

On the Relationship between Two Connections in n -dimensional $*g$ -UFT

By Hyun Woo Lee

Yonsei University, Seoul, Korea

I. Introduction

A. n -dimensional $*g^{\lambda\nu}$ -unified field theory

In the usual Einstein's unified field theory the generalized n -dimensional Riemannian space X_n is endowed with a real non-symmetric tensor $g_{\lambda\mu}$ which may be split into its symmetric part $h_{\lambda\mu}$ and its skew-symmetric part $k_{\lambda\mu}$:

$$(1.1) \quad g_{\lambda\mu} = h_{\lambda\mu} + k_{\lambda\mu},$$

where

$$\text{Det}(g_{\lambda\mu}) \neq 0, \quad \text{Det}(h_{\lambda\mu}) \neq 0.$$

On the other hand, Einstein's $*g^{\lambda\nu}$ -unified field theory ($*g$ -UFT) in the same space X_n is defined to be based upon the basic real tensor $*g^{\lambda\nu}$ defined by

$$(1.2) \quad g_{\lambda\mu} *g^{\lambda\nu} = \delta_{\mu}^{\nu}.$$

It may also be decomposed into its symmetric part $*h^{\lambda\nu}$ and its skew-symmetric part $*k^{\lambda\nu}$:

$$(1.3) \quad *g^{\lambda\nu} = *h^{\lambda\nu} + *k^{\lambda\nu}.$$

Since $\text{Det}(*h^{\lambda\nu}) \neq 0$, we may define an unique tensor $*h_{\lambda\mu}$ by

$$(1.4) \quad *h_{\lambda\mu} *h^{\lambda\nu} = \delta_{\mu}^{\nu}.$$

In our $*g$ -UFT we use both $*h_{\lambda\mu}$ and $*h^{\lambda\nu}$, instead of $h_{\lambda\mu}$ and $h^{\lambda\nu}$, as the tensors for raising and/or lowering indices of all tensors defined in X_n in the usual manner, with the exception of the tensors $g_{\lambda\mu}$, $h_{\lambda\mu}$, and $k_{\lambda\mu}$ in order to avoid the notational confusion. We then have, for example,

$$(1.5) \quad a \quad *k_{\lambda\mu} = *k^{\rho\sigma} *h_{\lambda\rho} *h_{\mu\sigma}, \quad *g_{\lambda\mu} = *g^{\rho\sigma} *h_{\lambda\rho} *h_{\mu\sigma},$$

so that

$$(1.5) \quad b \quad *g_{\lambda\mu} = *h_{\lambda\mu} + *k_{\lambda\mu}.$$

The basic connection $\Gamma_{\lambda\mu}^{\nu}$ in our $*g$ -UFT is given by the following Einstein's equation expressed in terms of $*g^{\lambda\nu}$:

$$(1.6) \quad a \quad \partial_{\sigma} *g^{\lambda\mu} + \Gamma_{\sigma\alpha}^{\lambda} *g^{\alpha\mu} + \Gamma_{\sigma\alpha}^{\mu} *g^{\lambda\alpha} = 0,$$

or equivalently

$$(1.6) \quad b \quad D_{\sigma} *g^{\lambda\mu} = -2S_{\sigma\alpha}^{\mu} *g^{\lambda\alpha},$$

where D_{σ} is the symbolic vector of the covariant derivative with respect to $\Gamma_{\lambda\mu}^{\nu}$ and

$$(1.7) \quad S_{\lambda\mu}^{\nu} (*g^{\lambda\sigma}) \triangleq \Gamma_{(\lambda\mu)}^{\nu}.$$

For the convenience of writing we shall make the following agreement in our further considerations:

Agreement (1.1). We omit the notation “*” from all starred tensors and operators in the remaining of the present paper, and understand that they are starred.

B. The problem to be solved

As given in (1.3) and (1.6), *g-UFT is based on a real tensor

$$(1.3) \quad g^{\lambda\nu} = h^{\lambda\nu} + k^{\lambda\nu}$$

and a real connection $\Gamma_{\lambda\mu}^{\nu}$ which satisfies the field equation

$$(1.6) \quad \partial_{\alpha} g^{\lambda\mu} + \Gamma_{\alpha\sigma}^{\lambda} g^{\sigma\mu} + \Gamma_{\alpha\sigma}^{\mu} g^{\lambda\sigma} = 0.$$

Consider another unified field theory based on a real tensor

$$(1.8) \quad \bar{g}^{\lambda\nu} \stackrel{\text{df}}{=} \bar{h}^{\lambda\nu} + \bar{k}^{\lambda\nu} \stackrel{\text{df}}{=} h^{\lambda\nu} - k^{\lambda\nu}; \quad \bar{h}^{\lambda\nu} = \bar{h}^{(\lambda\nu)}, \quad \bar{k}^{\lambda\nu} = \bar{k}^{(\lambda\nu)},$$

and a real connection $\bar{\Gamma}_{\lambda\mu}^{\nu}$ which satisfies

$$(1.9) \quad \partial_{\alpha} \bar{g}^{\lambda\mu} + \bar{\Gamma}_{\alpha\sigma}^{\lambda} \bar{g}^{\sigma\mu} + \bar{\Gamma}_{\alpha\sigma}^{\mu} \bar{g}^{\lambda\sigma} = 0.$$

The main purpose of the present paper is to find the relationship between $\Gamma_{\lambda\mu}^{\nu}$ given by (1.6) and $\bar{\Gamma}_{\lambda\mu}^{\nu}$ given by (1.9).

II. Preliminary results

In this chapter some results which are essential to the present paper will be briefly stated without proofs. The detailed proofs are given in [1].

The following tensors will be used in our further considerations:

$$(2.1) \text{ a} \quad {}^{(0)}k_{\lambda}^{\nu} \stackrel{\text{df}}{=} \delta_{\lambda}^{\nu}, \quad {}^{(1)}k_{\lambda}^{\nu} \stackrel{\text{df}}{=} k_{\lambda}^{\nu}, \quad {}^{(\rho)}k_{\lambda}^{\nu} \stackrel{\text{df}}{=} {}^{(\rho-1)}k_{\lambda}^{\alpha} k_{\alpha}^{\nu}.$$

$$(2.1) \text{ b} \quad A_{\alpha\mu\nu}^{\beta\gamma} \stackrel{\text{df}}{=} {}^{(\rho)}k_{\alpha}^{\alpha} {}^{(\rho)}k_{\mu}^{\beta} {}^{(\rho)}k_{\nu}^{\gamma},$$

$$(2.1) \text{ c} \quad 2 A_{\alpha\mu\nu}^{\beta\gamma} \stackrel{\text{df}}{=} A_{\alpha\mu\nu}^{\beta\gamma} + A_{\alpha\mu\nu}^{\beta\gamma}, \quad 2 A_{\alpha\mu\nu}^{\beta\gamma} \stackrel{\text{df}}{=} A_{\alpha\mu\nu}^{\beta\gamma} - A_{\alpha\mu\nu}^{\beta\gamma}.$$

Also, denoting the tensor $T_{\alpha\mu\nu}$ by T, the following abbreviations will be used:

$$(2.2) \text{ a} \quad T \stackrel{\text{df}}{=} T_{\alpha\mu\nu} \stackrel{\text{df}}{=} A_{\alpha\mu\nu}^{\beta\gamma} T_{\alpha\beta\gamma},$$

$$(2.2) \text{ b} \quad T \stackrel{\text{df}}{=} T_{\alpha\mu\nu} \stackrel{\text{df}}{=} T_{\alpha\mu\nu}^{\text{000}}.$$

The connection $\Gamma_{\lambda\mu}^{\nu}$ defined by (1.6) is given by

$$(2.3) \quad \Gamma_{\lambda\mu}^{\nu} = \left\{ \begin{matrix} \nu \\ \lambda\mu \end{matrix} \right\} + S_{\lambda\mu}^{\nu} + U_{\lambda\mu}^{\nu}$$

where $\left\{ \begin{matrix} \nu \\ \lambda\mu \end{matrix} \right\}$ are the Christoffel symbols defined by $h_{\lambda\mu}$ in the usual way and

$$(2.4) \quad U_{\nu\lambda\mu} \stackrel{\text{df}}{=} S_{(\lambda\mu)\nu}^{\text{(10)0}} + 2 S_{\nu(\lambda\mu)}^{\text{(10)0}}.$$

The system of equations (1.6) is equivalent to

$$(2.5) \quad S_{\alpha\beta\gamma} X_{\alpha\mu\nu}^{\beta\gamma} = B_{\alpha\beta\gamma},$$

where

$$(2.6) \quad X_{\alpha\mu\nu}^{\beta\gamma} \stackrel{\text{df}}{=} A_{(\alpha\mu)\nu}^{\text{000}} + A_{(\alpha\mu)\nu}^{\text{110}} + 2 A_{(\alpha\mu)\nu}^{\text{(10)1}}$$

$$(2.7) \quad B_{\alpha\mu\nu} \stackrel{\text{df}}{=} \frac{1}{2} (K_{\alpha\mu\nu} + 3 K_{\alpha(\mu\beta} k_{\nu)}^{\beta}),$$

$$(2.8) \quad K_{\alpha\mu\nu} \stackrel{\text{df}}{=} \nabla_{\alpha} k_{\lambda\mu} + \nabla_{\mu} k_{\alpha\lambda} + \nabla_{\lambda} k_{\alpha\mu}.$$

Here ∇_{α} is the symbolic vector of the covariant derivative with respect to $\left\{ \begin{matrix} \nu \\ \lambda\mu \end{matrix} \right\}$.

III. Main theorem

The first consequences of (1.3), (1.8), (2.6), (2.7), and (2.8) are formally stated in the

following two lemmas:

Lemma (3.1). *We have*

$$\begin{aligned} (3.1) \ a & \quad \bar{h}^{\lambda\nu} = \bar{h}^{\lambda\nu}, \quad \bar{h}_{\lambda\mu} = h_{\lambda\mu}, \\ (3.1) \ b & \quad \bar{k}^{\lambda\nu} = -k^{\lambda\nu}, \quad \bar{k}_{\lambda\mu} = -k_{\lambda\mu}, \quad \bar{k}_{\lambda}{}^{\nu} = -k_{\lambda}{}^{\nu}, \\ (3.1) \ c & \quad \bar{g}^{\lambda\nu} = g^{\nu\lambda}, \quad \bar{g}_{\lambda\mu} = g_{\mu\lambda}. \end{aligned}$$

Lemma(3.2). *We have*

$$\begin{aligned} (3.2) \ a & \quad \bar{K}_{\omega\rho\mu} = -K_{\omega\rho\mu}, \\ (3.2) \ b & \quad \bar{B}_{\omega\rho\nu} = -B_{\omega\rho\nu}, \\ (3.2) \ c & \quad X_{\omega\rho\nu}^{\alpha\beta\gamma} = X_{\omega\rho\nu}^{\alpha\beta\gamma}. \end{aligned}$$

Proof. (3.1) are direct consequences of the definition of $\bar{g}^{\lambda\mu}$. (3.2) follow from (2.6), (2.7), and (2.8) using (2.1) and (3.1)b.

Finally, we have the following main theorem which states that the connection $\bar{\Gamma}_{\lambda\mu}^{\nu}$ is pseudo-Hermitian symmetric in the indices λ, μ .

Theorem (3.3). *If the system (1.6) admits unique solution, then we have the following relation:*

$$(3.3) \quad \bar{\Gamma}_{\lambda\mu}^{\nu} = \Gamma_{\lambda\mu}^{\nu}.$$

Proof. Assume that (1.6) has a unique solution. Then $\Gamma_{\lambda\mu}^{\nu}$ is expressed as a function of $g^{\lambda\nu}$. Observing (1.6) and (1.9), we see that $\bar{\Gamma}_{\lambda\mu}^{\nu}$ is the same function of $\bar{g}^{\lambda\nu}$. Hence, comparing (2.5) with

$$(3.4) \quad \bar{S}_{\alpha\beta\gamma} X_{\omega\rho\nu}^{\alpha\beta\gamma} = \bar{B}_{\omega\rho\nu},$$

which is equivalent to (1.9), and using (3.2), we have

$$(3.5) \quad \bar{S}_{\omega\rho\nu} = -S_{\omega\rho\nu} = S_{\omega\rho\nu},$$

so that

$$(3.6) \quad \bar{U}_{\nu\lambda\mu} = U_{\nu\lambda\mu} = U_{\nu\mu\lambda}$$

according to (2.1), (2.2), (2.4), (3.1), and (3.5). On the other hand, we have

$$(3.7) \quad \left\{ \begin{array}{c} \bar{\nu} \\ \bar{\lambda\mu} \end{array} \right\} = \left\{ \begin{array}{c} \nu \\ \lambda\mu \end{array} \right\} = \left\{ \begin{array}{c} \nu \\ \mu\lambda \end{array} \right\}$$

in virtue of (3.1)a. Hence, according to (2.3),

$$\bar{\Gamma}_{\lambda\mu}^{\nu} = \left\{ \begin{array}{c} \bar{\nu} \\ \bar{\lambda\mu} \end{array} \right\} + \bar{S}_{\lambda\mu}{}^{\nu} + \bar{U}_{\lambda\mu}{}^{\nu} = \left\{ \begin{array}{c} \nu \\ \mu\lambda \end{array} \right\} + S_{\mu\lambda}{}^{\nu} + U_{\mu\lambda}{}^{\nu} = \Gamma_{\mu\lambda}^{\nu}.$$

References

- [1] K. T. Chung (1963), Einstein's connection in terms of $*g^{\lambda\nu}$, *Il Nuovo Cimento*, (X)27, 1297-1324.
- [2] V. Hlavatý (1957), *Geometry of Einstein's unified field theory*, P. Noordhoff Ltd, Groningen.