \varphi \bar{\xi}_k, \quad \bar{\mathcal{L}}(\varphi) = -\sum_{k=1}^n \bar{\xi}_k \varphi \xi_k.$$

These operators are independent of the Hermitian basis chosen to define them. We consider the operator $\mathcal{H} = [\mathcal{L}, \bar{\mathcal{L}}]$. Then they satisfy the following relations;

$$(1.5) \quad [\mathcal{L}, \bar{\mathcal{L}}] = \mathcal{H}, \quad [\mathcal{H}, \mathcal{L}] = 2\mathcal{L}, \quad [\mathcal{H}, \bar{\mathcal{L}}] = -2\bar{\mathcal{L}}.$$

Hence they define a representation of $sl(2, \mathbb{C})$, the Lie algebra of $SL(2, \mathbb{C})$, on $\mathbb{C}l(M)$. Since each of the operators $\mathcal{L}, \bar{\mathcal{L}}$ and \mathcal{H} commutes with \mathcal{J} , we can define the subspaces

$$\mathbb{C}l^{p,q}(M) = \{\varphi \in \mathbb{C}l(M) \mid \mathcal{H}\varphi = q\varphi, \mathcal{J}\varphi = p\varphi\}$$

and obtain a decomposition([10])

$$(1.6) \quad \mathbb{C}l(M) = \bigoplus_{p,q} \mathbb{C}l^{p,q}(M).$$

PROPOSITION 1.1([10]). *For each $\xi \in T^{1,0}(M)$, one has that $\xi \cdot \mathbb{C}l^{p,q} \subseteq \mathbb{C}l^{p+1,q+1}$ and $\bar{\xi} \cdot \mathbb{C}l^{p,q} \subseteq \mathbb{C}l^{p-1,q-1}$. Furthermore, if $\xi \neq 0$, the sequences*

$$\begin{aligned} \dots &\xrightarrow{\lambda_\xi} \mathbb{C}l^{p-1,q-1} \xrightarrow{\lambda_\xi} \mathbb{C}l^{p,q} \xrightarrow{\lambda_\xi} \mathbb{C}l^{p+1,q+1} \longrightarrow \dots \\ \dots &\xleftarrow{\lambda_{\bar{\xi}}} \mathbb{C}l^{p-1,q-1} \xleftarrow{\lambda_{\bar{\xi}}} \mathbb{C}l^{p,q} \xleftarrow{\lambda_{\bar{\xi}}} \mathbb{C}l^{p+1,q+1} \longleftarrow \dots, \end{aligned}$$

where λ_ξ denotes left Clifford multiplication by ξ , are exact.

Moreover, these subspaces $\mathbb{C}l^{p,q}$ have the following properties: If $q - s \neq p + r$, then $\mathbb{C}l^{p,q} \cdot \mathbb{C}l^{r,s} = \{0\}$. Otherwise, $\mathbb{C}l^{p,q} \cdot \mathbb{C}l^{r,q-p-r} \subseteq \mathbb{C}l^{p+r,q-r}$. In particular, $\mathbb{C}l^{p,q} \cdot \mathbb{C}l^{p,q} \subseteq \mathbb{C}l^{p,q}$, $\mathbb{C}l^{k,k} \cdot \mathbb{C}l^{\ell,-\ell} \subseteq \mathbb{C}l^{k+\ell,k-\ell}$ and $\mathbb{C}l^{k,\ell} \cdot \mathbb{C}l^{-k,\ell} \subseteq \mathbb{C}l^{0,\ell+k}$.

2. Clifford cohomology group

We recall some facts from [3]: Consider Hilbert spaces H_i ($0 \leq i \leq N$), $H_{N+1} := 0$ and closed operators $D_i : H_i \rightarrow H_{i+1}$, with D_i^* , the adjoint operator. Let $dom D_i$ be the domain of D_i and $ran D_i$ the range of D_i . We then assume that

$$ran D_i \subset dom D_{i+1} \quad \text{and} \quad D_{i+1} \circ D_i = 0.$$

Thus we obtain a complex

$$(2.1) \quad 0 \longrightarrow \text{dom}D_0 \xrightarrow{D_0} \text{dom}D_1 \xrightarrow{D_1} \dots \xrightarrow{D_{N-1}} \text{dom}D_N \longrightarrow 0$$

in the sense of homological algebra with additional functional analytic structure, which is called a *Hilbert complex*. We will abbreviate the complex (2.1) as $(\text{dom}D, D)$.

LEMMA 2.1([3])(THE WEAK HODGE DECOMPOSITION). *Let $(\text{dom}D, D)$ be a Hilbert complex. Then for each i , we have an orthogonal decomposition*

$$(2.2) \quad H_i = \hat{\mathcal{H}}_i \oplus \overline{\text{im}D_{i-1}} \oplus \overline{\text{im}D_i^*}$$

where $\hat{\mathcal{H}}_i := \text{Ker}D_i \cap \text{Ker}D_{i-1}^*$.

Put $\Delta_i := D_i D_i^* + D_i^* D_i$. Then we have

LEMMA 2.2([3]). $\hat{\mathcal{H}}_i = \text{Ker}\Delta_i$.

Now, let $E_i \rightarrow M$ ($0 \leq i \leq N$) be hermitian vector bundles over a Riemannian manifold M and $d_i := \Gamma_{\text{cpt}}(E_i) \rightarrow \Gamma_{\text{cpt}}(E_{i+1})$ differential operators such that $d_i \circ d_{i-1} = 0$. Denote the formal adjoint d_i^t by d_i^t . Then d_i has a closed extension $d_{i,\text{max}}$ in the Hilbert space $H_i := L^2(E_i)$ given by

$$d_{i,\text{max}} := (d_{i,\text{min}}^t)^*,$$

where $d_{i,\text{min}}$ is the minimal extension or the closure of d_i . Then we have

LEMMA 2.3([3]). *If $(\Gamma_{\text{cpt}}(E_i), d_i)$ is an elliptic complex, then*

$$\dots \longrightarrow \text{dom}d_{i-1,\text{max}} \xrightarrow{d_{i-1,\text{max}}} \text{dom}d_{i,\text{max}} \xrightarrow{d_{i,\text{max}}} \text{dom}d_{i+1,\text{max}} \longrightarrow \dots$$

is a Hilbert complex.

Suppose now that M is a complete Kähler manifold. We introduce two differential operators $\mathcal{D}, \bar{\mathcal{D}} : \Gamma\text{Cl}(M) \rightarrow \Gamma\text{Cl}(M)$ by the formulas

$$(2.3) \quad \mathcal{D} = \sum_j \xi_j \nabla_{\bar{\xi}_j}, \quad \bar{\mathcal{D}} = \sum_j \bar{\xi}_j \nabla_{\xi_j},$$

where ∇ is the canonical connection. Since ∇ preserves the subbundles $\Gamma\text{Cl}^{p,q}(M)$, we have

$$\mathcal{D}(\Gamma\text{Cl}^{p,q}) \subset \Gamma\text{Cl}^{p+1,q+1}, \quad \bar{\mathcal{D}}(\Gamma\text{Cl}^{p,q}) \subset \Gamma\text{Cl}^{p-1,q-1}$$

for all p and q . Then we have the following well known fact:

THEOREM 2.4([10]). *The operators \mathcal{D} and $\bar{\mathcal{D}}$ are formal adjoints of one another on $\Gamma_{cpt}\mathbb{C}l(M)$, the set of all sections with the compact support. And they satisfy*

$$\mathcal{D}^2 = \bar{\mathcal{D}}^2 = 0.$$

Furthermore, the complex

$$\dots \xrightarrow{\mathcal{D}} \Gamma\mathbb{C}^{p-1,q-1} \xrightarrow{\mathcal{D}} \Gamma\mathbb{C}^{p,q} \xrightarrow{\mathcal{D}} \Gamma\mathbb{C}^{p+1,q+1} \xrightarrow{\mathcal{D}} \dots$$

is elliptic.

Now we set

$$(2.4) \quad \Delta := \mathcal{D}\bar{\mathcal{D}} + \bar{\mathcal{D}}\mathcal{D}.$$

Then Δ is a formally self-adjoint elliptic operator. To understand Δ we introduce two “real” operators on $\mathbb{C}l(M)$:

$$(2.6) \quad D = \sum_j \{e_j \nabla_{e_j} + (Je_j) \nabla_{Je_j}\}, \quad D^c = \sum_j \{e_j \nabla_{Je_j} - (Je_j) \nabla_{e_j}\}.$$

The first operator is called the *Dirac operator*. Then we can easily see that

$$(2.7) \quad \mathcal{D} = \frac{1}{4}(D + iD^c), \quad \bar{\mathcal{D}} = \frac{1}{4}(D - iD^c).$$

Since $\mathcal{D}^2 = 0$, we have that $D^2 = (D^c)^2$ and $DD^c + D^cD = 0$. It follows that

$$(2.7) \quad \Delta = \frac{1}{4}D^2.$$

Since D is essentially self-adjoint, we have

$$(2.8) \quad \text{Ker } D = \text{Ker } D^2 = \text{Ker } \Delta.$$

Now, we consider the usual inner product

$$(2.9) \quad \ll \varphi_1, \varphi_2 \gg = \int_M \langle \varphi_1, \varphi_2 \rangle$$

for any $\varphi_1, \varphi_2 \in \Gamma_{cpt} \mathcal{Cl}(M)$. Let $L^2(\mathbb{C}l^{p,q}(M))$ be the completion of $\Gamma_{cpt} \mathbb{C}l^{p,q}$ with respect to \ll, \gg . We recall that the operators \mathcal{D} and $\bar{\mathcal{D}}$ are formal adjoint to one another with respect to \ll, \gg . Then \mathcal{D} and $\bar{\mathcal{D}}$ have closed extensions in $L^2(\mathbb{C}l^{p,q}(M))$ defined by

$$(2.10) \quad \mathcal{D}_{\max} := (\bar{\mathcal{D}}_{\min})^*, \quad \bar{\mathcal{D}}_{\max} := (\mathcal{D}_{\min})^*,$$

where $\bar{\mathcal{D}}_{\min}$ (resp. \mathcal{D}_{\min}) is a minimal extension of $\bar{\mathcal{D}}$ (resp. \mathcal{D}) and $(\)^*$ is the adjoint operator of $(\)$ with respect to \ll, \gg . Since Δ and D are essentially self-adjoints, we have $\mathcal{D}_{\max} = \mathcal{D}_{\min}$ and $\bar{\mathcal{D}}_{\max} = \bar{\mathcal{D}}_{\min}$ ([3]). And hence we denote the closed extensions as the same symbols. Consequently, from Lemma 2.3 and Theorem 2.4, we obtain the Hilbert complexes

$$(2.11) \quad \begin{aligned} & \dots \xrightarrow{\mathcal{D}} L^2(\mathbb{C}l^{p-1,q-1}(M)) \xrightarrow{\mathcal{D}} L^2(\mathbb{C}l^{p,q}(M)) \xrightarrow{\mathcal{D}} L^2(\mathbb{C}l^{p+1,q+1}(M)) \xrightarrow{\mathcal{D}} \dots, \\ & \dots \xleftarrow{\bar{\mathcal{D}}} L^2(\mathbb{C}l^{p-1,q-1}(M)) \xleftarrow{\bar{\mathcal{D}}} L^2(\mathbb{C}l^{p,q}(M)) \xleftarrow{\bar{\mathcal{D}}} L^2(\mathbb{C}l^{p+1,q+1}(M)) \xleftarrow{\bar{\mathcal{D}}} \dots. \end{aligned}$$

Now, we put

$$(2.12) \quad L^2\mathcal{H}^{p,q} := Ker\mathcal{D} / \overline{Im\bar{\mathcal{D}}} \cap L^2(\mathbb{C}l^{p,q}(M)),$$

$$(2.13) \quad L^2\hat{\mathcal{H}}^{p,q} := Ker\mathcal{D} \cap Ker\bar{\mathcal{D}} \cap L^2(\mathbb{C}l^{p,q}(M)),$$

$$(2.14) \quad L^2H^{p,q} := Ker\Delta \cap L^2(\mathbb{C}l^{p,q}(M)).$$

Here $L^2\mathcal{H}^{p,q}$ and $L^2H^{p,q}$ are called the *Clifford L^2 -cohomology group* and *L^2 -harmonic space*, respectively. Then we have

COROLLARY 2.5. *Let M be a complete Kähler manifold. Then we have*

$$L^2(\mathbb{C}l^{p,q}(M)) = L^2\hat{\mathcal{H}}^{p,q} \oplus \overline{Im\bar{\mathcal{D}}} \oplus \overline{Im\mathcal{D}},$$

and

$$L^2\mathcal{H}^{p,q} \cong L^2\hat{\mathcal{H}}^{p,q} \cong L^2H^{p,q}.$$

Proof. The first follows from Lemma 2.3 and Lemma 2.4. The second is obvious from [3, Lemma 3.2].

REMARK ([10]). We study the relationship between Dolbeault cohomology and Clifford cohomology. First, we prepare the some facts: Let $\Lambda^{r,s}(M)$ be the standard Dolbeault decomposition of $\Lambda^*(M) \otimes \mathbb{C}$. Then there are operators

$$\partial : \Gamma\Lambda^{r,s} \longrightarrow \Gamma\Lambda^{r+1,s}, \quad \bar{\partial} : \Gamma\Lambda^{r,s} \longrightarrow \Gamma\Lambda^{r,s+1}$$

given by the formulas;

$$(2.15) \quad \partial = \sum_j \bar{\xi}_j \wedge \nabla_{\xi_j}, \quad \bar{\partial} = \sum_j \xi_j \wedge \nabla_{\bar{\xi}_j},$$

where ∇ is the Kähler connection and $\{\xi_j, \bar{\xi}_j\}$ is as before. The formal adjoints of ∂ and $\bar{\partial}$ are given respectively by

$$(2.16) \quad \partial^* = - \sum_j i(\xi_j) \nabla_{\bar{\xi}_j}, \quad \bar{\partial}^* = - \sum_j i(\bar{\xi}_j) \nabla_{\xi_j},$$

where $i(\cdot)$ denotes the interior product. It is well known that under the isomorphism $\mathbb{C}l(M) \cong \Lambda^*(M) \otimes \mathbb{C}$, we have $\mathcal{D} \cong \bar{\partial} + \partial^*$ and $\bar{\mathcal{D}} \cong \partial + \bar{\partial}^*$. Note that the (p, q) -decomposition of $\mathbb{C}l(M)$ constructed above does not directly correspond to the Dolbeault decomposition. In fact,

$$(2.17) \quad \mathbb{C}l^{p,*}(M) \cong \bigoplus_{s-r=p} \Lambda^{r,s}(M),$$

where $\mathbb{C}l^{p,*}(M) = \bigoplus_q \mathbb{C}l^{p,q}(M)$. Moreover,

$$(2.18) \quad H^{s-r,n-r-s}(M) \cong H_{Dol}^{r,s}(M)$$

where $H_{Dol}^{r,s}(M) = H \cap \Lambda^{r,s}(M)$, H is the harmonic space. The relations (2.17) and (2.18) hold for the space of L^2 sections.

3. Vanishing theorems

In this section, we shall prove some vanishing theorems under various curvature conditions. Let M be a Kähler manifold and consider a hermitian vector bundle $S \rightarrow M$ of left modules over $\mathbb{C}l(M)$ with a hermitian metric $\langle \cdot, \cdot \rangle$ such that:

(1) Module multiplication by unit tangent vectors is unitary, i.e.,

$$(3.1) \quad \langle \xi \cdot \phi, \psi \rangle + \langle \phi, \bar{\xi} \cdot \psi \rangle = 0,$$

for any $\phi, \psi \in \Gamma(S)$ and $\xi \in \Gamma(TM) \otimes \mathbb{C}$

(2) With respect to the canonical hermitian connection, covariant differentiation is a derivation of module multiplication. That is, for $\phi \in \Gamma(\mathbb{C}l(M))$ and $s \in \Gamma(S)$, we have

$$(3.2) \quad \nabla(\phi \cdot s) = (\nabla\phi) \cdot s + \phi \cdot (\nabla s).$$

Now, we recall some basic results from [10]. For each j , we set $\omega_j = -\xi_j \bar{\xi}_j$, $\bar{\omega}_j = -\bar{\xi}_j \xi_j$. To each (possibly empty) subset $I = \{i_1, \dots, i_p\} \subseteq \{1, \dots, n\}$ with complementary subset $\{j_1, \dots, j_{n-p}\}$ we set $\omega_I = \omega_{i_1} \cdots \omega_{i_p} \bar{\omega}_{j_1} \cdots \bar{\omega}_{j_{n-p}}$ and we denote $|I| = p$. Then we have

$$(3.3) \quad 1 = \prod_{j=1}^n (\omega_j + \bar{\omega}_j) = \sum_{r=1}^n \pi_r,$$

where $\pi_r = \sum_{|I|=r} \omega_I$. Moreover, we have an orthogonal decomposition of the bundle

$$(3.4) \quad S = \bigoplus_{r=0}^n S^r, \quad S^r = \pi_r \cdot S.$$

Then the complex

$$(3.5) \quad 0 \rightarrow \Gamma_{cpt}(S^0) \xrightarrow{\mathcal{D}} \Gamma_{cpt}(S^1) \xrightarrow{\mathcal{D}} \cdots \xrightarrow{\mathcal{D}} \Gamma_{cpt}(S^n) \rightarrow 0$$

is elliptic. By Lemma 2.3, its completion becomes a Hilbert complex. Similarly with Corollary 2.5, we have

$$(3.6) \quad L^2\mathcal{H}^r(M, S) \cong L^2\hat{\mathcal{H}}^r(M, S) \cong L^2H^r(M, S).$$

Now, we define invariant operators on $\Gamma(S)$ by

$$(3.7) \quad \begin{aligned} \nabla^* \nabla &= - \sum_j \nabla_{\xi_j, \bar{\xi}_j}, & \bar{\nabla}^* \bar{\nabla} &= - \sum_j \nabla_{\bar{\xi}_j, \xi_j}, \\ \mathcal{R} &= \sum_{j,k} \xi_j \bar{\xi}_k R_{\bar{\xi}_j, \xi_k}, & \bar{\mathcal{R}} &= \sum_{j,k} \bar{\xi}_j \xi_k R_{\xi_j, \bar{\xi}_k}, \end{aligned}$$

where $R_{V,W} = \nabla_{V,W} - \nabla_{W,V}$ is the curvature tensor and where $\nabla_{V,W} = \nabla_V \nabla_W - \nabla_{\nabla_V W}$ is the invariant second covariant derivative. Then we have the following result([10]):

PROPOSITION 3.1. For any two sections $s_1, s_2 \in \Gamma(S)$, at least one of which has compact support, the following holds:

$$\int_M \langle \nabla^* \nabla s_1, s_2 \rangle = \int_M \langle \nabla s_1, \nabla s_2 \rangle,$$

where $\langle \nabla s_1, \nabla s_2 \rangle = \langle \nabla_{\bar{\xi}_i} s_1, \nabla_{\xi_i} s_2 \rangle$. Hence $\nabla^* \nabla$ is a formally self adjoint, nonnegative operator. Similarly, this holds for $\bar{\nabla}^* \bar{\nabla}$. Moreover, the zero order operators \mathcal{R} and $\bar{\mathcal{R}}$ are self-adjoint.

Moreover, by the straight calculation, we obtain the Bochner-Weitzenböck type formula([10]);

$$(3.8) \quad D\bar{D} + \bar{D}D = \nabla^* \nabla + \mathcal{R} = \bar{\nabla}^* \bar{\nabla} + \bar{\mathcal{R}}.$$

From this formula, we obtain the first important consequence

THEOREM 3.2. For any $s \in \text{dom}D\bar{D} \cap \text{dom}\bar{D}D$, we have

$$(3.9) \quad \|Ds\|^2 + \|\bar{D}s\|^2 = \|\nabla s\|^2 + \langle\langle \mathcal{R}s, s \rangle\rangle = \|\bar{\nabla}s\|^2 + \langle\langle \bar{\mathcal{R}}s, s \rangle\rangle,$$

where $\|\nabla s\|^2 = \langle\langle \nabla_{\bar{\xi}_j} s, \nabla_{\xi_j} s \rangle\rangle$ and $\|\bar{\nabla}s\|^2 = \langle\langle \nabla_{\xi_j} s, \nabla_{\bar{\xi}_j} s \rangle\rangle$.

Proof. First we consider a function ω_ℓ such that $0 \leq \omega_\ell(x) \leq 1$ for any $x \in M$, $\text{supp } \omega_\ell \subset B(x_0, 2\ell)$, $\omega_\ell(x) = 1$ for any $x \in B(x_0, \ell)$, $\lim_{\ell \rightarrow \infty} \omega_\ell = 1$ and $|d\omega_\ell| \leq C/\ell$ almost everywhere on M , where C is a positive constant independent of $\ell \in \mathbb{R}_+$, $x_0 \in M$ and $B(x_0, r)$ is the Riemannian open ball with radius r and center x_0 .

For any $s \in L^2(S)$, we calculate $\langle\langle D\bar{D}s + \bar{D}Ds, \omega_\ell^2 s \rangle\rangle$ on $B(2\ell)$. We choose $\{\xi_j, \bar{\xi}_j\}$ such that $(\nabla \xi_j)_x = (\nabla \bar{\xi}_j)_x = 0$. By the definition of \bar{D} and (3.2), we get

$$\langle\langle D\bar{D}s, \omega_\ell^2 s \rangle\rangle = 2 \langle\langle \omega_\ell \bar{D}s, \bar{\xi}_j (\nabla_{\xi_j} \omega_\ell) s \rangle\rangle + \|\omega_\ell \bar{D}s\|^2.$$

Using (3.1) and $\xi \bar{\xi} + \bar{\xi} \xi = -\|\xi\|^2$, we obtain

$$\|\xi \cdot s\|^2 + \|\bar{\xi} \cdot s\|^2 = \|\xi\|^2 \|s\|^2.$$

Hence we get $\|\xi \cdot s\| \leq \|\xi\| \|s\|$ for any $\xi \in TM \otimes \mathbb{C}$. Therefore, by this inequality and Schwarz inequality, we have

$$\begin{aligned} |\langle\langle \bar{D}s, \bar{\xi}_j (\nabla_{\xi_j} \omega_\ell) s \rangle\rangle| &\leq \|\bar{D}s\| \|\bar{\xi}_j (\nabla_{\xi_j} \omega_\ell) s\| \leq \|\bar{D}s\| \|(\nabla_{\xi_j} \omega_\ell) s\| \\ &\leq \|\bar{D}s\| \|\nabla_{\xi_j} \omega_\ell\| \|s\| \leq \frac{C}{\ell} \|\bar{D}s\| \|s\|. \end{aligned}$$

Since $\|s\|$ and $\|\bar{\mathcal{D}}s\|$ are finite, letting $\ell \rightarrow \infty$, we have $\ll \bar{\mathcal{D}}s, \bar{\xi}_j(\nabla_{\xi_j} \omega_\ell)s \gg \rightarrow 0$. This implies that $\ll \mathcal{D}\bar{\mathcal{D}}s, s \gg = \|\bar{\mathcal{D}}s\|^2$. Similarly, we get $\ll \bar{\mathcal{D}}\mathcal{D}s, s \gg = \|\mathcal{D}s\|^2$. On the other hand, by Proposition 3.1 and (3.2), we have

$$\ll \nabla^* \nabla s, \omega_\ell^2 s \gg = 2 \ll \omega_\ell \nabla s, \nabla \omega_\ell \cdot s \gg + \|\omega_\ell \nabla s\|^2.$$

By similar method, we have $|\ll \omega_\ell \nabla s, \nabla \omega_\ell \cdot s \gg| \rightarrow 0$ as $\ell \rightarrow \infty$. Hence we have $\ll \nabla^* \nabla s, s \gg = \|\nabla s\|^2$. Hence we complete the proof of the first equation of (3.9). For the second part, the proof is similar. \square

From Theorem 3.2, we have

$$2(\|\mathcal{D}s\|^2 + \|\bar{\mathcal{D}}s\|^2) = \|\nabla s\|^2 + \|\bar{\nabla}s\|^2 + \ll (\mathcal{R} + \bar{\mathcal{R}})s, s \gg.$$

Hence for any $s \in Ker\mathcal{D} \cap Ker\bar{\mathcal{D}}$, if $R = \mathcal{R} + \bar{\mathcal{R}}$ is non-negative, then we have $\|\nabla s\| = \|\bar{\nabla}s\| = 0$. This implies that s is a parallel section. In addition, if R is positive at some point, then $s = 0$. Hence we have

THEOREM 3.3. *Let M be a complete Kähler manifold and let S be any hermitian vector bundle of modules over $Cl(M)$. If R is non-negative and positive at some point of M , then the Clifford L^2 -cohomology group is trivial. This is,*

$$L^2\mathcal{H}^r(M, S) = \{0\}, \quad \text{for any } r = 0, 1, \dots, n.$$

Moreover, on $TM \subset Cl(M)$, we have ([8])

$$\mathcal{R} + \bar{\mathcal{R}} = \frac{1}{2} Ric.$$

Thus, from (3.6) and Theorem 3.3. we have

COROLLARY 3.4. *On the complete Kähler manifold, if the Ricci curvature is non-negative and positive at some point, then every L^2 -harmonic 1-form is necessary zero.*

Now, we shall consider some special cases of Theorem 3.3. To begin, we suppose that M is a Kähler spin manifold, i.e., we assume that

there exists a principal Spin-bundle, $P_{Spin}(M) \rightarrow M$, with a $Spin_{2n}$ -equivalent map $\tau : P_{Spin}(M) \rightarrow P_{SO}(M)$, to the bundle of real oriented orthonormal frame on M . The *bundle of spinors*, S , is then defined to be vector bundle associated to the unitary representation Δ of $Spin_{2n}$ given by the unique irreducible complex representation of Cl_{2n} , i.e., $S = P_{Spin} \times_{\Delta} \mathbb{C}^{2^n}$. This bundle is naturally a bundle of modules over $Cl(M)$ and carries a canonical connection induced from the lift of the riemannian connection on $P_{SO}(M)$. Since M is Kähler, this bundle S is naturally holomorphic and its connection is hermitian. On this bundle S , the curvature tensor R^S is given by

$$(3.10) \quad R_{V,W}^S = \frac{1}{4} \sum_{\alpha,\beta=1}^{2n} \langle R_{V,W} X_{\alpha}, X_{\beta} \rangle X_{\alpha} X_{\beta},$$

where X_1, \dots, X_{2n} is any real orthonormal basis of the tangent space. Choosing a basis $e_1, \dots, J e_n$, we can write R^S as

$$R_{V,W}^S = 2 \sum_{j,k=1}^n \langle R_{V,W} \xi_j, \bar{\xi}_k \rangle \bar{\xi}_j \xi_k + \sum_{j=1}^n \langle R_{V,W} \xi_j, \bar{\xi}_j \rangle .$$

Hence we have

$$(3.11) \quad \begin{aligned} \mathcal{R}^S &= \sum_{j,k=1}^n \xi_j \bar{\xi}_k R_{\xi_j, \xi_k}^S \\ &= \sum_{i,j,k=1}^n \langle R_{\xi_i, \bar{\xi}_i} \bar{\xi}_j, \xi_k \rangle \xi_j \bar{\xi}_k \\ &= -\frac{1}{2} \sum_{j,k=1}^n Ric(\bar{\xi}_j, \xi_k) \xi_j \bar{\xi}_k, \end{aligned}$$

where Ric is Ricci tensor on M ([10]). Since Ric is hermitian symmetric, we may choose our basis so that $Ric(\bar{\xi}_j, \xi_k) = 1/2 \lambda_j \delta_{jk}$, where $\lambda_j = Ric(e_j, e_j) = Ric(J e_j, J e_j)$, for $j = 1, \dots, n$, are the eigenvalues. Then we have

$$(3.12) \quad D\bar{D} + \bar{D}D = \nabla^* \nabla + \frac{1}{4} \sum_{j=1}^n \lambda_j \omega_j = \bar{\nabla}^* \bar{\nabla} + \frac{1}{4} \sum_{j=1}^n \lambda_j \bar{\omega}_j.$$

We note that $\nabla^* \nabla + \bar{\nabla}^* \bar{\nabla} = \frac{1}{2} \tilde{\nabla}^* \tilde{\nabla}$ where

$$(3.13) \quad \tilde{\nabla}^* \tilde{\nabla} = - \sum_j (\nabla_{e_j, e_j} + \nabla_{Je_j, Je_j})$$

is a self-adjoint, elliptic operator whose kernel is the space of parallel sections. We note that the scalar curvature κ of M is given by

$$(3.14) \quad \kappa = \text{trace}_R(\text{Ric}) = 2 \sum_j \lambda_j.$$

Hence we get

THEOREM 3.5([10]). *On the spinor bundle S , we have*

$$4(\mathcal{D}\bar{\mathcal{D}} + \bar{\mathcal{D}}\mathcal{D}) = \tilde{\nabla}^* \tilde{\nabla} + \frac{1}{4} \kappa,$$

where κ is the scalar curvature of M .

Summing up Theorem 3.3 and Theorem 3.5, we have

THEOREM 3.6. *Let M be a complete Kähler spin manifold. If $\kappa \geq 0$ for all $x \in M$ and $\kappa > 0$ for some point $x_0 \in M$, then there are no non-trivial L^2 -harmonic spinors.*

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