

**ON ADJOINT COMPLEX STRUCTURES  
IN THE DIFFERENTIABLE MANIFOLD OF DIMENSION  $2n+1$**

By Jae Koo Ahn

**Introduction.**

In the present paper, we define a  $(2n+1)$ -dimensional differentiable manifold which is liken to a manifold constituted by the  $2n$ -dimensional complex manifold and one-dimensional real manifold adjoining to it, and, we study this manifold as an analogue of the complex manifold.

In § 1, we define an adjoint complex manifold befitting to the complex manifold, introduce a Riemannian metric to this and define more an adjoint Hermitian manifold. In § 2, we calculate the Christoffel symbol and define an adjoint Kaehlerian manifold befitting to the Kaehlerian manifold. And in § 3, we calculate the curvature tensor in the adjoint Kaehlerian manifold. In § 4, we define an almost adjoint complex manifold introduced by three structure tensors, and find the form of the structure tensors when this manifold induces to the adjoint complex manifold, and a condition that this manifold admits an adjoint complex manifold. In § 5, and § 6, we define an almost adjoint Hermitian manifold and an almost adjoint Kaehlerian manifold, and find the conditions that they admit the adjoint Hermitian manifold and an adjoint Kaehlerian manifold respectively.

**§ 1. Adjoint complex manifold.**

Let us consider a  $(2n+1)$ -dimensional differentiable manifold  $X_{2n+1}$  of class  $C^\infty$ , which is covered by the coordinate neighborhood system  $(x^i)$ , where  $i, j, k$  is taken by the indices  $1, 2, \dots, n; \bar{1}, \bar{2}, \dots, \bar{n}; \infty$ . For arbitrary point  $P$  of  $X_{2n+1}$ , let  $P$  belong to  $U \cap U'$ , then it holds

$$(1.1) \quad x^{i'} = x^{i'}(x), \quad \left| \frac{\partial x^{i'}}{\partial x^i} \right| \neq 0,$$

where  $U$  and  $U'$  are coordinate neighborhoods represented by  $x^i$  and  $x^{i'}$  and  $x^{i'}(x)$  is real analytical functions of  $x^i$ .

Let us put

$$(1.2) \quad z^\alpha = x^\alpha + ix^{\bar{\alpha}}, \quad z^{\bar{\alpha}} = x^\alpha - ix^{\bar{\alpha}}, \quad z^\infty = x^\infty,$$

where  $\alpha, \beta, \gamma$  and  $\bar{\alpha}, \bar{\beta}, \bar{\gamma}$  are taken by the indices  $1, 2, \dots, n$  and  $1, 2, \dots, \bar{n}$  respectively, then  $(x^\alpha, x^{\bar{\alpha}}, x^\infty)$  assigns to  $(z^\alpha, z^{\bar{\alpha}}, z^\infty)$ , and conversely by the equations

$$(1.3) \quad x^\alpha = \frac{z^\alpha + z^{\bar{\alpha}}}{2}, \quad x^{\bar{\alpha}} = \frac{z^\alpha - z^{\bar{\alpha}}}{2i}, \quad x^\infty = z^\infty,$$

Form this, we have a one-to-one correspondence between  $(x^i)$  and  $(z^\alpha, z^{\bar{\alpha}}, z^\infty)$ , and  $(z^\alpha, z^{\bar{\alpha}}, z^\infty)$  may be considered to be a coordinate of point  $P$  in  $X_{2n+1}$ . And the transformation (1.1) may be always represented by the equations

$$(1.4) \quad z^{\alpha'} = z^{\alpha'}(z^\alpha, z^{\bar{\alpha}}, z^\infty), \quad z^{\bar{\alpha}'} = z^{\bar{\alpha}'}(z^\alpha, z^{\bar{\alpha}}, z^\infty),$$

$$z^{\infty'} = z^{\infty'}(z^\alpha, z^{\bar{\alpha}}, z^\infty);$$

$$\begin{vmatrix} \frac{\partial z^{\alpha'}}{\partial z^\alpha} & \frac{\partial z^{\alpha'}}{\partial z^{\bar{\alpha}}} & \frac{\partial z^{\alpha'}}{\partial z^\infty} \\ \frac{\partial z^{\bar{\alpha}'}}{\partial z^\alpha} & \frac{\partial z^{\bar{\alpha}'}}{\partial z^{\bar{\alpha}}} & \frac{\partial z^{\bar{\alpha}'}}{\partial z^\infty} \\ \frac{\partial z^{\infty'}}{\partial z^\alpha} & \frac{\partial z^{\infty'}}{\partial z^{\bar{\alpha}}} & \frac{\partial z^{\infty'}}{\partial z^\infty} \end{vmatrix} \neq 0,$$

Now, if it is possible to choose a coordinate neighborhood system, in such that, in the domain  $U \cap U'$  of two coordinate neighborhoods  $U(z^i), U'(z^{i'})$ , it holds

$$(1.5) \quad z^{\alpha'} = z^{\alpha'}(z^\alpha), \quad z^{\bar{\alpha}'} = z^{\bar{\alpha}'}(z^{\bar{\alpha}}), \quad z^{\infty'} = z^{\infty'}(z^\infty),$$

$$\left| \frac{\partial z^{\alpha'}}{\partial z^\alpha} \right| \cdot \left| \frac{\partial z^{\bar{\alpha}'}}{\partial z^{\bar{\alpha}}} \right| \cdot \frac{\partial z^{\infty'}}{\partial z^\infty} \neq 0,$$

we say that the manifold  $X_{2n+1}$  admits an *adjoint complex structure* and we call  $X_{2n+1}$  an *adjoint complex manifold*.

From (1.5), since  $\left| \frac{\partial z^{\alpha'}}{\partial z^\alpha} \right| \cdot \left| \frac{\partial z^{\bar{\alpha}'}}{\partial z^{\bar{\alpha}}} \right| > 0$ , the manifold is orientable, if and only if  $\frac{\partial z^{\infty'}}{\partial z^\infty} > 0$ .

We consider the components of a tensor  $T_i^h$ , i. e.,

$$T_i^h = (T_\beta^\alpha, T_\beta^{\bar{\alpha}}, \dots, T_\infty^\infty).$$

For coordinate transformation (1.5), they hold

$$(1.6) \quad T_{i'}^{h'} = \frac{\partial z^{h'}}{\partial z^h} \frac{\partial z^i}{\partial z^{i'}} T_i^h$$

From this and (1.5),  $(T_{\beta}^{\alpha}, \dots, 0)$ ,  $(0, T_{\beta}^{\bar{\alpha}}, \dots, 0)$ ,  $\dots$ ,  $(0, \dots, 0, T_{\infty}^{\infty})$  are the components of tensors. And from (1.6), we get

$$(1.7) \quad \bar{T}_{i'}^{\bar{h}'} = \frac{\partial z^{h'}}{\partial z^h} \frac{\partial z^i}{\partial z^{i'}} \bar{T}_i^{\bar{h}}$$

where, if  $h$  takes the indices  $\alpha, \bar{\alpha}, \infty$ ,  $\bar{h}$  takes  $\bar{\alpha}, \alpha, \infty$  respectively, and thus, if  $T_i^h$  are components of tensor, then  $\bar{T}_i^{\bar{h}}$  are the components of the tensor. We call  $\bar{T}_i^{\bar{h}}$  the *adjoint conjugate* of  $T_i^h$ , and if  $T_i^h = \bar{T}_i^{\bar{h}}$ , we say that  $T_i^h$  is *adjoint self-conjugate*. If  $T_i^h$  is adjoint self-conjugate, then the component  $T_{\infty}^{\infty}$  is real.

Now, let us assume that our adjoint complex manifold admits a Riemannian metric

$$(1.8) \quad ds^2 = g_{ji} dz^j dz^i,$$

where symmetric tensor  $g_{ji}$  is adjoint self-conjugate and satisfies

$$(1.9) \quad g_{ji} = \begin{pmatrix} 0 & g_{\beta\bar{\alpha}} & 0 \\ g_{\bar{\beta}\alpha} & 0 & 0 \\ 0 & 0 & g_{\infty\infty} \end{pmatrix}.$$

Then the metric form (1.8) can be written in the form

$$(1.10) \quad ds^2 = 2g_{\beta\bar{\alpha}} dz^{\beta} dz^{\bar{\alpha}} + g_{\infty\infty} dz^{\infty} dz^{\infty}.$$

Of course, this is real since  $ds^2 = \overline{ds^2}$ . We call this metric satisfying (1.9) an *adjoint Hermitian metric* and the adjoint complex manifold with the adjoint Hermitian metric an *adjoint Hermitian manifold*.

## § 2. The Christoffel symbol of the adjoint Hermitian manifold.

In the adjoint Hermitian manifold  $X_{2n+1}$  with the metric (1.9) and (1.10), let us calculate the contravariant fundamental tensor  $g^{ji}$ . Then we have the components of  $g^{ji}$  as follows:

$$(2.1) \quad g^{ji} = \begin{pmatrix} 0 & g^{\beta\bar{\alpha}} & 0 \\ g^{\bar{\beta}\alpha} & 0 & 0 \\ 0 & 0 & g^{\infty\infty} \end{pmatrix}.$$

$$(2.2) \quad g^{\gamma\bar{\beta}} g_{\bar{\beta}\alpha} = \delta_{\alpha}^{\gamma}, \quad g^{\bar{\gamma}\beta} g_{\beta\bar{\alpha}} = \delta_{\bar{\alpha}}^{\bar{\gamma}}, \quad g^{\infty\infty} g_{\infty\infty} = 1,$$

and  $g^{ji}$  is adjoint self-conjugate.

Next, let us calculate the Christoffel symbol

$$(2.3) \quad \Gamma_{ji}^h = \frac{1}{2} g^{hr} (\partial_j g_{ir} + \partial_i g_{jr} - \partial_r g_{ji}).$$

Then we obtain the components as follows:

$$(2.4) \quad \begin{aligned} \Gamma_{\gamma\beta}^{\alpha} &= \frac{1}{2} g^{\alpha\bar{\tau}} (\partial_{\gamma} g_{\beta\bar{\tau}} + \partial_{\beta} g_{\gamma\bar{\tau}}), \\ \Gamma_{\bar{\gamma}\beta}^{\alpha} &= \frac{1}{2} g^{\alpha\bar{\tau}} (\partial_{\bar{\gamma}} g_{\beta\bar{\tau}} - \partial_{\bar{\tau}} g_{\bar{\gamma}\beta}), \\ \Gamma_{\bar{\gamma}\beta}^{\infty} &= -\frac{1}{2} g^{\infty\infty} \partial_{\infty} g_{\bar{\gamma}\beta}, & \Gamma_{\infty\beta}^{\alpha} &= \frac{1}{2} g^{\alpha\bar{\tau}} \partial_{\infty} g_{\beta\bar{\tau}}, \\ \Gamma_{\infty\infty}^{\alpha} &= \frac{1}{2} g^{\alpha\bar{\tau}} \partial_{\bar{\tau}} g_{\infty\infty}, & \Gamma_{\beta\infty}^{\infty} &= \partial_{\beta} \log \sqrt{g_{\infty\infty}}, \\ \Gamma_{\infty\infty}^{\infty} &= \partial_{\infty} \log \sqrt{g_{\infty\infty}}, & \Gamma_{\bar{\gamma}\beta}^{\alpha} &= \Gamma_{\infty\bar{\beta}}^{\alpha} = \Gamma_{\gamma\beta}^{\infty} = 0, \end{aligned}$$

and the values of the other components are given by symmetric and adjoint self-conjugate properties.

From the equation of the coordinate transformation

$$(2.5) \quad \Gamma_{j'i'}^{h'} = \frac{\partial z^{h'}}{\partial z^h} \left( \frac{\partial z^j}{\partial z^{j'}} \frac{\partial z^i}{\partial z^{i'}} \Gamma_{ji}^h + \frac{\partial^2 z^h}{\partial z^{j'} \partial z^{i'}} \right),$$

If the  $(h', i', j')$  takes of the indices excepting  $(\alpha', \beta', \gamma')$ ,  $(\bar{\alpha}', \bar{\beta}', \bar{\gamma}')$  and  $(\infty', \infty')$ , then, since it holds

$$\frac{\partial z^{h'}}{\partial z^h} \frac{\partial^2 z^h}{\partial z^{j'} \partial z^{i'}} = 0,$$

the condition

$$(2.6) \quad \Gamma_{ji}^h = 0, \text{ excepting } \Gamma_{\gamma\beta}^{\alpha}, \Gamma_{\bar{\gamma}\beta}^{\bar{\alpha}} \text{ and } \Gamma_{\infty\infty}^{\infty},$$

is invariant under the coordinate transformation (1.5). And from (2.4), (2.6) is equivalent to

$$(2.7) \quad \begin{aligned} \partial_{\gamma} g_{\beta\bar{\alpha}} &= \partial_{\beta} g_{\gamma\bar{\alpha}}, & (\partial_{\bar{\gamma}} g_{\beta\bar{\alpha}} &= \partial_{\bar{\beta}} g_{\bar{\gamma}\alpha}), \\ \partial_{\infty} g_{\beta\bar{\tau}} &= 0, \\ \partial_{\tau} g_{\infty\infty} &= 0, & (\partial_{\bar{\tau}} g_{\infty\infty} &= 0). \end{aligned}$$

The condition (2.6) or (2.7) is called *adjoint Kaehlerian condition*, and if an adjoint Hermitian manifold  $X_{2n+1}$  admits this adjoint Kaehlerian condition, we call this manifold  $X_{2n+1}$  an *adjoint Kaehlerian manifold*.

From (2.7), we obtain the following:

**THEOREM 1.** *On an adjoint Kaehlerian manifold, for the metric tensor, it holds that  $g_{\beta\bar{\alpha}}$  depend only upon  $z^\alpha$  and  $z^{\bar{\alpha}}$  and  $g_{\infty\infty}$  only upon  $z^\infty$ .*

On the adjoint Kaehlerian manifold, the Christoffel symbols are represented by

$$(2.8) \quad \Gamma_{\gamma\beta}^\alpha = g^{\alpha\bar{\tau}} \partial_\gamma g_{\beta\bar{\tau}}, \quad \Gamma_{\bar{\gamma}\bar{\beta}}^{\bar{\alpha}} = g^{\alpha\bar{\tau}} \partial_{\bar{\gamma}} g_{\beta\bar{\tau}}, \quad \Gamma_{\infty\infty}^\infty = \partial_\infty \log \sqrt{g_{\infty\infty}},$$

and the others are zero.

### § 3. The curvature tensor in the adjoint Kaehlerian manifold.

In the adjoint Kaehlerian manifold  $X_{2n+1}$ , we calculate the Riemannian curvature tensor

$$(3.1) \quad R_{kji}^h = \partial_i \Gamma_{ji}^h - \partial_j \Gamma_{ki}^h + \Gamma_{ks}^h \Gamma_{ji}^s - \Gamma_{js}^h \Gamma_{ki}^s.$$

From (2.8), we have

$$(3.2) \quad \begin{aligned} R_{kj\bar{\beta}}^\alpha &= 0, & R_{kj\beta}^{\bar{\alpha}} &= 0, \\ R_{kj\infty}^\alpha &= 0, & R_{kj\beta}^\infty &= 0, \\ R_{kj\infty}^{\bar{\alpha}} &= 0, & R_{kj\bar{\beta}}^\infty &= 0, \end{aligned}$$

and for

$$R_{kjih} = R_{kji}^s g_{sh},$$

we obtain

$$(3.3) \quad \begin{aligned} R_{kj\beta\alpha} &= 0, & R_{kj\bar{\beta}\bar{\alpha}} &= 0, \\ R_{kj\beta\infty} &= 0, & R_{kj\bar{\beta}\infty} &= 0. \end{aligned}$$

Also

$$R_{kjih} = R_{ihkj}$$

implies

$$(3.4) \quad \begin{aligned} R_{\beta\alpha kj} &= 0, & R_{\bar{\beta}\bar{\alpha} kj} &= 0, \\ R_{\beta\infty kj} &= 0, & R_{\bar{\beta}\infty kj} &= 0. \end{aligned}$$

And from (2.4) and THEOREM 1, we obtain

$$(3.5) \quad R_{\delta\infty\infty}^\infty = \partial_\delta \Gamma_{\infty\infty}^\infty = 0,$$

$$R_{\delta\infty\infty}^{\infty} = \partial_{\delta}\Gamma_{\infty\infty}^{\infty} = 0.$$

Then, from (3.2)~(3.5), we can see that only the components of the form

$$(3.6) \quad R_{\delta\gamma\beta}^{\alpha}, \quad R_{\delta\bar{\gamma}\bar{\beta}}^{\bar{\alpha}}, \quad R_{\delta\bar{\gamma}\beta}^{\alpha}, \quad R_{\delta\gamma\bar{\beta}}^{\alpha},$$

and

$$(3.7) \quad R_{\delta\bar{\gamma}\beta\alpha}, \quad R_{\delta\bar{\gamma}\bar{\beta}\alpha}, \quad R_{\delta\gamma\beta\alpha}, \quad R_{\delta\bar{\gamma}\bar{\beta}\alpha},$$

can be different from zero. And they are represented by

$$(3.8) \quad R_{\delta\gamma\beta}^{\alpha} = \partial_{\delta}\Gamma_{\gamma\beta}^{\alpha}, \quad R_{\delta\bar{\gamma}\bar{\beta}}^{\bar{\alpha}} = \partial_{\delta}\Gamma_{\bar{\gamma}\bar{\beta}}^{\bar{\alpha}}.$$

If  $\Gamma_{\gamma\beta}^{\alpha}$  depend only upon  $z^{\alpha}$ , then  $\Gamma_{\bar{\gamma}\bar{\beta}}^{\bar{\alpha}}$  depend only upon  $z^{\bar{\alpha}}$  since  $\Gamma_{\gamma\beta}^{\alpha}$  are adjoint self-conjugate, and thus we obtain  $R_{kji}^h = 0$ , and hence we have the following:

**THEOREM 2.** *In the adjoint Kaehlerian manifold, if  $\Gamma_{\gamma\beta}^{\alpha}$  (or  $\Gamma_{\bar{\gamma}\bar{\beta}}^{\bar{\alpha}}$ ) depend only upon  $z^{\alpha}$  (or  $z^{\bar{\alpha}}$ ), then the manifold is flat.*

For the Ricci tensor  $R_{ji}$ , from (3.6), we have

$$(3.9) \quad R_{\gamma\beta} = 0, \quad R_{\bar{\gamma}\bar{\beta}} = 0, \quad R_{\gamma\infty} = 0, \quad R_{\bar{\gamma}\infty} = 0, \quad R_{\infty\infty} = 0,$$

$$R_{\bar{\gamma}\beta} = R_{i\bar{\gamma}\beta}^i = R_{\alpha\bar{\gamma}\beta}^{\alpha} = -\partial_{\bar{\gamma}}\Gamma_{\alpha\beta}^{\alpha} = -\partial_{\bar{\gamma}}\partial_{\beta}\log\sqrt{g}.$$

And since it holds that

$$(3.10) \quad g = |g_{ji}| = g_{\infty\infty}|g_{\beta\bar{\alpha}}|^2,$$

and  $g_{\infty\infty}$  depends only upon  $z^{\infty}$ , we have

$$(3.11) \quad R_{\bar{\gamma}\beta} = -\partial_{\bar{\gamma}}\partial_{\beta}\log|g_{\beta\bar{\alpha}}|,$$

and the others are all zero.

#### § 4. The structure of almost adjoint complex manifold.

In an adjoint complex manifold, there exist three mixed tensor  $F_i^h$ ,  $G_i^h$  and  $H_i^h$  which have the numerical components

$$(4.1) \quad (F_i^h) = \begin{pmatrix} i\delta_{\beta}^{\alpha} & 0 & 0 \\ 0 & -i\delta_{\beta}^{\bar{\alpha}} & 0 \\ 0 & 0 & i \end{pmatrix}, \quad (G_i^h) = \begin{pmatrix} -i\delta_{\beta}^{\alpha} & 0 & 0 \\ 0 & i\delta_{\beta}^{\bar{\alpha}} & 0 \\ 0 & 0 & i \end{pmatrix},$$

$$(H_i^h) = \begin{pmatrix} i\delta_\beta^\alpha & 0 & 0 \\ 0 & i\delta_\beta^\alpha & 0 \\ 0 & 0 & -i \end{pmatrix}, \quad (i = \sqrt{-1})$$

in all coordinate system (1.5), and which satisfy

$$(4.2) \quad F_i^h F_j^i = -A_j^h, \quad G_i^h G_j^i = -A_j^h, \quad H_i^h H_j^i = -A_j^h$$

$$(4.3) \quad F_i^h + G_i^h + H_i^h = iA_i^h,$$

$$(4.4) \quad F_i^h G_j^i + G_i^h H_j^i + H_i^h F_j^i = A_j^h,$$

where  $A_j^h$  is unit tensor. From (4.3), we can see that, if two among three tensors in (4.1) are given, then another is determined by their two tensors.

Let us put

$$(4.5) \quad \begin{aligned} B_i^h &= \frac{1}{2}(A_i^h + iG_i^h), & C_i^h &= \frac{1}{2}(A_i^h + iF_i^h), \\ D_i^h &= \frac{1}{2}(A_i^h + iH_i^h), \\ \bar{B}_i^h &= \frac{1}{2}(A_i^h - iG_i^h), & \bar{C}_i^h &= \frac{1}{2}(A_i^h - iF_i^h), \\ \bar{D}_i^h &= \frac{1}{2}(A_i^h - iH_i^h). \end{aligned}$$

Then their components are represented by

$$(4.6) \quad \begin{aligned} (B_i^h) &= \begin{pmatrix} \delta_\beta^\alpha & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}, & (C_i^h) &= \begin{pmatrix} 0 & 0 & 0 \\ 0 & \delta_\beta^\alpha & 0 \\ 0 & 0 & 0 \end{pmatrix}, & (D_i^h) &= \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix}, \\ (\bar{B}_i^h) &= \begin{pmatrix} 0 & 0 & 0 \\ 0 & \delta_\beta^\alpha & 0 \\ 0 & 0 & 1 \end{pmatrix}, & (\bar{C}_i^h) &= \begin{pmatrix} \delta_\beta^\alpha & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix}, & (\bar{D}_i^h) &= \begin{pmatrix} \delta_\beta^\alpha & 0 & 0 \\ 0 & \delta_\beta^\alpha & 0 \\ 0 & 0 & 0 \end{pmatrix}. \end{aligned}$$

If, in a  $(2n+1)$ -dimensional differentiable manifold  $X_{2n+1}$  of class  $C^\infty$ , there exist three mixed tensors  $F_i^h$ ,  $G_i^h$  and  $H_i^h$  satisfying

$$(4.6) \quad F_i^h F_j^i = -A_j^h, \quad G_i^h G_j^i = -A_j^h, \quad H_i^h H_j^i = -A_j^h,$$

$$(4.7) \quad F_i^h G_j^i + G_i^h H_j^i + H_i^h F_j^i = A_j^h,$$

we say that the manifold admits an *almost adjoint complex structure* and we call such a manifold an *almost adjoint complex manifold*.

In the almost adjoint complex manifold, if there exists a coordinate system (1.4) with respect to which the two tensors among  $F_i^h$ ,  $G_i^h$  and  $H_i^h$  have the components as (4.1), then the manifold is an adjoint complex manifold. In fact, in a domain in which two such coordinate systems  $(z^i)$ ,  $(z^{i'})$  are valid, we have, for  $F_i^h$  and  $G_i^h$ ,

$$(4.9) \quad \frac{\partial z^{i'}}{\partial z^i} F_{i'}^{h'} = \frac{\partial z^{h'}}{\partial z^h} F_i^h, \quad \frac{\partial z^{i'}}{\partial z^i} G_{i'}^{h'} = \frac{\partial z^{h'}}{\partial z^h} G_i^h,$$

from which we obtain

$$\frac{\partial z^{\beta'}}{\partial z^{\alpha}} = 0, \quad \frac{\partial z^{\beta'}}{\partial z^{\infty}} = 0, \quad \frac{\partial z^{\beta'}}{\partial z^{\alpha}} = 0, \quad \frac{\partial z^{\beta'}}{\partial z^{\infty}} = 0, \quad \frac{\partial z^{\infty'}}{\partial z^{\alpha}} = 0, \quad \frac{\partial z^{\infty'}}{\partial z^{\infty}} = 0,$$

and thus, it follows that

$$z^{\alpha'} = z^{\alpha'}(z^{\alpha}), \quad z^{\bar{\alpha}'} = z^{\bar{\alpha}'}(z^{\bar{\alpha}}), \quad z^{\infty'} = z^{\infty'}(z^{\infty}).$$

In this case, we say that the almost adjoint complex structure is induced by an adjoint complex structure.

Next, we find a condition that an almost adjoint complex manifold admits an adjoint complex structure.

If an almost adjoint complex manifold with structure tensors  $F_i^h$ ,  $G_i^h$ ,  $H_i^h$  admits an adjoint complex structure, then

$$(4.10) \quad B_i^h dz^i = 0, \quad C_i^h dz^i = 0, \quad D_i^h dz^i = 0$$

are completely integrable. Conversely, let us assume that (4.10) are completely integrable, where  $B_i^h$ ,  $C_i^h$ ,  $D_i^h$  are defined in the same form by  $F_i^h$ ,  $G_i^h$ ,  $H_i^h$ , and denote the solutions of (4.10) by

$$z^{\alpha'} = z^{\alpha'}(z^i) = \text{const.}, \quad z^{\bar{\alpha}'} = z^{\bar{\alpha}'}(z^i) = \text{const.}, \quad z^{\infty'} = z^{\infty'}(z^i) = \text{const.}$$

respectively. Solving for  $z^i$ , we have

$$z^h = z^h(z^{\alpha'}, z^{\bar{\alpha}'}, z^{\infty'}),$$

and thus,

$$dz^h = \frac{\partial z^h}{\partial z^{\alpha'}} dz^{\alpha'} + \frac{\partial z^h}{\partial z^{\bar{\alpha}'}} dz^{\bar{\alpha}'} + \frac{\partial z^h}{\partial z^{\infty'}} dz^{\infty'}.$$

From  $dz^{\alpha'} = 0$ ,  $dz^{\bar{\alpha}'} = 0$ ,  $dz^{\infty'} = 0$ , since  $dz^h$  satisfies (4.10) respectively, we get

$$(4.11) \quad \frac{\partial z^h}{\partial z^{\bar{\alpha}'}} = -iG_i^h \frac{\partial z^i}{\partial z^{\bar{\alpha}'}} , \quad \frac{\partial z^h}{\partial z^{\infty'}} = -iG_i^h \frac{\partial z^i}{\partial z^{\infty'}} ,$$



$$\begin{aligned}\frac{\partial z^h}{\partial z^{\alpha'}} &= -iF_i^h \frac{\partial z^i}{\partial z^{\alpha'}}, & \frac{\partial z^h}{\partial z^{\infty'}} &= -iF_i^h \frac{\partial z^i}{\partial z^{\infty'}}, \\ \frac{\partial z^h}{\partial z^{\alpha'}} &= -iH_i^h \frac{\partial z^i}{\partial z^{\alpha'}}, & \frac{\partial z^h}{\partial z^{\bar{\alpha}'}} &= -iH_i^h \frac{\partial z^i}{\partial z^{\bar{\alpha}'}}.\end{aligned}$$

from which we have

$$(4.12) \quad \begin{aligned}(G_i^h) &= \begin{pmatrix} \sigma_\beta^\alpha & 0 & 0 \\ 0 & i\delta_\beta^{\bar{\alpha}} & 0 \\ 0 & 0 & i \end{pmatrix}, & (F_i^h) &= \begin{pmatrix} i\delta_\beta^\alpha & 0 & 0 \\ 0 & \tau_\beta^{\bar{\alpha}} & 0 \\ 0 & 0 & i \end{pmatrix}, \\ (H_i^h) &= \begin{pmatrix} i\delta_\beta^\alpha & 0 & 0 \\ 0 & i\delta_\beta^{\bar{\alpha}} & 0 \\ 0 & 0 & \lambda \end{pmatrix}.\end{aligned}$$

From (4.8), we obtain

$$(4.13) \quad \sigma_\beta^\alpha = -i\delta_\beta^\alpha, \quad \tau_\beta^{\bar{\alpha}} = -i\delta_\beta^{\bar{\alpha}}, \quad \lambda = -i,$$

which shows that  $F_i^h$ ,  $G_i^h$  and  $H_i^h$  have the components (4.1) with respect to the coordinate system  $(z^{\alpha'}, z^{\bar{\alpha}'}, z^{\infty'})$ . Thus we have the following:

**THEOREM 3.** *In order that an almost adjoint complex manifold admits an adjoint complex structure, it is necessary and sufficient that (4.10) are completely integrable.*

### § 5. Almost adjoint Hermitian manifold.

In adjoint complex manifold, let us take three operators  $O, O^*$  and  $O^{**}$  such that

$$(5.1) \quad \begin{aligned}O : O_{ji}^{ml} &= \frac{1}{4} (\bar{B}_j^m C_i^l + C_j^m \bar{B}_i^l + B_j^m \bar{C}_i^l + \bar{C}_j^m B_i^l) \\ &= A_j^m A_i^l + \frac{1}{2} (G_j^m F_i^l + F_j^m G_i^l),\end{aligned}$$

$$(5.2) \quad \begin{aligned}O^* : O_{ji}^{*ml} &= \frac{1}{4} (\bar{C}_j^m D_i^l + D_j^m \bar{C}_i^l + C_j^m \bar{D}_i^l + \bar{D}_j^m C_i^l) \\ &= A_j^m A_i^l + \frac{1}{2} (F_j^m H_i^l + H_j^m F_i^l),\end{aligned}$$

$$(5.3) \quad \begin{aligned}O^{**} : O_{ji}^{**ml} &= \frac{1}{4} (\bar{D}_j^m B_i^l + B_j^m \bar{D}_i^l + D_j^m \bar{B}_i^l + \bar{B}_j^m D_i^l) \\ &= A_j^m A_i^l + \frac{1}{2} (H_j^m G_i^l + G_j^m H_i^l).\end{aligned}$$

Then, for any covariant tensor  $g_{ji}$ , they are operated respectively by the forms

$$(5.4) \quad (O_{ji}^{ml} g_{ml}) = \begin{pmatrix} 2g_{\beta\alpha} & 0 & g_{\beta\infty} \\ 0 & 2g_{\bar{\beta}\alpha} & g_{\bar{\beta}\infty} \\ g_{\infty\alpha} & g_{\infty\bar{\alpha}} & 0 \end{pmatrix},$$

$$(5.5) \quad (O^*_{ji}{}^{ml} g_{ml}) = \begin{pmatrix} 0 & g_{\beta\alpha} & g_{\beta\infty} \\ g_{\bar{\beta}\alpha} & 2g_{\bar{\beta}\alpha} & 0 \\ g_{\infty\alpha} & 0 & 2g_{\infty\infty} \end{pmatrix},$$

$$(5.6) \quad (O^{**}_{ji}{}^{ml} g_{ml}) = \begin{pmatrix} 2g_{\beta\alpha} & g_{\beta\alpha} & 0 \\ g_{\bar{\beta}\alpha} & 0 & g_{\bar{\beta}\infty} \\ 0 & g_{\infty\bar{\alpha}} & 2g_{\infty\infty} \end{pmatrix},$$

from which we have

$$(5.7) \quad O \cdot O^* \cdot O^{**}(g) = (0).$$

If an almost adjoint complex manifold has a Riemannian metric  $ds^2 = g_{ji} dz^j dz^i$  satisfying

$$(5.8) \quad O_{ji}^{ml} g_{ml} = 0,$$

or

$$(5.9) \quad g_{ji} = -g_{ml} F_{(j}^m G_{i)},$$

where the bracket denotes the symmetric part of the indices, then the manifold is called an *almost adjoint Hermitian manifold* and we call this metric an *almost adjoint Hermitian metric*.

From this definition, we have the following instantly:

**THEOREM 4.** *In an adjoint complex manifold  $X_{2n+1}$ , if it admits an almost adjoint Hermitian metric, then it is an adjoint Hermitian manifold, and conversely.*

In fact, it follows that, from (5.4) and (5.8),

$$g_{\beta\alpha} = g_{\bar{\beta}\alpha} = 0, \quad g_{\beta\infty} = g_{\bar{\beta}\infty} = 0,$$

and, since  $ds^2$  is real,  $g_{ji}$  is adjoint self-conjugate.

And, we obtain the following theorem is valid:

**THEOREM 5.** *In an adjoint complex manifold, it is always possible to define an adjoint Hermitian metric.*

In fact, let  $a_{ji}$  be a tensor which defines a Riemannian metric in an adjoint complex manifold, and let us put

$$(5.10) \quad g_{ji} = O^{*ml}_{ji} O^{*ts}_{ml} a_{ts},$$

then  $g_{ji}$  defines another positive definite Riemannian metric and satisfies (5.8).

### § 6. Almost adjoint Kaehlerian manifold.

In an almost adjoint Hermitian manifold  $X_{2n+1}$ , we define the tensors  $G_{ji}$ ,  $F_{ji}$ , and  $H_{ji}$  as

$$(6.1) \quad G_{ji} = g_{jh} G_i^h, \quad F_{ji} = g_{jh} F_i^h, \quad H_{ji} = g_{jh} H_i^h,$$

and  $G_{kji}$ ,  $F_{kji}$  and  $H_{kji}$  as

$$(6.2) \quad \begin{aligned} G_{kji} &= \nabla_k G_{ji} + \nabla_j G_{ik} + \nabla_i G_{kj}, \\ F_{kji} &= \nabla_k F_{ji} + \nabla_j F_{ik} + \nabla_i F_{kj}, \\ H_{kji} &= \nabla_k H_{ji} + \nabla_j H_{ik} + \nabla_i H_{kj}, \end{aligned}$$

In an adjoint Hermitian manifold, let us calculate the above tensors. Form (1.9) and (4.1), we have

$$(6.3) \quad \begin{aligned} (G_{ji}) &= \begin{pmatrix} 0 & ig_{\beta\bar{\alpha}} & 0 \\ -ig_{\bar{\beta}\alpha} & 0 & 0 \\ 0 & 0 & ig_{\infty\infty} \end{pmatrix}, & (F_{ji}) &= \begin{pmatrix} 0 & -ig_{\beta\bar{\alpha}} & 0 \\ ig_{\bar{\beta}\alpha} & 0 & 0 \\ 0 & 0 & ig_{\infty\infty} \end{pmatrix}, \\ (H_{ji}) &= \begin{pmatrix} 0 & ig_{\beta\bar{\alpha}} & 0 \\ ig_{\bar{\beta}\alpha} & 0 & 0 \\ 0 & 0 & -ig_{\infty\infty} \end{pmatrix}. \end{aligned}$$

And we can see that for the tensors  $G_{kji}$ ,  $F_{kji}$  and  $H_{kji}$ , their essential components are

$$(6.4) \quad G_{\gamma\beta\alpha}, G_{\gamma\beta\bar{\alpha}}, G_{\gamma\beta\infty}, G_{\gamma\bar{\beta}\alpha}, G_{\gamma\bar{\beta}\infty}, \\ G_{\gamma\infty\infty}, G_{\bar{\gamma}\bar{\beta}\alpha}, G_{\bar{\gamma}\bar{\beta}\infty}, G_{\gamma\infty\infty}, G_{\infty\infty\infty},$$

and  $F_{kji}$  and  $H_{kji}$  are similar. In consequence that  $G_{kji}$  are represented by

$$(6.5) \quad G_{kji} = \partial_k G_{ji} + \partial_j G_{ik} + \partial_i G_{kj} \\ - \Gamma_{kj}^m (G_{mi} + G_{im}) - \Gamma_{ji}^m (G_{mk} + G_{km}) - \Gamma_{ik}^m (G_{mj} + G_{jm}),$$

and the others are similar. Using of (2.4) and (6.3) we obtain the followings:

$$(6.7) \quad G_{\gamma\beta\alpha} = 0, \quad G_{\gamma\beta\bar{\alpha}} = i(\partial_\gamma g_{\beta\bar{\alpha}} - \partial_\beta g_{\gamma\bar{\alpha}}) \\ G_{\gamma\beta\infty} = 0, \quad G_{\gamma\bar{\beta}\alpha} = i(\partial_\alpha g_{\gamma\bar{\beta}} - \partial_{\bar{\beta}} g_{\gamma\alpha}), \\ G_{\gamma\bar{\beta}\infty} = 2i \partial_\infty g_{\gamma\bar{\beta}}, \quad G_{\gamma\infty\infty} = -i \partial_\gamma g_{\infty\infty}, \\ G_{\bar{\gamma}\bar{\beta}\alpha} = 0, \quad G_{\bar{\gamma}\bar{\beta}\infty} = 0, \\ G_{\bar{\gamma}\infty\infty} = -i \partial_{\bar{\gamma}} g_{\infty\infty}, \quad G_{\infty\infty\infty} = 0,$$

$$(6.8) \quad F_{\gamma\beta\alpha} = 0, \quad F_{\gamma\beta\bar{\alpha}} = -i(\partial_\gamma g_{\beta\bar{\alpha}} - \partial_\beta g_{\gamma\bar{\alpha}}), \\ F_{\gamma\beta\infty} = 0, \quad F_{\gamma\bar{\beta}\alpha} = -i(\partial_\alpha g_{\gamma\bar{\beta}} - \partial_{\bar{\beta}} g_{\gamma\alpha}), \\ F_{\gamma\bar{\beta}\infty} = 2i \partial_\infty g_{\gamma\bar{\beta}}, \quad F_{\gamma\infty\infty} = -i \partial_\gamma g_{\infty\infty}, \\ F_{\bar{\gamma}\bar{\beta}\alpha} = 0, \quad F_{\bar{\gamma}\bar{\beta}\infty} = 0, \\ F_{\bar{\gamma}\infty\infty} = -i \partial_{\bar{\gamma}} g_{\infty\infty}, \quad F_{\infty\infty\infty} = 0,$$

$$(6.9) \quad H_{\gamma\beta\alpha} = 0, \quad H_{\gamma\beta\bar{\alpha}} = 0, \quad H_{\gamma\beta\infty} = 0, \\ H_{\gamma\bar{\beta}\alpha} = 0, \quad H_{\gamma\bar{\beta}\infty} = 0, \quad H_{\gamma\infty\infty} = i \partial_\gamma g_{\infty\infty}, \\ H_{\bar{\gamma}\bar{\beta}\alpha} = 0, \quad H_{\bar{\gamma}\bar{\beta}\infty} = 0, \quad H_{\bar{\gamma}\infty\infty} = i \partial_{\bar{\gamma}} g_{\infty\infty}, \\ H_{\infty\infty\infty} = 0.$$

Form these and the adjoint Kaehlerian condition (2.7), we can see that  $G_{kji}=0$  or  $F_{kji}=0$  is necessary and sufficient in order that an adjoint Hermitian manifold be an adjoint Kaehlerian.

If, in an almost adjoint Hermitian manifold, it satisfies

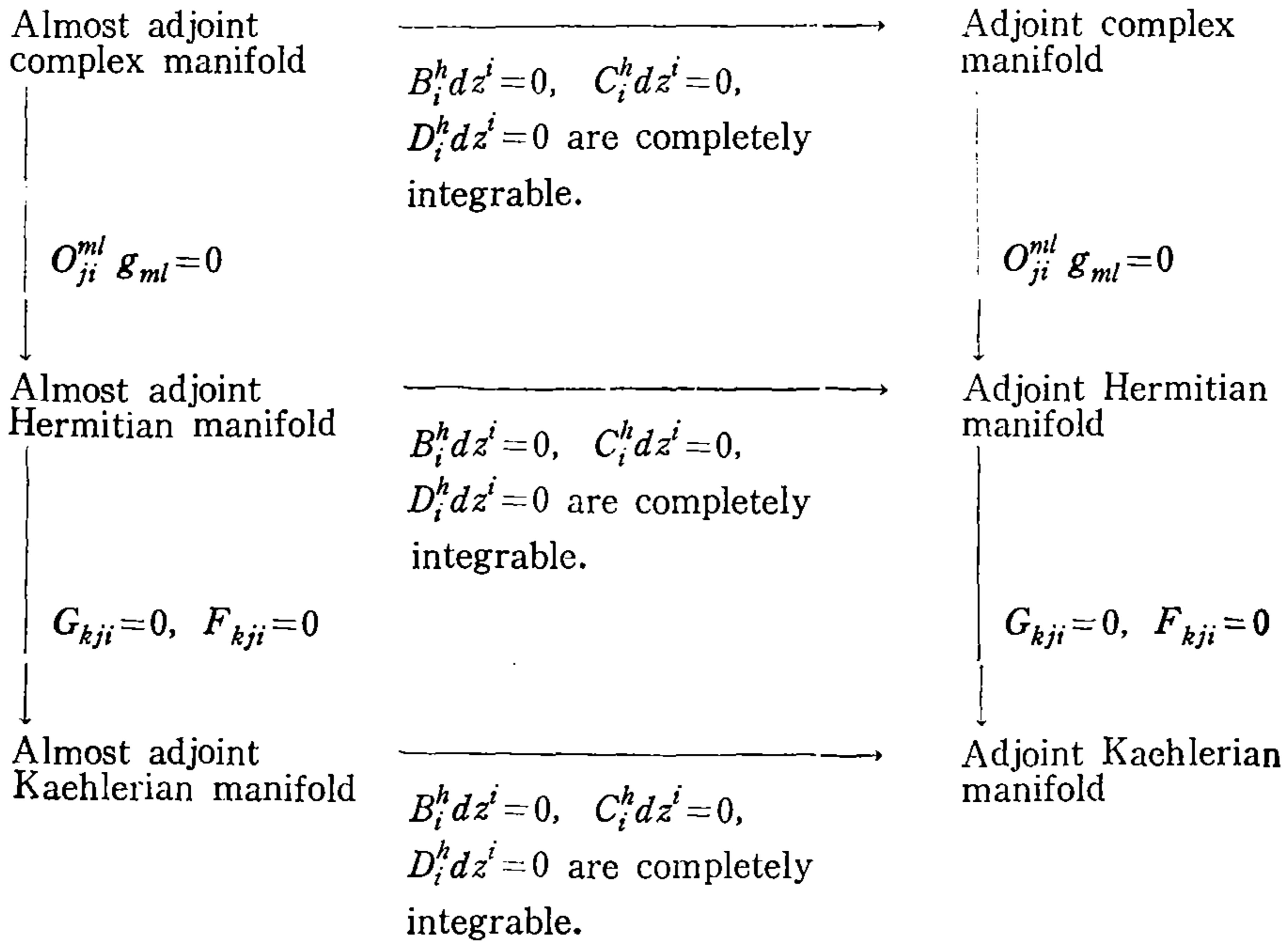
$$(6.10) \quad G_{kji}=0, \quad F_{kji}=0,$$

we say that the manifold admits an *almost adjoint Kaehlerian structure* and we call it an *almost adjoint Kaehlerian manifold*. And thus we have the following:

THEOREM 6. *If an adjoint Hermitian manifold admits an almost adjoint Kählerian structure, then it is an adjoint Kählerian manifold, and conversely.*

§ 7. Diagram.

Summarizing the theorems in the present paper, we obtain the following diagram:



July, 1964  
 Mathematical Department  
 Kyungpook University  
 Taegu, Korea

**REFERENCES**

- [1] K. Yano; *The theory of Lie derivatives and its applications*, Amsterdam, (1955).
- [2] K. Yano; *On analytical vectors*, Sugaku, Vol.8, No.4, (1957).
- [3] K. Yano and S. Bochner; *Curvature and Betti Numbers*, Princeton, (1953).