ACTIVE GALACTIC NUCLEUS INTERACTION WITH THE HOT GAS ENVIRONMENT: UNDERSTANDING FROM THE RADIO AND X-RAY DATA

DHARAM V. LAL
National Centre for Radio Astrophysics (TIFR), Pune University Campus, Ganeshkhind P.O., Pune 411 007, India
E-mail: dharam@ncra.tifr.res.in
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ABSTRACT

Recognition of the role of radio galaxies in the universe has been increasing in recent years. Their colossal energy output over huge volumes is now widely believed to play a key role not only in the formation of galaxies and their supermassive black holes, but also in the evolution of clusters of galaxies and, possibly, the cosmic web itself. In this regard, we need to understand the interaction of radio bubbles in the hot gas atmospheres of clusters and the importance of the role that radio galaxies play in the overall energy budget of the intracluster medium. Here, we present results from X-ray and radio band observations of the hot gas atmospheres of powerful, nearby radio galaxies in poor clusters.

Key words: galaxies: individual (3C 31, 3C 270.1, 3C 449) – hydrodynamics – intergalactic medium – X-rays: galaxies: clusters

1. INTRODUCTION

All galaxies and quasars appear to be sources of radio emission at some level. Normal spiral galaxies such as our own galactic system are near the low end of the radio luminosity function and have radio luminosities near \(10^{37}\) erg s\(^{-1}\). Some Seyfert galaxies, starburst galaxies, and the nuclei of active elliptical galaxies are 100 to 10\(^3\) times more luminous. Radio galaxies (and some quasars), where the radio emission is powered by active galactic nucleus (AGN) activity rather than star formation, are powerful radio sources at the high end of the luminosity function with radio band luminosities up to \(10^{45}\) erg s\(^{-1}\) (Peterson, 1997).

Extended radio galaxies constitute the largest known physical structures in the universe. Their energy content is very large, \(\sim10^{60}\) erg. Several attempts at understanding the origin of this (accretion) energy and the manner in which it is converted into relativistic particles and magnetic fields have been made, including the ‘spin paradigm’ (Blandford & Znajek 1977). High-resolution radio images generally show a very compact component which is coincident with an AGN, called the ‘central engine’. Long collimated ‘jets’ extend away from the compact central core toward the outer radio lobes; in some sources these jets eventually end up terminating at ‘hotspots’. The primary observables for a radio galaxy at a known redshift are, for the purpose of this article, to be regarded as the morphology and the luminosity at some common rest frame frequency derived from the spectral shape and the redshift in some chosen cosmological model.

The study of radio galaxies has evolved somewhat unsystematically, e.g., complex classification schemes, objects which were once thought to be different are now being found to be related, etc. Attempts at morphological classification and statistical studies include ‘head-tail radio galaxies’ – (Miley et al. 1972, Fanaroff & Riley 1974) a classification scheme for strong and weak radio galaxies with powerful, edge-brightened (with highly collimated jets and compact ‘hotspots’) structure and less powerful, centre-brightened structure (with less collimated jets), respectively, being FR class I and class II radio galaxies.

1.1. Low-/High-excitation Radio Galaxies

In addition, with the advent of detailed X-ray observations using Chandra and XMM-Newton, and well established optical surveys, clear evidence for a fundamental dichotomy related to the accretion onto the central supermassive black hole, called ‘high-excitation’ and ‘low-excitation’ modes, is also seen. Briefly, Table 1 (reproduced from Garofalo, Evans & Sambruna 2010) summarises the broad properties of low- and high-excitation radio galaxies.

1.2. Emission Processes

The energetics of radio sources are important not only because of their relation to the emission process but because they also play an important role in all considerations of how radio sources are held together or confined. The emission in the radio band from radio galaxies is synchrotron emission, as inferred from its very
Table 1

<table>
<thead>
<tr>
<th>Definition</th>
<th>FR classification</th>
<th>X-ray spectra</th>
<th>Accretion flow type</th>
<th>Optical constraints</th>
</tr>
</thead>
<tbody>
<tr>
<td>No narrow optical line emission.</td>
<td>Almost all FRIs are LERGs at (z\approx0).</td>
<td>Jet-related unabsorbed power law only. Upper limits only to ‘hidden’ accretion-related emission.</td>
<td>Highly sub-Eddington. Likely radiatively inefficient.</td>
<td>Strong radio/optical/soft X-ray correlations.</td>
</tr>
<tr>
<td>Prominent optical emission lines, either narrow (NLRG) or broad (BLRG), or quasar.</td>
<td>Significant population of FRIs at (z\approx0.5).</td>
<td>Jet-related unabsorbed power law + significant accretion contribution (heavily absorbed in NLRGs).</td>
<td>Reasonable fraction of Eddington. Likely standard accretion disk.</td>
<td>Optical emission is jet-related.</td>
</tr>
<tr>
<td>Low-excitation radio galaxy (LERG)</td>
<td>Most FRIs are HERGs, as are a handful of FRIs (e.g., Cen A).</td>
<td>Jet-related unabsorbed power law</td>
<td>Jet-related unabsorbed power law</td>
<td>Strong radio/optical/soft X-ray correlations.</td>
</tr>
<tr>
<td>High-excitation radio galaxy (HERG)</td>
<td></td>
<td>+ significant accretion contribution (heavily absorbed in NLRGs).</td>
<td></td>
<td>Optical emission is jet-related.</td>
</tr>
</tbody>
</table>

smooth, broad-band nature and strong polarization. In the X-ray band, the population of relativistic electrons that gives rise to the radio emission can also produce X-rays. There are two methods: 1) the synchrotron process and 2) the inverse-Compton process. The former simply states that the synchrotron process produces not only radio emission, but also X-ray emission. In the latter, the relativistic electrons in the radio lobes, via the inverse-Compton process, scatter the cosmic microwave background (CMB) photons with frequency \(\nu\)~10^{13}(1+z) (\(z\) is redshift) and boost them up to a frequency \(\nu_{\text{inverse-Compton}}\approx\gamma^2\nu\) (where \(\gamma\) is the Lorentz factor), i.e., into the X-ray and gamma-ray regions. Since the CMB energy density \(u_{\text{CMB}}=4.1\times10^{-13}(1+z)^4\) ergs s^{-1} is precisely known, a comparison of the synchrotron radio emission and the inverse-Compton X-ray flux gives estimates of energy densities in the magnetic field in radio lobes and lets us examine energy equipartition between particles and fields.

High resolution imaging using Chandra has shown X-ray emission from the jets of radio galaxies, which is interpreted as synchrotron emission. Synchrotron emission is often observed from jets of weak radio galaxies, such as Centaurus A (Hardcastle et al., 2007) and M 87 (Perlman & Wilson, 2005), as well as from some terminal hotspots in powerful radio galaxies (Hardcastle et al., 2007). Similarly, Chandra has also shown that the hotspots of several radio galaxies, e.g. Cygnus A, while undetected at optical wavelengths, radiate significantly in the X-ray (Wilson et al., 2000), which is interpreted as evidence that some hotspots emit X-rays by a synchrotron self-Compton process. Finally, extended X-ray emission from the radio lobes of powerful, FR II radio galaxies and quasars is thought to be produced by inverse-Compton scattering off the CMB (Feigelson et al., 1995; Hardcastle et al., 2002a, 2005).

2. RADIO AND X-RAY STUDY OF RADIO GALAXIES

Therefore, X-ray observations of radio galaxies can provide important clues about the dynamics and particle acceleration in the relativistic jets and outflows. Or in other words, X-ray imaging has opened new vistas on the structure and evolution of groups and clusters and their member galaxies. The aim of this paper is to provide three examples to understand the nature of X-ray emission from three parts of radio galaxies. These three aspects reviewed in this contribution are

- the nature of X-ray emission from a radio jet close to the core of the weak FR I, 3C 31 radio galaxy (Hardcastle et al., 2002b; Lanz et al., 2011),
- the nature of X-ray emission from the radio lobes of the strong FR II radio galaxy, 3C270.1 (Wilkes et al., 2012) and
- the nature of the hot gas around the canonical FR I radio galaxy 3C 449 (Lal et al., 2013).

3. RADIO GALAXY

3.1. Jet Emission

Our first example is a case study of 3C 31. It is a nearby twin-jet radio source hosted within the early-type galaxy NGC 383 at a redshift of 0.0168. The twin jets are aligned roughly along a north-south axis. The northern jet is straight, narrow and brighter in the central 20 kpc, beyond this distance the jet 
ares. We clearly detect X-ray emission using Chandra, from the jet component close to the radio core in 3C 31. The presence of X-ray emission suggests that there is on-going particle acceleration in these regions. Often radio, infrared, optical, and X-ray emission from FR I jets can be modeled as synchrotron emission from a single electron energy distribution described by a broken power law. In this example, we model the broadband spectra of 3C 31 with clearly measured jet flux densities in radio, infrared, optical and X-ray in the region closest to the nucleus (please refer Figure 4 of Lanz et al. 2011). The broken power-law model of synchrotron emission from a broken power-law electron energy distribution gives an acceptable fit to the data. Therefore, here again, radio through X-ray spectra in these regions demonstrate that the emission, which can be interpreted as synchrotron emission from a broken power-law distribution of electron energies (Hardcastle et al. 2002b; Lanz et al. 2011).
3.2. Radio Lobe Emission

The radio galaxy 3C 270.1 at a redshift of 1.532 has the double-lobed radio structure characteristic of FR II. We model the unperturbed, static and weak but significant X-ray emission associated with the southeastern extension of the northern radio lobe, called the ‘N-C-jet-bend’ feature (please refer Figure 1 of Wilkes et al., 2012) detected using Chandra. The emission can be explained by the predictions of inverse-Compton scattering of CMB photons, can be brought into agreement with the observations if the magnetic field is 3 ±1 nT. This field is a factor of ~7–10 lower than the equipartition magnetic field, which is a larger than departure (Croston et al., 2005).

Here, one should note that the low energy end of the spectrum of the relativistic particles in radio galaxies is essentially unconstrained. This is true for the 3C 270.1 radio source and other sources in general, and the effect is more severe for high-redshift radio galaxies than for low-redshift radio galaxies. In the example presented above, the electron energy required to scatter a photon to X-ray energies depends on the initial energy of the photon: if the photon is one of the radio photons emitted by the electron population itself, then Lorentz factors of 10^4 are required, whereas if the photon comes from the all-pervading CMB radiation, Lorentz factors of 10^3 are required. The key fact about the inverse-Compton process is that the energy radiated in scattered photons (for a given electron and photon energy) is the product of the number density of electrons and photons, hence, the observations of inverse-Compton emission are always combined with the constraints from the synchrotron emission. However, it is difficult to constrain the synchrotron emission for Lorentz factors below 100–10^3 and often one is forced to make assumptions, e.g. a flattening or a low energy cut-off in the spectrum at these energies.

3.3. Radio Galaxy Interaction with the Hot Gas Environment

Although X-ray observations of several radio galaxies exists in a variety of environments including cluster outskirts and the field, here we present an example of the canonical FR I radio galaxy 3C 449. 3C 449 exhibits a complex radio morphology on spatial scales from parsecs to hundreds of kiloparsecs. The source is characterised by an unresolved core, two symmetrically opposed jets, and very extended radio lobes. Each jet is remarkably straight from the nucleus before deviating sharply towards the west, dropping in surface brightness and terminating in large diffuse radio lobes. Since both jets bend in the same direction, rotation effects cannot cause this morphology (Birkinshaw et al., 1981).

Several panels in Figure 1 show following features:

- There is an elongated diffuse structure extending ~110 kpc to the northeast of the core, labeled 'cool-filament'.
- There are two arc-shaped tangential edges in the surface brightness ~33 kpc to the southeast and to the west of the core.
- There is an X-ray cavity, i.e. a decrement in X-ray emission, at the location of the southern inner radio lobe. This cavity is surrounded by a rim of enhanced brightness.
- There is a tunnel-like feature in the X-ray emission connecting the southern cavity and the group core, marked with a dashed line, which coincides with the southern radio jet.

We further discuss one key feature, the presence of surface brightness edges mentioned above. A careful treatment of each of these surface brightness discontinuities suggests that the gas temperature and density change discontinuously but the pressure typically remains at or near equilibrium across the front, hence, these are probably shearing cold fronts due to a merger (Lal et al., 2013). Additionally, the straight inner jet flares at approximately the position where it crosses the contact edge, suggesting that the jet is entraining and thermalising some of the hot gas as it crosses the edge. This is consistent with hydrodynamic simulations of jets crossing density discontinuities with shear flow which show that well collimated jets can become partially disrupted (Loken et al., 1995); the disruption of the jets of 3C 449 by this type of edge was first predicted by Katz-Stone & Rudnick (1997). It is therefore not unreasonable to interpret the fact that jets crossing density edges can become partially disrupted and inflate to form inner radio lobes as a consequence of observed signatures of interactions between an AGN and its surrounding medium (Lal et al., 2013).

4. LOOKING AHEAD

The statistics, i.e. general conclusions about the population as a whole are already available in the literature from optical surveys and also to a large extent from radio surveys, essentially only at these two key wavebands. This approach of drawing general conclusions about the population as a whole is lacking in the X-ray band. Some of the results presented here are primarily from the published results of observations of interesting individual objects.

In the meantime, hopefully, the low-frequency array (LOFAR), which is already making enormous progress, is expected to provide a major step in our understanding of the physics of radio sources. LOFAR will reach unprecedented sensitivities and resolution in the frequency range 40-240 MHz and this will immediately allow us to study the electron spectrum of Lorentz factors of ~10^3, as discussed in Sec. 3.2. This is an appreciable step forward, however it is still not enough to catch the bulk of the particle energy in radio sources and to constrain the physics of particle acceleration in these sources. Furthermore, in terms of current X-ray observing facilities, Chandra and XMM-Newton will not last for ever, their immediate replacements, including Athena1 and SMART-X2 though they are still at the design stage, may become a reality before current facilities fade away. It thus seems that the future is bright, and radio and

1https://www.mpe.mpg.de/Athena
2http://smart-x.cfa.harvard.edu/index.html
X-ray band observations will continue to play an important role for radio galaxies.

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REFERENCES


Figure 1. Left panel: Combined Chandra/ACIS-S counts image of 3C 449 in the 0.5-2.0 kilo-electronvolts (keV) band. The image is binned to 4' pixels and smoothed with a Gaussian ($\sigma = 6''$). All point sources, other than the active nucleus of 3C 449, have been removed. We see diffuse thermal emission in the soft band extending over the entire field of view, 12' ($\approx$ 240 kpc) along the east-west axis (Lal et al 2013). Middle panel: Background-subtracted, exposure-corrected, Gaussian-smoothed ($\sigma = 4''$) Chandra image in the energy range 0.5-2.0 keV. Features of interest are labeled including lopsided gas emission near core, the cool filament, the eastern and western edges, the cavity associated with the southern inner radio lobe, and the tunnel associated with the southern inner jet. Right panel: Projected thermodynamic temperature map using the merged data. The temperature map has radio contours overlaid. The color bar gives the temperature scale in keV.